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# Recovery Techniques For Finite Element Methods And Their Applications

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**RECOVERY TECHNIQUES FOR FINITE ELEMENT METHODS  
AND THEIR APPLICATIONS**

by

**HAILONG GUO**

**DISSERTATION**

Submitted to the Graduate School

of Wayne State University,

Detroit, Michigan

in partial fulfillment of the requirements

for the degree of

**DOCTOR OF PHILOSOPHY**

2015

MAJOR: MATHEMATICS

Approved by:

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Advisor

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Date

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# DEDICATION

*To My Parents*

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## CHAPTER 1 Introduction

Post-processing is an important technique in scientific computing, where it is necessary to draw some useful information that have physical meanings such as velocity, flux, stress, etc., from the primary results of the computing. These quantities of interest usually involve derivatives of the primary data. This type of post-processing methods are called recovery techniques. Typical recovery techniques include gradient recovery and Hessian recovery.

As for gradient recovery, it is extensively studied in the literature. Most of them are based on averaging methods [19, 63], local or global projections [12, 52, 90], or local least square fittings [58, 79, 107, 110, 111, 112]. However, all the methods above are only limited to  $C^0$  finite element methods. According to [4], a good recovery operator should satisfy consistency condition, localization condition, and boundedness and linearity condition. For conforming elements, it is well known that polynomial preserving recovery (PPR) [107] satisfies all the above conditions. Unfortunately, there is no such type gradient recovery for the Crouzeix-Raviart element in the literature. To bridge the gap, we propose and analyze a gradient recovery method for the Crouzeix-Raviart element in this work.

Hessian matrix is particularly significant in adaptive mesh design, since it can indicate the direction where the function changes the most and guide us to construct anisotropic meshes to cope with the anisotropic properties of the solution of the underlying partial differential equation [59]. It also plays an important role in finite element approximation of second order non-variational elliptic problems [65], numerical solution of some fully nonlinear equations such as Monge-Ampère equation [66, 83], and designing nonlocal finite element technique [45]. There have been some works in literature on this subject. In 1998, Lakhany-Whiteman



used a simple averaging method twice at edge centers of the regular uniform triangular mesh to produce a superconvergent Hessian [64]. Later, some other researchers such as Agouzal et al.[2], Bank et al.[13] and Owall [88] also studied Hessian recovery. Comparison studies of existing Hessian recovery techniques can be found in Vallet et al. [94] and Picasso et al. [89]. However, there is no systematic theory guaranteeing convergence under general circumstances. Moreover, there are certain technical difficulties in obtaining rigorous convergence proof for meshes other than the regular pattern triangular mesh. In a very recent work, Kamenski-Huang argued that it is not necessary to have very accurate or even convergent Hessian in order to obtain a good mesh [61]. Our goal is not targeted on the direction of adaptive mesh refinement; instead, we emphasize on obtaining accurate Hessian matrices via recovery techniques. We propose an effective Hessian recovery method and establish a solid theoretical analysis for such method.

The applications of recovery techniques include adaptive finite element methods and eigenvalue problems. In this work, we discuss the application of the proposed gradient recovery for the Crouzeix-Raviart element into adaptive nonconforming finite element method. Also, we apply the polynomial preserving recovery into efficient eigenvalue computation to propose several superconvergent two-grid schemes and multilevel adaptive methods for elliptic eigenvalue problems.

The rest of this dissertation is organized as follows:

Chapter 2 begins with introducing some basic notation of Sobolev spaces and model problems. Then we review some preliminary knowledges of finite element methods and their approximations of the model problems.

Chapter 3 focuses on recovery of the first order derivative for the Crouzeix-Raviart ele-

ment. In this chapter, we propose and analyze a gradient recovery method for the Crouzeix-Raviart element. The proposed method is proved to preserving quadratic polynomials and to be a bounded linear operator. Its application in adaptive nonconforming finite element is also discussed. This chapter is based on our published paper [49].

Chapter 4 is devoted to recovery of the second order derivative of Lagrange element of arbitrary order. A new Hessian recovery strategy is proposed and its mathematical theory is established. This chapter is based on our submitted paper [50].

Chapter 5 concentrates on application of recovery technique. In partial, we employ polynomial preserving recovery to design some fast and efficient algorithms for eigenvalue problems. Our new algorithms compare favorably with some existing algorithms and enjoy superconvergence property. This chapter is based on our submitted paper [51].

## CHAPTER 2 Preliminaries

### 2.1 Sobolev spaces

In this work, we assume  $\Omega \subset \mathbb{R}^2$  is a bounded polygonal domain with Lipschitz continuous boundary  $\partial\Omega$ . Let  $\int_{\Omega} f(z)dz$  denote the Lebesgue integral [1, 25, 35] for some Lebesgue measure function  $f$  on the domain  $\Omega$ . For  $1 \leq p \leq \infty$ , let

$$\|f\|_{p,\Omega} = \left( \int_{\Omega} |f(z)|^p dz \right)^{1/p},$$

and for the case  $p = \infty$ ,

$$\|f\|_{0,\infty,\Omega} = \text{ess sup}\{|f(z) : z \in \Omega\}.$$

Define the Lebesgue spaces

$$L^p(\Omega) = \{f : \|f\|_{p,\Omega} < \infty\}.$$

A multi-index is a 2-tuple of non-negative integers  $\alpha_i, i = 1, 2$ . The length of  $\alpha$  is given by

$$|\alpha| = \sum_{i=1}^2 \alpha_i.$$

The weak partial derivative  $D^\alpha v$ , see [1, 25, 35], is then defined as

$$D^\alpha v = \left( \frac{\partial}{\partial x} \right)^{\alpha_1} \left( \frac{\partial}{\partial y} \right)^{\alpha_2}.$$

Also,  $D^k u$  with  $|\alpha| = k$  is the vector of all partial derivatives of order  $k$ . Let  $W^{k,p}(\Omega) = \{v :$

$D^\alpha v \in L^p(\Omega), |\alpha| \leq k\}$  be the classical Sobolev spaces with norms

$$\|v\|_{k,p,\Omega} = \left( \sum_{|\alpha| \leq k} \int_{\Omega} |D^\alpha v(z)|^p dz \right)^{\frac{1}{p}}, \quad 1 \leq p < \infty,$$

$$\|v\|_{k,\infty,\Omega} = \operatorname{ess\,sup}_{|\alpha| \leq k, z \in \Omega} |D^\alpha v(z)|, \quad p = \infty;$$

and seminorms

$$|v|_{k,p,\Omega} = \left( \sum_{|\alpha|=k} \int_{\Omega} |D^\alpha v(z)|^p dz \right)^{\frac{1}{p}}, \quad 1 \leq p < \infty,$$

$$|v|_{k,\infty,\Omega} = \operatorname{ess\,sup}_{|\alpha|=k, z \in \Omega} |D^\alpha v(z)|, \quad p = \infty.$$

When  $p = 2$ , let  $H^k(\Omega) = W^{k,2}(\Omega)$  and the index  $p$  is omitted in their corresponding norms and seminorms.

For any positive integer  $n$ , we say a vector function  $\vec{v} = (v_1, v_2, \dots, v_n)^T \in W^{k,p}(\Omega)^n$  provided that  $v_i \in W^{k,p}(\Omega)$  for  $i = 1, 2, \dots, n$ . Its norm is defined as

$$\|\vec{v}\|_{k,p,\Omega} = \left( \sum_{i=1}^n \|v_i\|_{k,p,\Omega}^p \right)^{\frac{1}{p}}, \quad 1 \leq p < \infty,$$

$$\|\vec{v}\|_{k,\infty,\Omega} = \operatorname{ess\,sup}_{|\alpha| \leq k, z \in \Omega, 1 \leq i \leq n} |D^\alpha v_i(z)|, \quad p = \infty;$$

and similarly the seminorm of  $\vec{v}$  is defined as

$$|\vec{v}|_{k,p,\Omega} = \left( \sum_{i=1}^n |v_i|_{k,p,\Omega}^p \right)^{\frac{1}{p}}, \quad 1 \leq p < \infty,$$

$$|\vec{v}|_{k,\infty,\Omega} = \operatorname{ess\,sup}_{|\alpha|=k, z \in \Omega, 1 \leq i \leq n} |D^\alpha v_i(z)|, \quad p = \infty.$$

The subscript  $p$  is omitted in case of  $p = 2$ .

For a subdomain  $\mathcal{A}$  of  $\Omega$ , let  $\mathbb{P}_m(\mathcal{A})$  be the space of polynomials of degree less than or equal to  $m$  over  $\mathcal{A}$  and  $n_m$  be the dimension of  $\mathbb{P}_m(\mathcal{A})$  with  $n_m = \frac{1}{2}(m+1)(m+2)$ . Similarly, we use  $W^{k,p}(\mathcal{A})$  to denote the restriction of classical Sobolev space on  $\mathcal{A}$  with norm  $\|\cdot\|_{k,p,\mathcal{A}}$  and seminorm  $|\cdot|_{k,p,\mathcal{A}}$ . When  $p = 2$ , we simply denote  $H^k(\mathcal{A}) = W^{k,2}(\mathcal{A})$  and the subscript  $p$  is omitted.

Throughout this article, the letter  $C$  or  $c$ , with or without subscript, denotes a generic constant which is independent of  $h$  and may not be the same at each occurrence. To simplify notation, we denote  $x \leq Cy$  by  $x \lesssim y$ .

## 2.2 Model problems and their variation problems

In the work, we consider both second order elliptic equation and Stokes equation.

### 2.2.1 Elliptic equation

Our first model problem will be the following homogeneous elliptic equation

$$\begin{cases} -\nabla(\mathcal{D}\nabla u) + cu = f, & \text{in } \Omega, \\ u = 0, & \text{on } \partial\Omega; \end{cases} \quad (2.1)$$

where  $D$  is a  $2 \times 2$  symmetric positive definite matrix and  $c$  as well as  $f$  are scalars.

The variational form is to find  $u \in H_0^1(\Omega)$  such that

$$\mathcal{B}(u, v) = L(v), \quad \forall v \in H_0^1(\Omega). \quad (2.2)$$

where

$$\begin{aligned}\mathcal{B}(u, v) &= \int_{\Omega} [(\mathcal{D}\nabla u)\nabla v + cuv]dz, \\ L(v) &= \int_{\Omega} fvdz.\end{aligned}$$

We assume that  $\mathcal{D}$  and  $f$  in  $L^\infty(\Omega)$ . In addition, assume there exist two constant  $\bar{c}$  and  $\underline{c}$  such that  $0 < \underline{c} \leq c(z) \leq \bar{c} < \infty$ . Then it is to check that the bilinear form  $\mathcal{B}$  and linear functional  $L$  satisfies the following three conditions:

1.  $\mathcal{B}$  is continuous, i.e.

$$\mathcal{B}(u, v) \lesssim \|u\|_{1,\Omega} \|v\|_{1,\Omega}, \quad \forall u, v \in H_0^1(\Omega);$$

2.  $\mathcal{B}$  is coercive, i.e.

$$\|u\|_{1,\Omega}^2 \lesssim \mathcal{B}(u, u), \quad \forall u \in H_0^1(\Omega);$$

3.  $L$  is continuous, i.e.

$$|L(v)| \lesssim \|v\|_{1,\Omega}, \quad \forall v \in H_0^1(\Omega).$$

According to Lax-Milgram Lemma [25, 35], variational problem (2.2) has an unique solution. Furthermore, the above conditions 1 and 2 implies  $\mathcal{B}$  is an inner production on  $H_0^1(\Omega)$ . Define the energy norm as  $|||\cdot|||_{\Omega} = \sqrt{\mathcal{B}(\cdot, \cdot)}$ . Then  $|||\cdot|||_{\Omega}$  and  $\|\cdot\|_{1,\Omega}$  are two equivalent norms in  $H_0^1(\Omega)$ .

### 2.2.2 Stokes equation

Our second model problem is the following Stokes equation

$$\begin{cases} -\Delta \vec{u} + \nabla p = \vec{f}, & \text{in } \Omega, \\ \operatorname{div} \vec{u} = 0, & \text{in } \Omega, \\ \vec{u} = 0, & \text{on } \partial\Omega; \end{cases} \quad (2.3)$$

which describes the motion of an incompressible viscous fluid in  $\Omega$  [17, 41]. Here  $\vec{u} : \Omega \rightarrow \mathbb{R}^2$  is the velocity field and  $p : \Omega \rightarrow \mathbb{R}$  is the pressure.

Let  $V = H_0^1(\Omega)^2$  and  $M = L_0^2(\Omega) = \{q \in L^2(\Omega) : \int_{\Omega} q dx = 0\}$ . Then the variational formulation of (2.3) reads as: Find  $(\vec{u}, p) \in V \times M$  such that

$$\begin{cases} a(\vec{u}, \vec{v}) + b(\vec{v}, p) = (\vec{f}, \vec{v}), & \forall \vec{v} \in V, \\ b(\vec{u}, q) = 0, & \forall q \in M; \end{cases} \quad (2.4)$$

where

$$\begin{aligned} a(\vec{v}, \vec{w}) &= \int_{\Omega} \nabla \vec{v} : \nabla \vec{w} dz, \quad \forall \vec{v} \in V, \\ b(\vec{v}, q) &= \int_{\Omega} \operatorname{div} \vec{v} q dz \quad \forall q \in M; \end{aligned}$$

with  $\nabla \vec{v} : \nabla \vec{w} = \sum_{i=1}^2 \frac{\partial v_i}{\partial x_j} \frac{\partial w_i}{\partial x_j}$ .

Suppose  $\vec{f} \in L^2(\Omega)^2$ , then it is easy to see that  $|(\vec{f}, \vec{v})| \leq \|\vec{f}\|_{0,\Omega} \|\vec{v}\|_{1,\Omega}$  for any  $\vec{v} \in V$ . In addition, one can prove the following coercive condition

$$\|v\|_{1,\Omega} \lesssim a(\vec{v}, \vec{v}) \quad \forall \vec{v} \in V, \quad (2.5)$$

and Babuska-Brezzi condition (or inf-sup condition) [18, 26, 36]

$$\inf_{q \in M} \sup_{\vec{v} \in V} \frac{(\operatorname{div} \vec{v}, q)}{\|\vec{v}\|_{1,\Omega} \|q\|_{0,\Omega}} \geq C. \quad (2.6)$$

Thus, the variational problem (2.4) exists an unique solution  $(\vec{u}, p)$ .

## 2.3 Finite element spaces

The finite element methods solve variational problems associated with boundary value problem on some finite dimensional spaces, which are called the finite element spaces. To construct a finite element space, first, a triangulation  $\mathcal{T}_h$  is established on  $\bar{\Omega}$ ; second, we define a finite element on element  $T$  for each  $T \in \mathcal{T}_h$ . According to [35], a finite element in  $\mathbb{R}^2$  is a triple  $(T, P, \Sigma)$  where

1.  $T$  is a close subset of  $\mathbb{R}^2$  with non empty interior and Lipschitz-continuous boundary;
2.  $P$  is a space of real-valued functions defined over the set  $T$ ;
3.  $\Sigma$  is a finite set of linearly independent linear forms which is  $P$ -unisolvent defined over the space  $P$ . Those linear forms are called degrees of freedom.

There are many different finite elements. In this work, we concentrate on  $C^0$  Lagrange elements and the Crouzeix-Raviart element.

### 2.3.1 Finite Element meshes

A finite element mesh (or triangulation )  $\mathcal{T}_h$  is a partition of  $\Omega$  into finitely many subdomains, called elements. Elements can be triangles or quadrilaterals. In this work, we restrict our



discussion on conforming meshes consisting of simplexes, i.e. triangles. Here, a conforming mesh means that the intersection of any two elements is empty, a common vertex, or a common edge.

For any  $T \in \mathcal{T}_h$ , let  $h_T$  denote the diameter of  $T$  and  $\rho_T$  denote the supremum of the diameter of the spheres inscribed in  $T$ . The mesh size of  $\mathcal{T}_h$  is denoted by  $h = \max\{h_T : T \in \mathcal{T}_h\}$ . We say  $\mathcal{T}_h$  is regular if there exists a constant  $\sigma$  such that

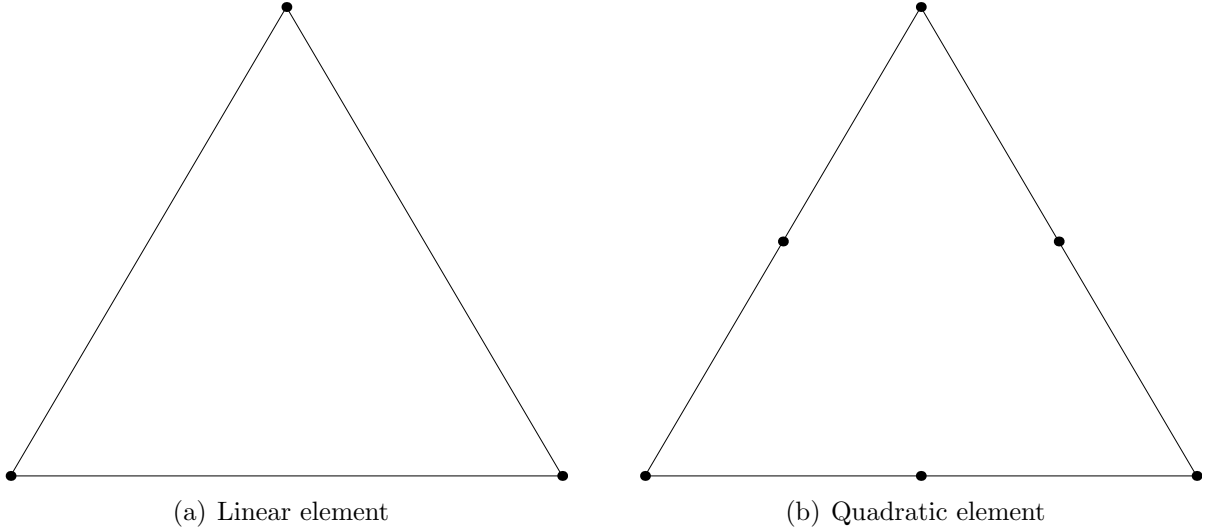
$$\frac{h_T}{\rho_T} \leq \sigma, \quad \forall T \in \mathcal{T}_h.$$

In the sequel, we always assume mesh  $\mathcal{T}_h$  is regular.

A mesh  $\mathcal{T}_h$  is called quasi-uniform mesh if there exists a constant  $\nu \geq 0$  such that

$$\frac{h}{h_T} \leq \nu, \quad \forall T \in \mathcal{T}_h.$$

A very important type of quasi-uniform mesh is uniform mesh. Uniform mesh plays an important role in superconvergence analysis [32, 71, 109]. Meshes can also be categorized as structured or unstructured. Structured meshes have a uniform topological structure that unstructured meshes lack. In this work, we shall consider both structured meshes and unstructured meshes.

Figure 1:  $C_0$  Lagrange element

### 2.3.2 $C^0$ Lagrange Elements

For any  $T \in \mathcal{T}$ , let  $a_j, 1 \leq j \leq 3$ , be the vertices of  $T$ . For any  $r > 0$ , let

$$\Sigma_r(T) = \left\{ x = \sum_{j=1}^3 \lambda_j a_j; \sum_{j=1}^3 \lambda_j = 1, \lambda_j \in \left\{ 0, \frac{1}{r}, \dots, \frac{r-1}{r}, 1 \right\}, 1 \leq j \leq 3 \right\}.$$

Then the  $C^0$  Lagrange element of degree  $r$  is defined as  $(T, \mathbb{P}_r(T), \Sigma_r(T))$ . Typical examples of  $C^0$  Lagrange elements include linear element and quadratic element, see Fig 1. For linear element, degrees of freedom only contains vertices. However, for quadratic element, degrees of freedom includes both vertices and edge centers.

The  $C^0$  finite element space of order  $r$  associated with mesh  $\mathcal{T}_h$  is defined as

$$S^{h,r} = \{v \in C(\bar{\Omega}) : v|_T \in \mathbb{P}_r(T), \forall T \in \mathcal{T}_h\}.$$

Note that the choice of  $\Sigma_r(T)$  guarantees the continuity of  $v$  across the boundaries of elements

in  $\mathcal{T}_h$ . Let  $\mathcal{N}_h$  denote the the set of all mesh nodes. Then for any function  $v \in S^{h,r}$ , it can be written as

$$v = \sum_{z \in \mathcal{N}_h} v(z) \phi_z$$

where  $\phi_z$  is called the nodes shape function associated with  $z$  is defined by

$$\phi_z(z') = \begin{cases} 1 & \text{if } z' = z; \\ 0 & \text{if } z' \in \mathcal{N}_h \setminus \{z\}. \end{cases}$$

It is easy to see that the set  $\{\phi_z : z \in \mathcal{N}_h\}$  forms a basis of  $S^{h,r}$ , which is called nodal basis.

Let  $I_h : C(\bar{\Omega}) \rightarrow S^{h,r}$  denote the standard Lagrange interpolation operator, i.e.

$$I_h u = \sum_{z \in \mathcal{N}_h} u(z) \phi_z, \tag{2.7}$$

for any  $u \in C(\bar{\Omega})$ .

For piecewise defined function, [25, 33, 35] prove the following smooth result:

**Theorem 2.1.** *Let  $\mathcal{T}_h$  be a mesh of  $\Omega$ . Let  $k \geq 1$ . Then a piecewise infinitely differentiable function  $v : \bar{\Omega} \rightarrow \mathbb{R}$  over the mesh  $\mathcal{T}_h$  belongs to  $H^k(\Omega)$  if and only if  $v \in C^{k-1}(\Omega)$ .*

Theorem 2.1 implies  $S^{h,r} \subset H^1$ . In addition, let  $S_0^{h,r} = S^{h,r} \cap H_0^1(\Omega)$ . It means that the finite element space  $S_0^{h,r}$  is a subspace of  $H_0^1$  where the model problem (2.1) is posed. For this reason, the  $C^0$  Lagrange elements are often referred to as conforming elements.

### 2.3.3 The Crouzeix-Raviart element

For the conforming finite elements, it is assumed that the finite element spaces lie in the function space in which the variational problem is posed. However, there are too many limitations of conforming finite elements. For example, for fourth order elliptic differential equations, conforming finite elements require  $C^1$  elements, and this leads to extremely large systems of equations. Thus, people try to relax the condition by not requiring finite element spaces to be subspaces of corresponding function spaces. Those types of finite elements are called nonconforming elements.

In this work, we consider the simplest nonconforming element, i.e. the Crouzeix-Raviart element. It is also referred to as the nonconforming  $P_1$  element. Differentiating from the conforming linear element, the degrees of freedom are on edge centers; see Fig 2.

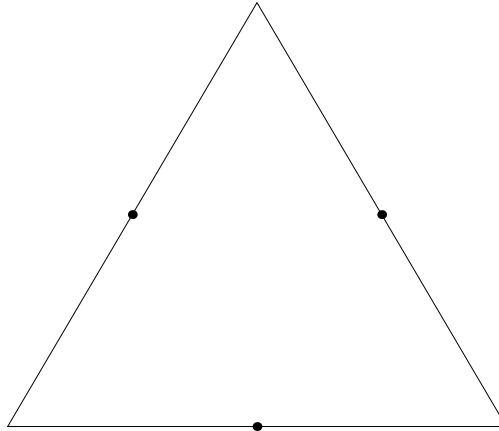


Figure 2: The Crouzeix Raviart element

Let  $\mathcal{T}_h$  be a shape regular triangulation of  $\Omega$ . Denote the set of all edges and edge centers by  $\mathcal{E}_h$  and  $\mathcal{M}_h$ , respectively. For any edge  $e \in \mathcal{E}$ , let  $\mathcal{M}(e)$  be the middle point of edge  $e$ .

Define the Crouzeix-Raviart finite element space as follows:

$$X^{\text{nc}} := \{v \in L_2(\Omega) : v|_T \in \mathbb{P}_1(T) \text{ and } v \text{ is continuous at } \mathcal{M}_h\},$$

$$X_0^{\text{nc}} := \{v \in X^{\text{nc}} : v(\mathcal{M}(e)) = 0 \text{ for all } e \in \mathcal{E}_h \cap \partial\Omega\}.$$

## 2.4 Finite element approximation of model problems

In the subsection, we firstly introduce the finite element approximation for second order elliptic problem (2.2) using both  $C^0$  Lagrange elements and the Crouzeix-Raviart element. Then we discretize the Stokes equation with stable nonconforming finite element.

### 2.4.1 Conforming finite element approximation of elliptic equation

The conforming finite element approximation of the model problem (2.2) is to find  $u_h \in S_0^{h,r}$  such that

$$\mathcal{B}(u_h, v_h) = L(v_h), \quad \forall v_h \in S_0^{h,r}. \quad (2.8)$$

Note that  $S_0^{h,r}$  is a subspace of  $H_0^1(\Omega)$ . Then equations (2.2) and (2.8) implies the following

Theorem:

**Theorem 2.2. ( C ea's Lemma )** *There exist a constant  $C > 0$ , independent of the  $S_0^{h,r}$ , such that*

$$\|u - u_h\|_{1,\Omega} \leq C \inf_{v \in S_0^{h,r}} \|u - v\|_{1,\Omega}.$$

Taking  $v$  as the Lagrange interpolation of  $u$ , see (2.7), we can get the following  $H^1$  error estimate:

**Theorem 2.3.** *If the solution  $u$  of (2.2) is in the space  $H^{r+1}(\Omega) \cap H_0^1(\Omega)$  and  $u_h \in S_0^{h,r}$  is*

the solution of (2.8), then

$$\|u - u_h\|_{1,\Omega} \lesssim h^r |u|_{r+1,\Omega}.$$

Using the duality argument (or Nitsche' Trick) [17, 25, 35], we can prove the following  $L^2$  error estimate:

**Theorem 2.4.** *Under the same hypothesis of Theorem 2.3, we have the following error estimate:*

$$\|u - u_h\|_{0,\Omega} \lesssim h^{r+1} |u|_{r+1,\Omega}.$$

**Remark.** Theorems 2.3 and 2.4 imply that the optimal convergence rate of  $L^2$  Error is two, the optimal convergence rate of  $H^1$  is one and hence the optimal convergence rate of piecewise  $H^2$  error is zero when we using linear element. It means an approximation of second order derivatives using piecewise linear element should not converge. In chapter 4, we will propose an effective post-processing method which can produce superconvergent or even ultraconvergent Hessian matrix.

#### 2.4.2 Nonconforming finite element approximation of elliptic equation

The nonconforming finite element approximation of (2.2) consists of finding  $u_h \in X_0^{\text{nc}}$  such that

$$\mathcal{B}_h(u_h, v_h) = (f, v_h) \quad \forall v_h \in X_0^{\text{nc}}, \quad (2.9)$$

where

$$\mathcal{B}_h(w, v) := \sum_{T \in \mathcal{T}_h} \int_T (\mathcal{D} \nabla w \cdot \nabla v + cuv) dx,$$

for all  $w, v \in X^{\text{nc}}$ . For a subdomain  $\mathcal{A}$  of  $\Omega$ , define the broken semi-norm on  $\mathcal{A}$  as  $\|v\|_{1,h,\mathcal{A}}^2 := \sum_{T \in \mathcal{T}_h \cap \mathcal{A}} |v|_{1,T}^2$ . Using the second Strang Lemma [17, 25, 35], we can prove the following error estimate

**Theorem 2.5.** *If the solution  $u$  of (2.2) is in the space  $H^2(\Omega) \cap H_0^1(\Omega)$  and  $u_h \in X_0^{\text{nc}}$  is the solution of (2.9), then*

$$\|u - u_h\|_{1,h,\Omega} \lesssim h|u|_{2,\Omega}.$$

**Remark.** Theorem 2.5 implies that the optimal convergence rate in discrete  $H^1$  norm is 1. In the chapter 3, we will propose a gradient recovery operator for the Crouzeix-Raviart element which is numerically verified to be superconvergent to the exact gradient.

### 2.4.3 Nonconforming finite element approximation of Stokes Equation

In this subsection, we restrict our discussion on nonconforming finite element approximation of Stokes equation. Define  $\widetilde{M}^h = \{q \in L^2(\Omega) : q|_T \in \mathbb{P}_0, T \in \mathcal{T}\}$  and  $M^h = \widetilde{M}^h \cap L_0^2(\Omega)$ . Let  $V^h = X^{\text{nc}} \times X^{\text{nc}}$  and  $V_0^h = X_0^{\text{nc}} \times X_0^{\text{nc}}$ . The nonconforming finite element approximation reads as finding  $(\vec{u}_h, p_h) \in V_0^h \times M^h$  such that

$$\begin{cases} a_h(\vec{u}_h, \vec{v}_h) + b_h(\vec{v}_h, p_h) = (\vec{f}, \vec{v}_h), & \forall \vec{v}_h \in V_0^h, \\ b_h(\vec{u}_h, q_h) = 0, & \forall q_h \in M^h; \end{cases} \quad (2.10)$$

where

$$\begin{aligned} a_h(\vec{v}_h, \vec{w}_h) &= \int_{\Omega} \nabla_h \vec{v}_h : \nabla_h \vec{w}_h dx = \sum_{T \in \mathcal{T}_h} \int_T \nabla \vec{v}_h : \nabla \vec{w}_h dx, \\ b(\vec{v}_h, q_h) &= \int_{\Omega} \nabla_h \cdot \vec{v}_h q_h dx = \sum_{T \in \mathcal{T}_h} \int_T \nabla \cdot \vec{v}_h q_h dx; \end{aligned}$$

for any  $\vec{v}_h, \vec{w}_h \in V^h$  and  $q_h \in M^h$ . Here  $\nabla_h$  is called broken gradient operator. [16] prove the following discrete inf-sup condition

$$\inf_{q_h \in M^h} \sup_{\vec{v}_h \in V_0^h} \frac{b_h(\vec{v}_h, q_h)}{\|\vec{v}_h\|_{1,h,\Omega} \|q_h\|_{0,\Omega}} \geq \beta \geq 0, \quad (2.11)$$

where the constant  $\beta$  is independent of  $h$  and  $\|\vec{v}_h\|_{1,h,\mathcal{A}}^2 = \|v_1\|_{1,h,\mathcal{A}}^2 + \|v_2\|_{1,h,\mathcal{A}}^2$  for any  $\mathcal{A} \subset \Omega$ .

Therefore the discrete variational problem (2.10) is well posed and the following discrete  $H^1$  error estimate holds:

**Theorem 2.6.** *Let the solution  $(\vec{u}, p)$  of the Stokes problem (2.4) satisfy*

$$\vec{u} \in (H^2(\Omega) \cap H_0^1(\Omega))^2, \quad p \in H^1(\Omega) \cap L_0^2(\Omega).$$

*Then*

$$\|\vec{u} - \vec{u}_h\|_{1,h,\Omega} + \|p - p_h\|_{0,\Omega} \leq h(|u|_{2,\Omega} + |p|_{1,\Omega}).$$

**Remark.** Theorem 2.6 implies that  $H^1$  error of velocity field is of order  $O(h)$ . Our numerical experiments in Chapter 3 indicates that the gradient recovery method proposed in Chapter 3 can also apply to each component of velocity to get superconvergent gradient.



## CHAPTER 3 Gradient recovery for the Crouzeix-Raviart element

The Crouzeix-Raviart element was first proposed in [37] to solve stationary Stokes problem. It has many useful properties [21], such as commutative relations with respect to the gradient, divergence, and curl operators, and the existence of an interpolation operator that can be defined on  $H^s(\Omega)$  for  $s \geq \frac{1}{2}$ . The applications of the Crouzeix-Raviart element can be found in solids [24, 43], fluids [37] and electromagnets [22, 23], which can be called as a universal element [21].

In this chapter, we concentrate on post-processing of the Crouzeix-Raviart element and its application. Specifically speaking, we propose a good gradient recovery method for the Crouzeix-Raviart element, which is based on the standard of [4], and apply it to a posteriori error estimates for adaptive nonconforming finite element method.

We provide the main definition of the gradient recovery operator in Section 3.1. Some illustrations of the proposed gradient recovery operator are given in Section 3.2. The properties are investigated in Section 3.3. Two numerical examples are presented to illustrate superconvergence of the proposed gradient recovery method in Section 3.4 and its application to adaptive finite element method will be discussed in Section 3.5. Finally, some conclusions will be drawn in Section 3.6.

### 3.1 Definition of Gradient Recovery Operator

Gradient recovery is a method providing a better approximation of  $\nabla u$ . We introduce a new gradient recovery operator  $G_h : X^{\text{nc}} \rightarrow X^{\text{nc}} \times X^{\text{nc}}$ . The structure of  $G_h u_h$  relies on the fact

that every function in  $X^{\text{nc}}$  is uniquely defined by its values at edge centers. Given a finite element solution  $u_h \in X^{\text{nc}}$ , we only need to define  $G_h u_h$  at edge centers. After determining values of  $G_h u_h$  at all edge centers, we obtain  $G_h u_h \in X^{\text{nc}} \times X^{\text{nc}}$  on the whole domain by interpolation, i.e.

$$G_h u_h := \sum_{z \in \mathcal{M}} (G_h u_h)(z) \phi_z,$$

where  $\{\phi_z : z \in \mathcal{M}_h\}$  is nodal basis of  $X^{\text{nc}}$ .

For any edge center  $z = (x, y)$  and  $1 \leq n \in \mathbb{N}$ , define the union of elements around  $z$  in the first  $n$  layers as follows

$$\mathcal{L}(z, n) := \cup\{T : T \in \mathcal{T}_h, T \cap \mathcal{L}(z, n-1) = e \text{ for some } e \in \mathcal{E}_h\},$$

with  $\mathcal{L}(z, 0) := \{e \in \mathcal{E}_h : z \text{ is the middle point of } e\}$ .

Let  $z$  be the middle point of some interior edge  $e$ , i.e.  $z = \mathcal{M}(e)$ . Intuitively, it is natural to use information on triangle  $T_1$  and  $T_2$  to recover gradient at  $z$  where  $T_1 \cap T_2 = e$ , i.e.  $T_1, T_2 \in \mathcal{L}(z, 1)$ . Let  $|T_i|$  be the area of  $T_i$  ( $i = 1, 2$ ) and  $|\omega|$  be the sum of  $|T_1|$  and  $|T_2|$ . There are several possible ways to define  $(G_h u_h)(z)$ .

1. Simple averaging:

$$(G_h u_h)(z) = \frac{1}{2}(\nabla u_h|_{T_1} + \nabla u_h|_{T_2}).$$

2. Weighted averaging:

$$(G_h u_h)(z) = \left(\frac{|T_1|}{|\omega|} \nabla u_h|_{T_1} + \frac{|T_2|}{|\omega|} \nabla u_h|_{T_2}\right).$$

3. Fitting a linear polynomial:

$$(G_h u_h)(z) = \nabla p_z(z)$$

where

$$p_z = \arg \min_{p \in \mathbb{P}_1(L(z,1))} \sum_{\tilde{z} \in \mathcal{M} \cap L(z,1)} |(u_h - p)(\tilde{z})|^2;$$

and  $\mathcal{M}_h$  is the set of edge centers.

**Remark.** Method 3 can be viewed as nonconforming counterpart of SCR proposed in [58].

On regular pattern uniform mesh, the above three methods give the same result.

**Remark.** Hu and Ma prove  $O(h^{1.5})$  superconvergence for the simple averaging method on uniform mesh of regular pattern in [54] by using the equivalence between the Crouzeix-Ravart element and Raviart-Thomas element [76] and the superconvergence result of Raviart-Thomas element [20].

However, just as we know for  $C^0$  Lagrange elements, the approximation property of the above three methods depends heavily on symmetry of local element patch  $L(z, 1)$  with respect to  $z$ . Simple calculation combined with Taylor expansion reveals that they exhibit superconvergence only on regular pattern uniform meshes.

Inspired by the idea of Polynomial Preserving Recovery (PPR) [79, 80, 107], the new gradient recovery method fits a quadratic polynomial at every edge center. Let  $z_i = (x_i, y_i) \in \mathcal{M}_h$  be an edge center and  $\mathcal{K}_{z_i}$  denote a patch of elements around  $z_i$ . Let  $p_{z_i} \in \mathbb{P}_2(\mathcal{K}_{z_i})$  be the quadratic polynomial that best fits  $u_h$  at the edge centers in  $\mathcal{K}_{z_i}$  in discrete least-squares sense, i.e.,

$$p_{z_i} = \arg \min_{p \in \mathbb{P}_2(K_{z_i})} \sum_{\tilde{z} \in \mathcal{M}_h \cap K_{z_i}} |(u_h - p)(\tilde{z})|^2,$$

and define the recovered gradient at  $z_i$  as

$$(G_h u_h)(z_i) = \nabla p_{z_i}(x_i, y_i).$$

Let  $z_{i_0}, z_{i_1}, \dots, z_{i_m}$  denote all the edge centers in  $\mathcal{K}_{z_i}$ . Without loss of generality, let  $z_i = z_{i_0}$  and  $h_i = \max\{|z_{i_j} - z_i| : 1 \leq j \leq m\}$ . To avoid the computational instability resulting from small  $h_i$ , we introduce coordinate transformation

$$F : (x, y) \rightarrow (\xi, \eta) = \frac{(x, y) - (x_i, y_i)}{h_i},$$

All computations are carried out on the local element patch  $\widehat{K}_{z_i} = F(K_{z_i})$ . Then we can rewrite the fitting polynomial as

$$p_{z_i}(x, y) = P^T a = \widehat{P}^T \widehat{a}.$$

with

$$\begin{aligned} P^T &= (1, x, y, x^2, xy, y^2), & \widehat{P}^T &= (1, \xi, \eta, \xi^2, \xi\eta, \eta^2); \\ a^T &= (a_0, a_1, a_2, a_3, a_4, a_5), & \widehat{a}^T &= (a_0, h_i a_1, h_i a_2, h_i^2 a_3, h_i^2 a_4, h_i^2 a_5). \end{aligned}$$

The coefficient vector  $\widehat{a}$  can be obtained by solving

$$A^T A \widehat{a} = A^T b.$$

with  $b^T = (u_h(z_{i_0}), u_h(z_{i_1}), \dots, u_h(z_{i_m}))$  and

$$A = \begin{pmatrix} 1 & \xi_0 & \eta_0 & \xi_0^2 & \xi_0\eta_0 & \eta_0^2 \\ 1 & \xi_1 & \eta_1 & \xi_1^2 & \xi_1\eta_1 & \eta_1^2 \\ 1 & \xi_2 & \eta_2 & \xi_2^2 & \xi_2\eta_2 & \eta_2^2 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 1 & \xi_m & \eta_m & \xi_m^2 & \xi_m\eta_m & \eta_m^2 \end{pmatrix}.$$

Then the recovered gradient at  $z_i$  is defined as

$$(G_h u_h)(z_i) = \nabla p_{z_i}(0,0) = \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \frac{1}{h_i} \begin{pmatrix} \hat{a}_1 \\ \hat{a}_2 \end{pmatrix} \quad (3.1)$$

Since  $p_{z_i}$  has 6 unknowns, it requires that  $\mathcal{K}_{z_i}$  contains at least 6 edge centers. This naturally leads to  $\mathcal{K}_{z_i} = \mathcal{L}(z_i, 2)$ . Also, notice that  $\mathcal{K}_{z_i}$  containing more than 6 middle points can not guarantee the uniqueness of  $p_{z_i}$ . If  $p_{z_i}$  is not unique and  $\mathcal{L}(z_i, n) = \mathcal{K}_{z_i} \subsetneq \mathcal{L}(z_i, n+1)$ , then set  $\mathcal{K}$  equal  $\mathcal{L}(z_i, n+1)$  and recompute  $p_{z_i}$ . The patch  $\mathcal{K}_{z_i}$  is defined to be the first  $\mathcal{L}(z_i, n)$  such that  $p_{z_i}$  is unique.

**Remark.** In superconvergent patch recovery (SPR) [111, 112] and PPR, if  $z_i$  lies on an edge between two vertices  $z_{i_1}$  and  $z_{i_2}$ , then recovered gradient is defined at  $z_i$  as  $(G_h u_h)(z_i) = \beta \nabla p_{z_{i_1}}(z_i) + (1 - \beta) \nabla p_{z_{i_2}}(z_i)$  where  $\beta$  is determined by the ratio of distances of  $z_i$  to  $z_{i_1}$  and  $z_{i_2}$ .

**Remark.** For the most usual cases,  $\mathcal{K}_{z_i} = \mathcal{L}_{z_i,2}$  can guarantee the uniqueness of  $p_{z_i}$ . However, we can use a slightly large patch to get an litter improved results. For the numerical results in the chapter, we use  $\mathcal{K}_{z_i} = \mathcal{L}_{z_i,2}$  for elliptic equation and  $\mathcal{K}_{z_i} = \mathcal{L}_{z_i,3}$  for stokes equation.

**Remark.** From the above definition, we can see that the proposed gradient recovery method is problem independent. In the following sections, we will apply the gradient recovery to second order elliptic problem (2.1) and Stokes equation (2.3) to show that it can produce superconvergent post-processing gradients.

## 3.2 Illustrations of the Proposed Gradient Recovery Method

To illustrate the essential idea of the above gradient recovery operator, we study regular pattern and chevron pattern uniform meshes in detail. Note that  $G_h$  is a linear operator from  $X^{\text{nc}}$  to  $X^{\text{nc}} \times X^{\text{nc}}$ . To simplify notation, sometimes  $G_h$  is rewritten as the following component form

$$G_h = \begin{pmatrix} G_h^x \\ G_h^y \end{pmatrix}, \quad (3.2)$$

where  $G_h^x$  and  $G_h^y$  are linear operators from  $X^{\text{nc}}$  to  $X^{\text{nc}}$ .

### 3.2.1 Regular Pattern

All edge centers can be classified into three cases: vertical edge center, horizontal edge center and diagonal edge center. We select two typical cases as showed in Figure 3 and 4. Other cases can be treated as an orthogonal rotation of the above two cases.

First, consider the case that  $z$  is a vertical edge center, i.e. the dotted point in Figure 3. In order to investigate the approximation property of recovery operator, we replace the

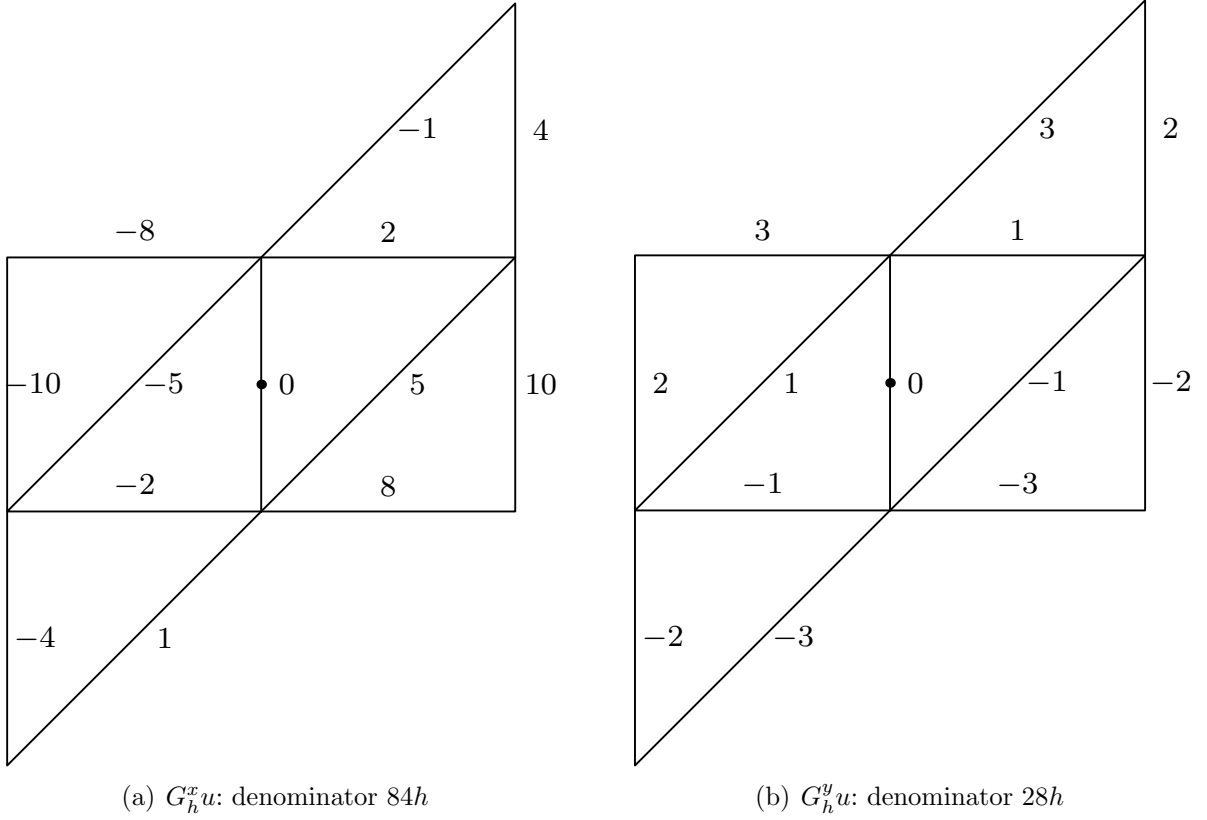


Figure 3: Recovery at vertical edge center of Regular Pattern

finite element solution  $u_h$  by the exact solution  $u$  in the sequel. Given

$$\vec{\xi} = (0, 1, 2, 1, -1, -2, -1, -1, 1, 2, 1, -2, -1)^T,$$

$$\vec{\eta} = (0, 0, 0, 1, 1, 0, 0, -1, -1, 2, 2, -2, -2)^T,$$

$$b = (u_0, u_1, u_2, u_3, u_4, u_5, u_6, u_7, u_8, u_9, u_{10}, u_{11}, u_{12})^T,$$

where  $u_i = u(F^{-1}(\xi_i, \eta_i))$  for  $0 \leq i \leq 12$ , we want to fit a polynomial

$$\hat{p}_2(\xi, \eta) = (1, \xi, \eta, \xi^2, \xi\eta, \eta^2)(\hat{a}_0, \hat{a}_1, \hat{a}_2, \hat{a}_3, \hat{a}_4, \hat{a}_5)^T,$$

in the least square sense with respect to  $(\xi, \eta)$ . Let  $\vec{e} = (1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1)^T$  and

$$A = (\vec{e}, \vec{\xi}, \vec{\eta}, \vec{\xi} \circ \vec{\xi}, \vec{\xi} \circ \vec{\eta}, \vec{\eta} \circ \vec{\eta})$$

where  $\circ$  is Hadamard product for matrices. Simple calculation shows that

$$B = \begin{pmatrix} \frac{17}{53} & 0 & 0 & -\frac{5}{53} & \frac{6}{53} & -\frac{6}{53} \\ \frac{12}{53} & \frac{5}{84} & -\frac{1}{28} & -\frac{233}{5512} & \frac{375}{5512} & -\frac{37}{424} \\ -\frac{3}{53} & \frac{5}{42} & -\frac{1}{14} & \frac{157}{1378} & -\frac{93}{1378} & -\frac{1}{106} \\ \frac{12}{53} & \frac{1}{42} & \frac{1}{28} & -\frac{339}{5512} & \frac{587}{5512} & -\frac{37}{424} \\ 0 & -\frac{2}{21} & \frac{3}{28} & \frac{3}{104} & -\frac{19}{104} & \frac{1}{8} \\ -\frac{3}{53} & -\frac{5}{42} & \frac{1}{14} & \frac{157}{1378} & -\frac{93}{1378} & -\frac{1}{106} \\ \frac{12}{53} & -\frac{5}{84} & \frac{1}{28} & -\frac{233}{5512} & \frac{375}{5512} & -\frac{37}{424} \\ \frac{12}{53} & -\frac{1}{42} & -\frac{1}{28} & -\frac{339}{5512} & \frac{587}{5512} & -\frac{37}{424} \\ 0 & \frac{2}{21} & -\frac{3}{28} & \frac{3}{104} & -\frac{19}{104} & \frac{1}{8} \\ -\frac{3}{53} & \frac{1}{21} & \frac{1}{14} & \frac{51}{1378} & \frac{119}{1378} & -\frac{1}{106} \\ 0 & -\frac{1}{84} & \frac{3}{28} & -\frac{3}{104} & -\frac{7}{104} & \frac{1}{8} \\ -\frac{3}{53} & -\frac{1}{21} & -\frac{1}{14} & \frac{51}{1378} & \frac{119}{1378} & -\frac{1}{106} \\ 0 & \frac{1}{84} & -\frac{3}{28} & -\frac{3}{104} & -\frac{7}{104} & \frac{1}{8} \end{pmatrix}$$

where  $B^T = (A^T A)^{-1} A^T$ . Thus we have  $\hat{a} = B^T b$ . Recall that

$$(\hat{a}_0, \hat{a}_1, \hat{a}_2, \hat{a}_3, \hat{a}_4, \hat{a}_5) = (a_0, ha_1, ha_2, h^2 a_3, h^2 a_4, h^2 a_5),$$



and hence it holds that

$$p_2(x, y) = \hat{a}_0 + \frac{1}{h}(\hat{a}_1x + \hat{a}_2y) + \frac{1}{h^2}(\hat{a}_3x^2 + \hat{a}_4xy + \hat{a}_5y^2).$$

Differentiating with respect to  $x$  and  $y$  and evaluate at  $(0, 0)$ , it follows that

$$(G_h^x u)(z) = \frac{1}{84h}(5u_1 + 10u_2 + 2u_3 - 8u_4 - 10u_5 - 5u_6 - 2u_7 + 8u_8 + 4u_9 - u_{10} - 4u_{11} + u_{12}), \quad (3.3)$$

and

$$(G_h^y u)(z) = \frac{1}{28h}(-u_1 - 2u_2 + u_3 + 3u_4 + 2u_5 + u_6 - u_7 - 3u_8 + 2u_9 + 3u_{10} - 2u_{11} - 3u_{12}). \quad (3.4)$$

By using computer algebra system such as **Mathematica**, it is very easy to get the following

Taylor expansion for  $(G_h u)(z)$

$$(G_h u)(z) = \begin{pmatrix} u_x(z) + \frac{h^2}{42}(21u_{xxx}(z) + 12u_{xxy} + 19u_{xyy} + 3u_{yyy}(z)) + o(h^2) \\ u_y(z) + \frac{h^2}{42}(39u_{xxy} + 39u_{xyy}(z) + 22u_{yyy}(z)) + o(h^2) \end{pmatrix}.$$

It obviously demonstrates that  $(G_h u)(z)$  provides a second order approximation to  $\nabla u$ .

Then we consider that  $z$  is a diagonal edge center as seen in Figure 4. Repeating the same procedure as above, we obtain

$$(G_h^x u)(z) = \frac{1}{84h}(5u_1 + 10u_2 + 8u_3 - 2u_4 - 10u_5 - 5u_6 - 8u_7 + 2u_8 + u_9 - 4u_{10} - u_{11} + 4u_{12}), \quad (3.5)$$

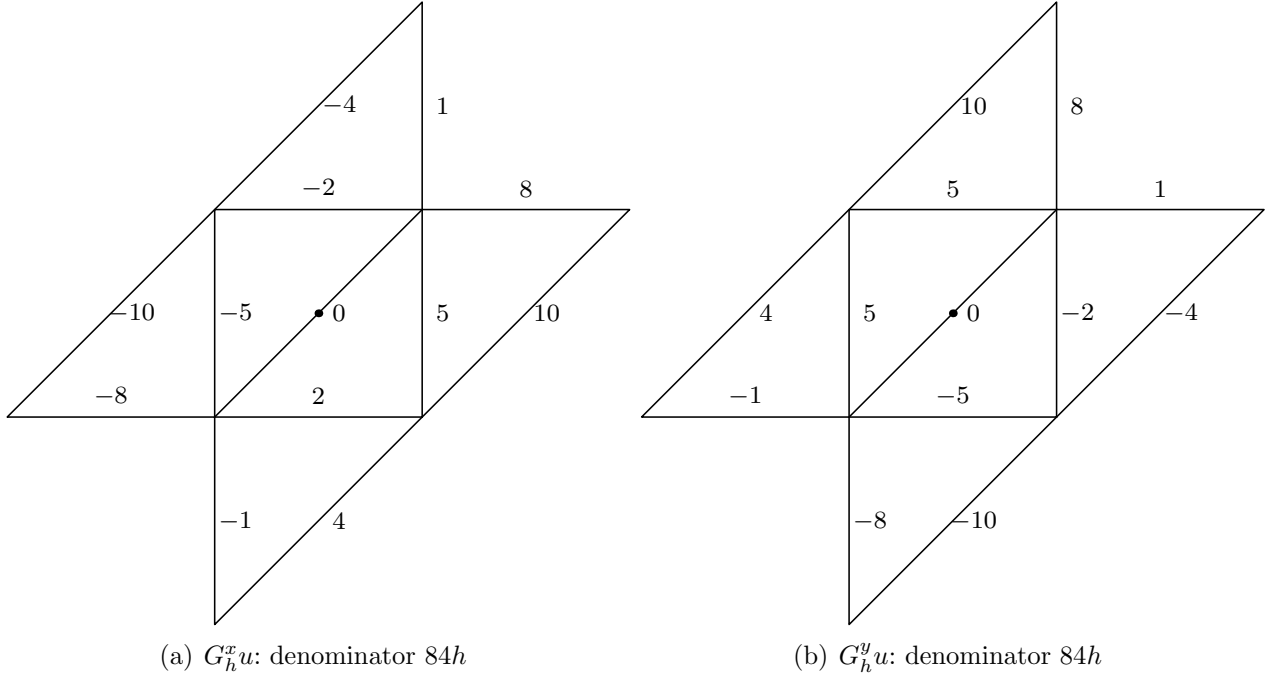


Figure 4: Recovery at diagonal edge center of Regular Pattern

and

$$(G_h^y u)(z) = \frac{1}{84h} (-2u_1 - 4u_2 + u_3 + 5u_4 + 4u_5 + 5u_6 - u_7 - 5u_8 + 8u_9 + 10u_{10} - 8u_{11} - 10u_{12}). \quad (3.6)$$

Also, we have the following Taylor expansion

$$(G_h u)(z) = \begin{pmatrix} u_x(z) + \frac{h^2}{42}(25u_{xxx}(z) + 17u_{xxy} + 10u_{xyy} - 3u_{yyy}(z)) + o(h^2) \\ u_y(z) - \frac{h^2}{42}(3u_{xxx} - 10u_{xxy}(z) - 17u_{xxy} - 25u_{yyy}(z)) + o(h^2) \end{pmatrix}.$$

Again (3.5) and (3.6) provide a second order approximation to  $\nabla u$ .

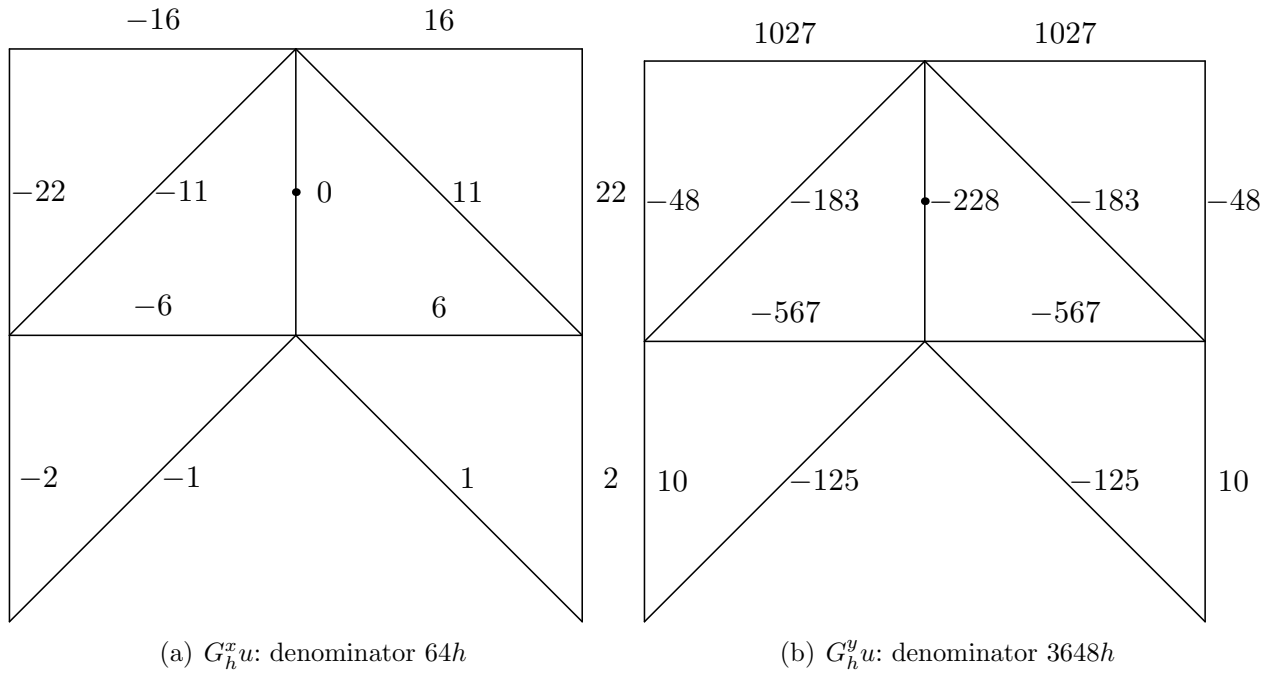


Figure 5: Recovery at vertical edge center of Chevron Pattern

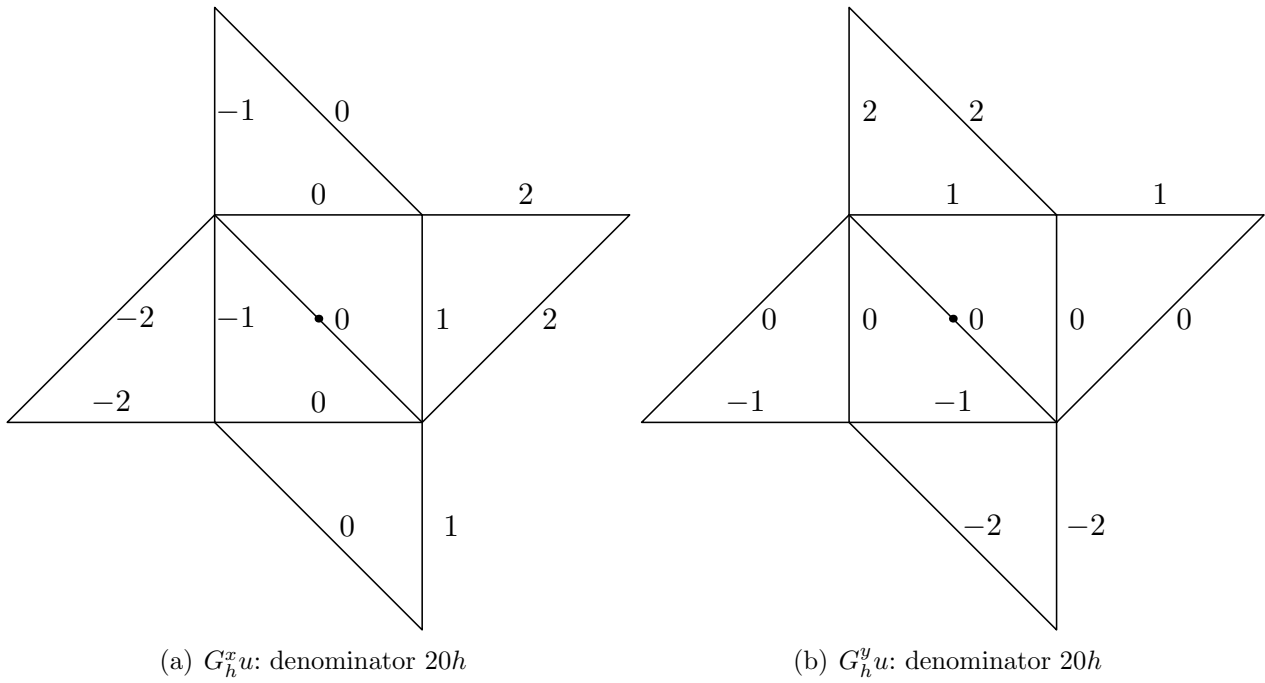


Figure 6: Recovery at diagonal edge center of Chevron Pattern

### 3.2.2 Chevron Pattern

Similar to Regular pattern, the gradient recovery operator are different when  $z$  are horizontal edge center, vertical edge center and diagonal edge center. Again, we choose two typical edge centers as depicted in Figure 5 and Figure 6. First, consider the case when  $z$  is a vertical edge center. Following the same procedure as regular pattern, we get

$$(G_h^x u)(z) = \frac{1}{164h} (11u_1 - 11u_2 - 6u_3 + 6u_4 + 22u_5 + 16u_6 - 16u_7 - 22u_8 - 2u_9 - u_{10} + u_{11} + 2u_{12}), \quad (3.7)$$

and

$$(G_h^y u)(z) = \frac{1}{3648h} (-228u_0 - 183u_1 - 183u_2 - 567u_3 - 567u_4 - 48u_5 + 1027u_6 + 1027u_7 - 48u_8 + 10u_9 - 125u_{10} - 125u_{11} + 10u_{12}). \quad (3.8)$$

It is straightforward to verify that

$$(G_h u)(z) = \begin{pmatrix} u_x(z) + \frac{h^2}{246} (113u_{xxx}(z) + 63u_{xyy}(z)) + o(h^2) \\ u_y(z) + \frac{h^2}{1824} (882u_{xxy}(z) + 419u_{yyy}(z)) + o(h^2) \end{pmatrix}.$$

which clearly indicates that  $G_h u(z)$  converges to  $\nabla u$  at rate of  $O(h^2)$ .

Then we consider  $z$  is diagonal edge center in Chevron pattern uniform mesh. The sampling points in patch  $K_z$  is displayed in Figure 6. Simple calculation verifies that

$$(G_h u)(z) = \frac{1}{20h} \begin{pmatrix} u_1 - u_3 + 2u_5 + 2u_6 - u_8 - 2u_9 - 2u_{10} + u_{12} \\ u_2 - u_4 + u_6 + 2u_7 + 2u_8 - u_{10} - 2u_{11} - 2u_{12} \end{pmatrix}. \quad (3.9)$$

Using Taylor expansion, we get

$$(G_h u)(z) = \begin{pmatrix} u_x(z) + \frac{h^2}{30}(17u_{xxx}(z) + 9u_{xxy}(z) + 12u_{xyy}(z) - 3u_{yyy}(z)) + o(h^2) \\ u_y(z) + \frac{h^2}{30}(3u_{xxx}(z) + 12u_{xxy}(z) - 9u_{xyy}(z) + 17u_{yyy}(z)) + o(h^2) \end{pmatrix}.$$

It approximates  $\nabla u$  with second order accuracy.

### 3.3 Properties of the Gradient Recovery Operator

The two examples in previous section show that  $G_h$  provides a finite difference scheme with 2nd order accuracy. Moreover, we can show  $G_h$  has 2nd order accuracy in case of Criss-cross pattern and Unionjack pattern uniform meshes. In general, the following theorem holds.

**Theorem 3.1.** *The recovery operator  $G_h$  preserves polynomials of degree two on  $\mathcal{K}_z$  for an arbitrary mesh.*

*Proof.* Suppose  $u \in \mathbb{P}_2(\mathcal{K}_z)$ . Then clearly the least square fitting of a polynomial of degree two will reproduce  $u$ , i.e.  $p_z = u$  on  $K_z$ . Thus  $G_h u = \nabla u$  on  $K_z$ .  $\square$

A direct application of Theorem 3.1 and Bramble-Hilbert Lemma implies the following superconvergence result.

**Theorem 3.2.** *Let  $u \in W_\infty^3(K_z)$ . Then*

$$\|\nabla u - G_h u\|_{0,\infty,\mathcal{K}_z} \leq Ch^2 |u|_{3,\infty,\mathcal{K}_z}.$$

*Proof.* It is similar to the proof of Theorem 2.2 in [107].  $\square$

From now on, we assume  $v \in X^{\text{nc}}$  and  $z_i$  is the middle point of an arbitrary edge. By

(3.1), we know

$$(G_h v)(z_i) = \begin{pmatrix} (G_h^x v)(z_i) \\ (G_h^y v)(z_i) \end{pmatrix} = \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \frac{1}{h_i} \begin{pmatrix} \widehat{a}_1 \\ \widehat{a}_2 \end{pmatrix} = \frac{1}{h_i} \begin{pmatrix} \sum_{j=0}^m b_j^1 v_{i_j} \\ \sum_{j=0}^m b_j^2 v_{i_j} \end{pmatrix},$$

where  $v_{i_j} = v(z_{i_j})$  and  $b_j^1$  and  $b_j^2$  are some constants that independent of  $h$  and  $h_i$  for  $0 \leq j \leq m$ . In addition, let  $z_{i_j}$  ( $0 \leq j \leq m$ ) be all edge centers in  $\mathcal{K}_{z_i}$ ,  $z_i = z_{i_1}$ , and  $v_i = v(z_i)$ . By setting  $u \equiv v_i$  in Theorem 3.1, we can conclude that

$$G_h v_i = \begin{pmatrix} 0 \\ 0 \end{pmatrix}.$$

Thus, we can rewrite  $(G_h v)(z_i)$  as

$$(G_h v)(z_i) = \begin{pmatrix} (G_h^x v)(z_i) \\ (G_h^y v)(z_i) \end{pmatrix} = \frac{1}{h_i} \begin{pmatrix} \sum_{j=1}^m b_j^1 (v_{i_j} - v_i) \\ \sum_{j=1}^m b_j^2 (v_{i_j} - v_i) \end{pmatrix}.$$

Notice that for any  $z_{i_j}$ , we can find  $z_i = z_{j_0}, \dots, z_{j_{n_j}} = z_{i_j}$  such that  $z_{j_\ell}$  and  $z_{j_{\ell+1}}$  belong to the same triangle  $T \in \mathcal{T}_h$  and  $z_{j_\ell} \in \mathcal{K}_{z_i} \cap \mathcal{M}$ . Then

$$(G_h^x v)(z_i) = \sum_{j=1}^m b_j^1 \sum_{\ell=0}^{n_j-1} \frac{(v_{j_\ell} - v_{j_{\ell+1}})}{h_i}.$$

Let  $e_{j_\ell}$  denote the line segment connecting  $v_{j_\ell}$  and  $v_{j_{\ell+1}}$ , and  $h_{j_\ell}$  denote the length of  $e_{j_\ell}$ .

Then we have

$$(G_h^x v)(z_i) = \sum_{j=1}^m b_j^1 \sum_{\ell=0}^{n_j-1} \frac{h_{j_\ell}}{h_i} \frac{(v_{j_\ell} - v_{j_{\ell+1}})}{h_{j_\ell}}.$$

Since  $v \in X^{nc}$ , it follows that  $\frac{(v_{j_\ell} - v_{j_\ell+1})}{h_{j_\ell}} = \frac{\partial v}{\partial t_{j_\ell}}$ , where  $t_j$  denotes the unit tangent vector of  $e_{k_j}$ . Notice that  $\frac{h_{j_\ell}}{h_i}$  are bounded by a constant independent of  $h$  and  $h_i$ , then we get

$$|(G_h^x v)(z_i)| \lesssim |v|_{1,\infty,\mathcal{K}_{z_i}}.$$

The same argument yields that

$$|(G_h^y v)(z_i)| \lesssim |v|_{1,\infty,\mathcal{K}_{z_i}}.$$

For any triangle  $T \in \mathcal{T}_h$ , let  $z_1, z_2, z_3$  be the three edge centers of  $T$  and  $K = \mathcal{K}_{z_1} \cup \mathcal{K}_{z_2} \cup \mathcal{K}_{z_3}$ . Then the above result can be summarized as the following theorem

**Theorem 3.3.**  $G_h : X^{nc} \rightarrow X^{nc} \times X^{nc}$  is a linear operator, and there exists a constant  $C$  independent of  $h$  such that

$$\|G_h v\|_{0,\infty,T} \leq C |v|_{1,\infty,K} \quad \forall T \in \mathcal{T}_h,$$

for any  $v \in X^{nc}$ .

With aid of an inverse estimate, we can prove the following corollary

**Corollary 3.4.** There is a constant  $C$  independent of  $h$  such that

$$\|G_h v\|_{0,2,T} \leq C |v|_{1,2,K} \quad \forall T \in \mathcal{T}_h$$

for any  $v \in X^{nc}$ .

Let  $G_X : X \rightarrow X \times X$  be any gradient recovery operator from  $X \rightarrow X \times X$  satisfying consistency condition, localization condition, and boundness and linearity condition in [4] where  $X$  is some finite element space. A classical way to prove the superconvergence of gradient recovery operator is to rewrite  $\nabla u - G_X(u_h)$  as

$$\nabla u - G_X(u_h) = \nabla u - G_X(u_I) + G_X(u_I - u_h) \quad (3.10)$$

where  $u_I$  is the interpolation of the exact solution  $u$  in the finite element space  $X$ . According to [4], we can prove  $O(h^{(1+r)})$  superconvergence result provided that there is  $O(h^{(1+r)})$  supercloseness [12, 71, 101] between the gradient of the finite element solution  $u_h$  and the gradient of the interpolation  $u_I$ . However, there is no such type supercloseness for the Crouzeix-Raviart element. Actually, [69] proved that the best convergence rate of  $|u_h - u_I|_{1,h,\Omega}$  is at best of  $O(h)$ .

It is worth to point out that supercloseness between the gradient of the finite element solution  $u_h$  and the gradient of the interpolation  $u_I$  is only a sufficient condition. It is well known that there is no supercloseness for Lagrange linear element when the mesh is uniform mesh of Criss-cross pattern or Unionjack pattern. But PPR and SPR can still produce superconvergent approximate gradient on this two type of uniform meshes; see [107]. The same thing applies to the proposed gradient recovery method.

### 3.4 Numerical Experiment

In this section, we present numerical examples to demonstrate superconvergence of our gradient recovery method. First, we consider the second order elliptic problem (2.2) and its non-



conforming finite element approximation (2.9). Then, we study the superconvergent property of gradient recovery operator  $G_h$  applied to each component of velocity field for nonconforming finite element approximation (2.10) of Stokes equation (2.4).

In order to identify the performance of the gradient recovery operator, we split mesh vertices  $\mathcal{N}$  into interior vertices  $\mathcal{N}_{h,1}$  and near boundary vertices  $\mathcal{N}_{h,2}$ , where  $\mathcal{N}_{h,2} = \{z \in \mathcal{N}_h : \text{dist}(z, \partial\Omega) \leq L\}$  and  $\mathcal{N}_{h,1} = \mathcal{N}_h \setminus \mathcal{N}_{h,2}$ . Let

$$\Omega_{h,1} = \bigcup \{T \in \mathcal{T}_h : T \text{ has all of its vertices in } \mathcal{N}_{h,1}\},$$

and  $\Omega_{h,2} = \Omega \setminus \Omega_{h,1}$ .

For elliptic equation, let  $u$  be solution of (2.2),  $u_h$  be the solution be the solution of (2.9),  $u_I$  is the interpolation of  $u$  in the Crouzeix-Raviart space  $X^{\text{nc}}$ . The notation used is the following:

$E_h = \|u - u_h\|_{1,h,\Omega}$  denotes broken  $H_1$ -semi error of finite element solution  $u_h$ ;

$E_i = \|u_I - u_h\|_{1,h,\Omega}$  denotes broken  $H_1$ -semi error between  $u_h$  and  $u_I$ ;

$E_r = \|\nabla u - G_h u_h\|_{0,\Omega_{h,2}}$  denotes  $L_2$  error of recovered gradient in the interior domain.

To make notation consist, we use the similar notation for Stokes equation without repeating the definition. In the following numerical test, we take  $L = 0.1$ .

### 3.4.1 Elliptic equation

Consider the following Poisson equation

$$-\Delta u = 2\pi^2 \sin \pi x \sin \pi y, \quad \text{in } \Omega = (0, 1)^2,$$

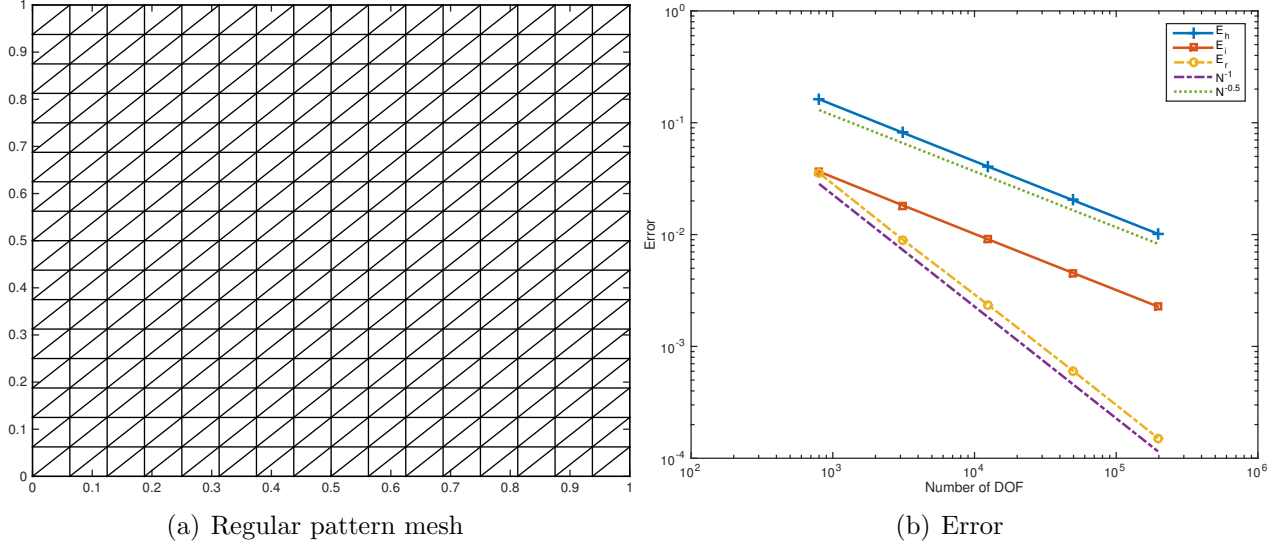


Figure 7: Numerical result of elliptic equation on regular pattern mesh

with  $u = 0$  on  $\partial\Omega$ . The exact solution is

$$u(x, y) = \sin \pi x \sin \pi y.$$

First, uniform meshes are considered. In Figure 7, we report the numerical results for regular pattern uniform meshes. The meshes are obtained by decomposing the unit square into  $16 \times 16$ ,  $32 \times 32$ ,  $64 \times 64$ , and  $128 \times 128$  subsquares and then dividing each subsquare into triangles with regular pattern. We observe  $\|u - u_h\|_{1,h,\Omega}$  is  $O(h)$  as proved in Theorem 2.3. The recovered gradient superconverges at rate of  $O(h^2)$  which doubles the convergence rate of finite element solution. Notice that  $\|u_I - u_h\|_{1,h,\Omega}$  converge at the same order of  $\|u - u_h\|_{1,h,\Omega}$ , which means there is no supercloseness result for the Crouzeix-Raviart element, but it is much smaller than  $\|u - u_h\|_{1,h,\Omega}$ . The numerical results of Chevron, Criss-cross and Unionjack pattern are displayed in Figure 8, 9 and 10, respectively.  $O(h^2)$  superconvergence of recovered gradient is also observed.

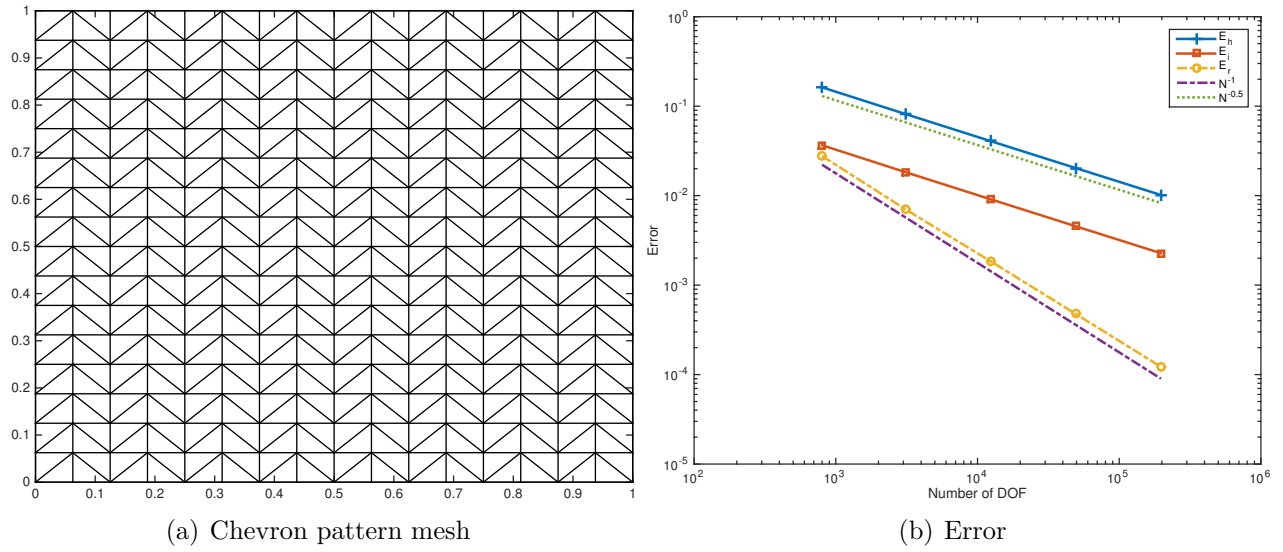


Figure 8: Numerical result of elliptic equation on chevron pattern mesh

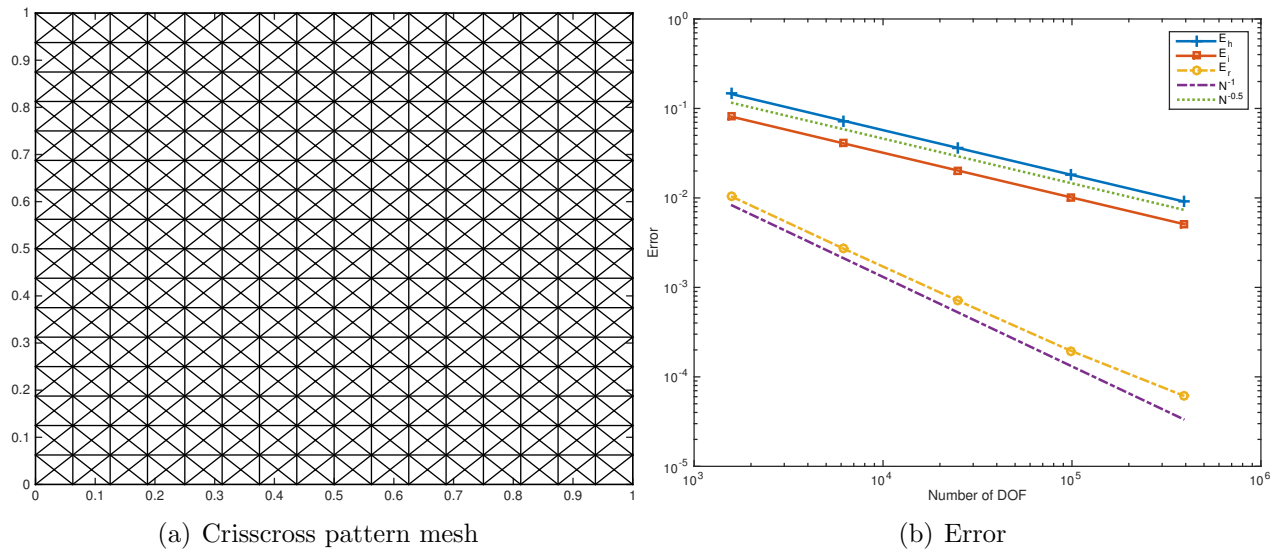


Figure 9: Numerical result of elliptic equation on crisscross pattern mesh

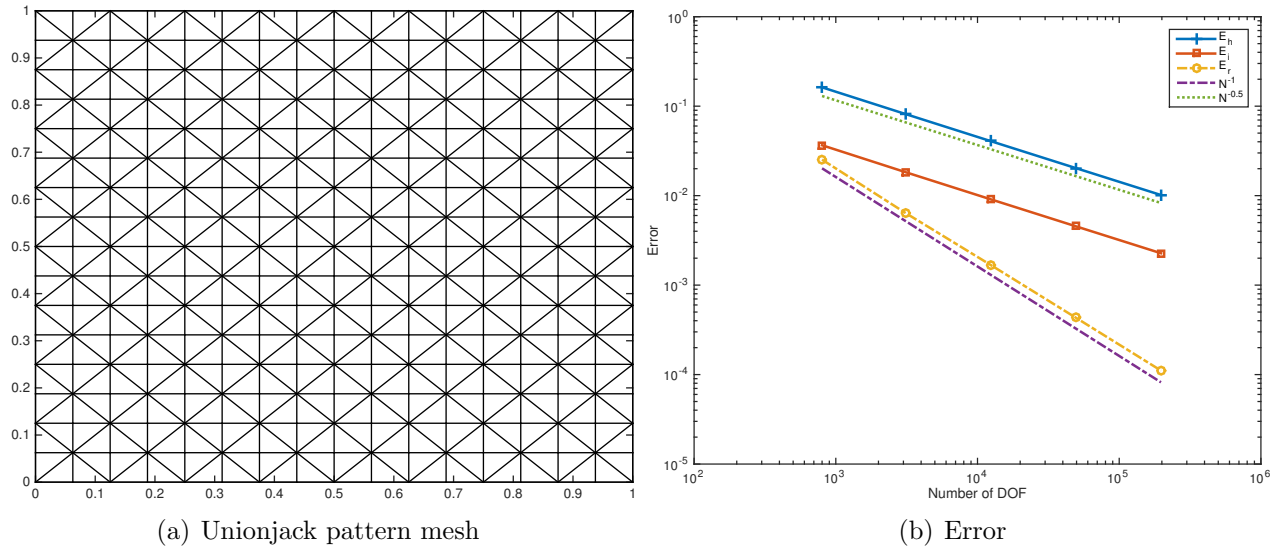


Figure 10: Numerical result of elliptic equation on unionjack pattern mesh

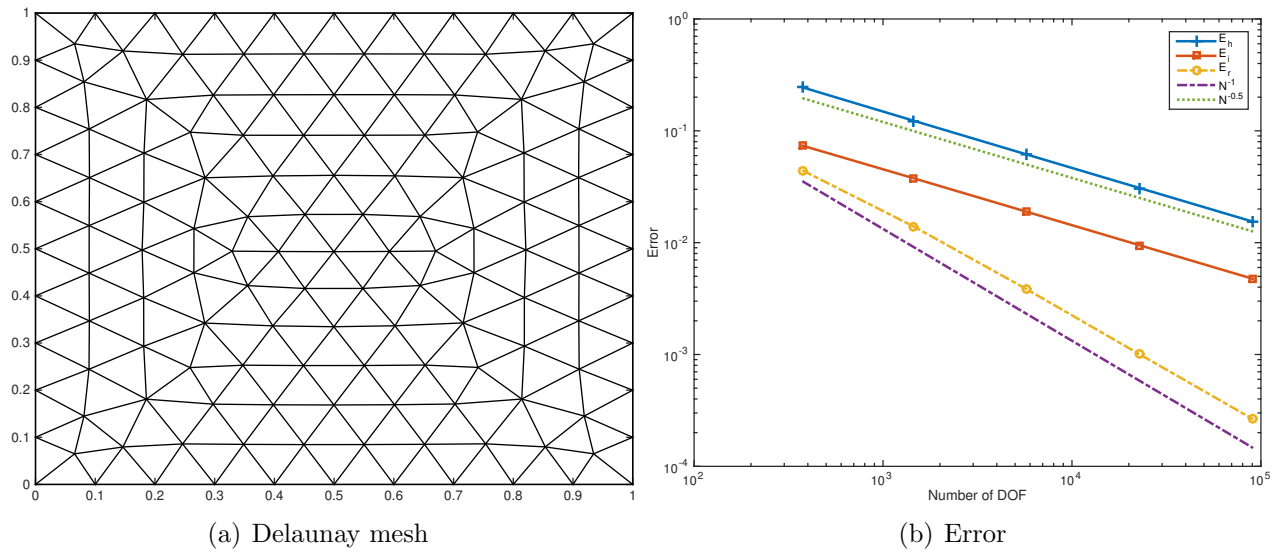


Figure 11: Numerical result of elliptic equation on delaunay mesh

Then we turn to unstructured mesh. We start from an initial mesh generated by EasyMesh [84], followed by four levels of refinement using bisection. As we can see in Figure 11, the rate of convergence for the recovered gradient in the  $L^2$  norm is very close to 2 and hence implies a superconvergent recovery.

From the above numerical results, we can observe the superconvergence for the proposed gradient recover operator on both structured meshes and unstructured mesh. Hence, it serves an asymptotically exact posteriori error estimators for adaptive finite element for the Crouzeix-Raviart element which will be studied at next section.

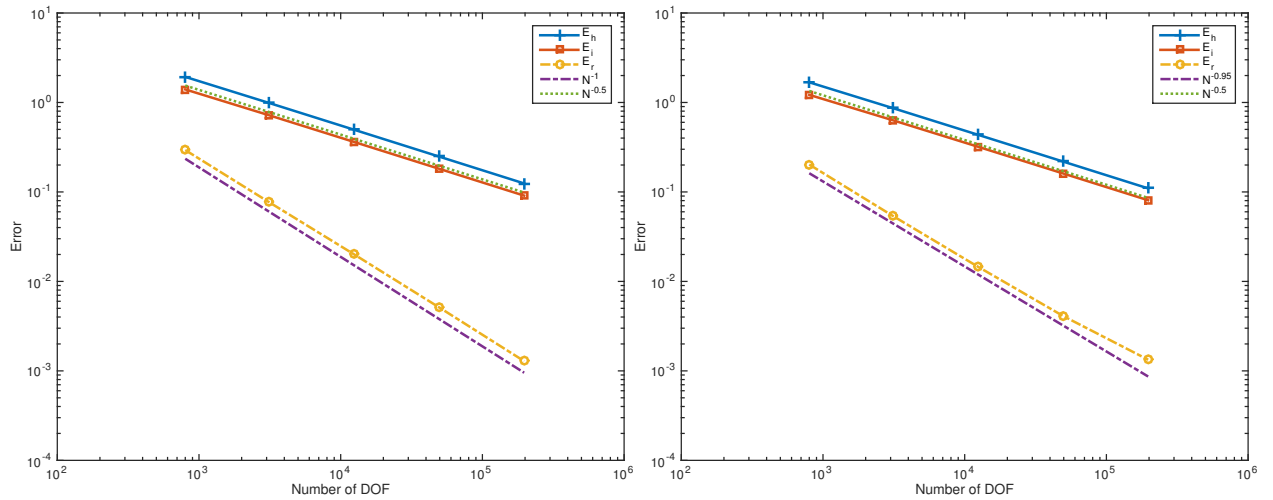


Figure 12: Numerical result of Stokes equation on regular pattern mesh  
 Figure 13: Numerical result of Stokes equation on chevron pattern mesh

### 3.4.2 Stokes equation

In the subsection, we consider the Stokes equation (2.3) on the unit square  $\Omega = [0, 1] \times [0, 1]$  with exact solution

$$\vec{u}(x, y) = \begin{pmatrix} 20xy^3 \\ 5x^4 - 5y^4 \end{pmatrix},$$

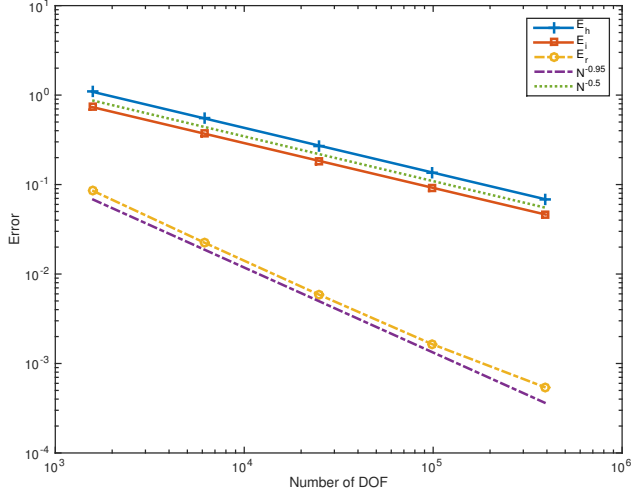


Figure 14: Numerical result of Stokes equation on crisscross pattern mesh

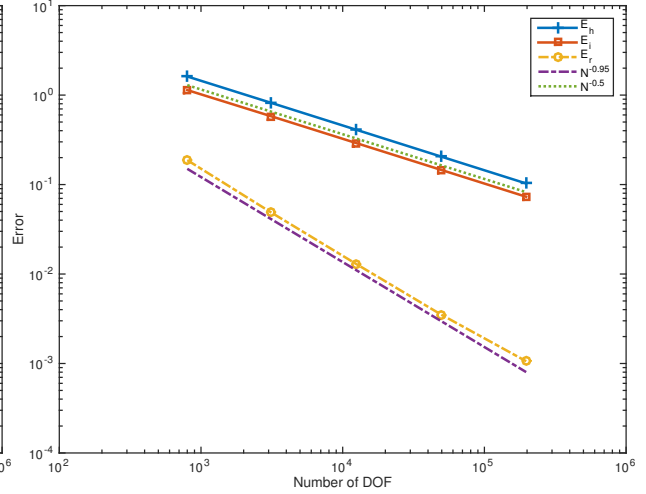


Figure 15: Numerical result of Stokes equation on unionjack pattern mesh

and  $p(x, y) = 60x^2y - 20y^3 - 5$ . Let  $(\vec{u}_h, p_h)$  be the nonconforming finite element approximation of the variational problem (2.4), i.e.  $(\vec{u}_h, p_h)$  is the solution of the discrete variational problem (2.10). Here we focus on the gradient recovery of velocity field  $\vec{u}_h$ .  $G_h\vec{u}_h$  means that the gradient recovery operator  $G_h$  is applied to each component of  $\vec{u}_h$ . According to Theorem 2.6, the optimal convergence rate of  $\|\vec{u}_h - \vec{u}\|_{1,h,\Omega}$  is  $O(h)$ . We can get superconvergence results by gradient recovery.

The numerical result for regular pattern uniform mesh is reported in Figure 12.  $O(h)$  convergence are observed for  $\|\vec{u} - \vec{u}_h\|_{1,h,\Omega}$  and  $\|\vec{u}_I - \vec{u}_h\|_{1,h,\Omega}$  for all kinds of meshes while  $O(h^2)$  superconvergence can be observed for  $\|G_h\vec{u}_h - \nabla\vec{u}\|_{0,\Omega_{h,1}}$ .

Figure 13 listed the numerical result for chevron pattern uniform mesh. The discrete  $H^1$  error and the discrete  $H^1$  norm of the difference between finite element solution  $\vec{u}_h$  and the nonconforming finite element interpolation  $\vec{u}_I$  of the exact solution  $\vec{u}$  are both about  $O(h)$ . Concerning the performance of gradient recovery operator,  $O(h^{1.9})$  superconvergence can be observed. Compared to the result on regular pattern uniform mesh, the superconvergence

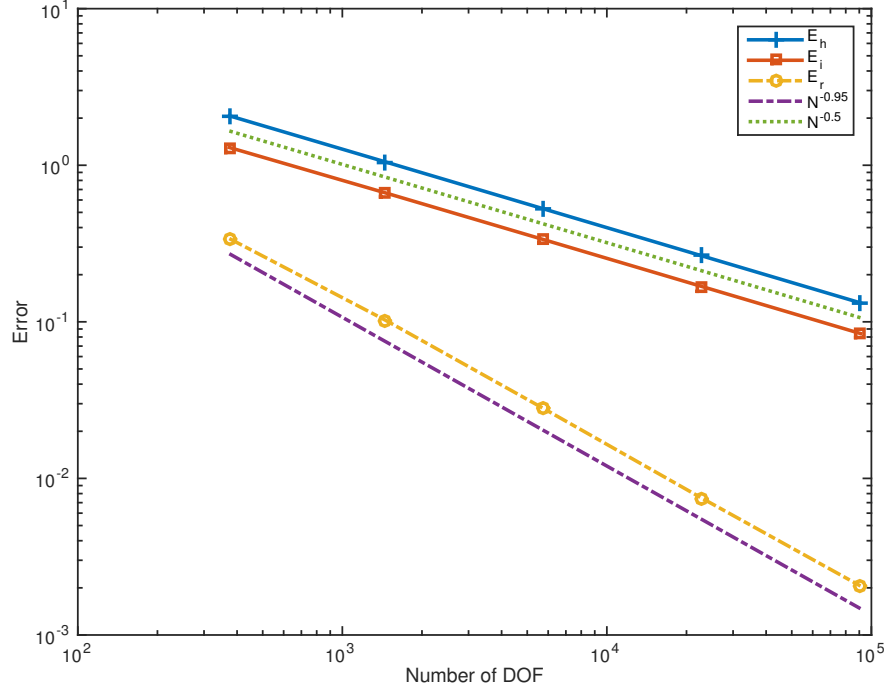


Figure 16: Numerical result of Stokes equation on Delaunay mesh

rate is a bit low.

Figure 14 and 15 present the numerical result for other two uniform meshes and Figure 16 shows the numerical result for Delaunay mesh. These numerical results are similar to chevron pattern uniform mesh. The recovered gradient superconverges at rate of  $O(h^{1.9})$  in the interior of the domain.

It is worth to point that  $\|\vec{u}_I - \vec{u}_h\|_{1,h,\Omega}$  is almost the same as  $\|\vec{u} - \vec{u}_h\|_{1,h,\Omega}$  for Stokes equation. Thus, supercloseness result is not true for nonconforming finite element approximation Stokes equation (2.10). Even in this case, we can get a superconvergent gradient by post-processing.

### 3.5 Application to adaptive methods

One of the most important applications of gradient recovery is in adaptive finite element methods. Adaptive finite element method (AFEM) is characterized by the loop of the form [86, 87]

**SOLVE**  $\rightarrow$  **ESTIMATE**  $\rightarrow$  **MARK**  $\rightarrow$  **REFINE**

More precisely, given an initial mesh  $\mathcal{T}_0$ , set  $k = 0$  and iterate

- **SOLVE.** Compute the solution  $u_k$  of discrete variational problem (2.9) on the Crouzeix-Raviart element space  $X^{\text{nc}}$  defined on the mesh  $\mathcal{T}_k$ .
- **ESTIMATE.** Compute the local error estimator  $\{\eta(u_k, T)\}_{T \in \mathcal{T}_k}$  using  $u_k$  and (or)  $\mathcal{T}_h$ .
- **MARK.** Collect a subset  $\mathcal{M}_k \subset \mathcal{T}_h$  of marked elements according the above posterior estimator and some marking strategy. In this work, only bulk marking strategy [38] is considered.
- **REFINE** Refine  $\mathcal{T}_k$  into a shape regular mesh and conforming mesh  $\mathcal{T}_{k+1}$  using bisection [15, 86, 87] in such a way that each element in  $\mathcal{M}_k$  is bisected at least once and, finally, increment  $k$ .

The essential part of AFEM is the **ESTIMATE** step. A posterior error estimators can be categorized into two classes: residue type and recovery type. There are extensive investigations of residue type a posteriori error estimators including both conforming finite element [4, 9, 10, 11, 15, 86, 87, 95, 96] and the Crouzeix-Raviart element [3, 39]. Concerning recovery type a posteriori estimators for conforming finite element method, the theory is also



relative mature, see [4, 29, 79, 98, 101, 110, 111, 112]. In particular, SPR and PPR become a standard part of several commercial finite element softwares. The study of recovery type a posteriori error estimator for nonconforming finite element methods is limited [27, 29, 42].

In this section, we apply the proposed gradient recovery technique to a recovery type a posteriori error estimate. Define a local a posteriori error estimator on the element  $T$  as

$$\eta(u_h, T) = \|\mathcal{D}^{\frac{1}{2}}(G_h u_h - \nabla u_h)\|_{0,T}, \quad (3.11)$$

and a global error estimator as

$$\eta(u_h, \Omega) = \left( \sum_{T \in \mathcal{T}_h} \eta(u_h, T) \right)^{\frac{1}{2}}. \quad (3.12)$$

In the context of adaptive finite element methods for boundary value problems, the effectivity index  $\kappa$  is used to measure the quality of an error estimator [4, 11]. This index is defined by the ratio between the estimated error and the true error

$$\kappa = \frac{\|G_h u_h - \nabla_h u_h\|_{0,\Omega}}{\|\nabla u - \nabla_h u_h\|_{0,\Omega}} \quad (3.13)$$

where  $\nabla_h$  is the broken gradient operator. To test the robustness of the error estimator (5.26), we use the first three examples in [30] as our benchmark problems. Readers are referred to [30] for computational comparison of several a posteriori error estimators for the Crouzeix-Raviart element.

**Example 2.5.1.** Let us consider the Laplace equation on the L-shaped domain  $\Omega = (-1, 1) \times (-1, 1) \setminus (0, 1) \times (-1, 0)$  with the Dirichlet boundary condition which is chosen

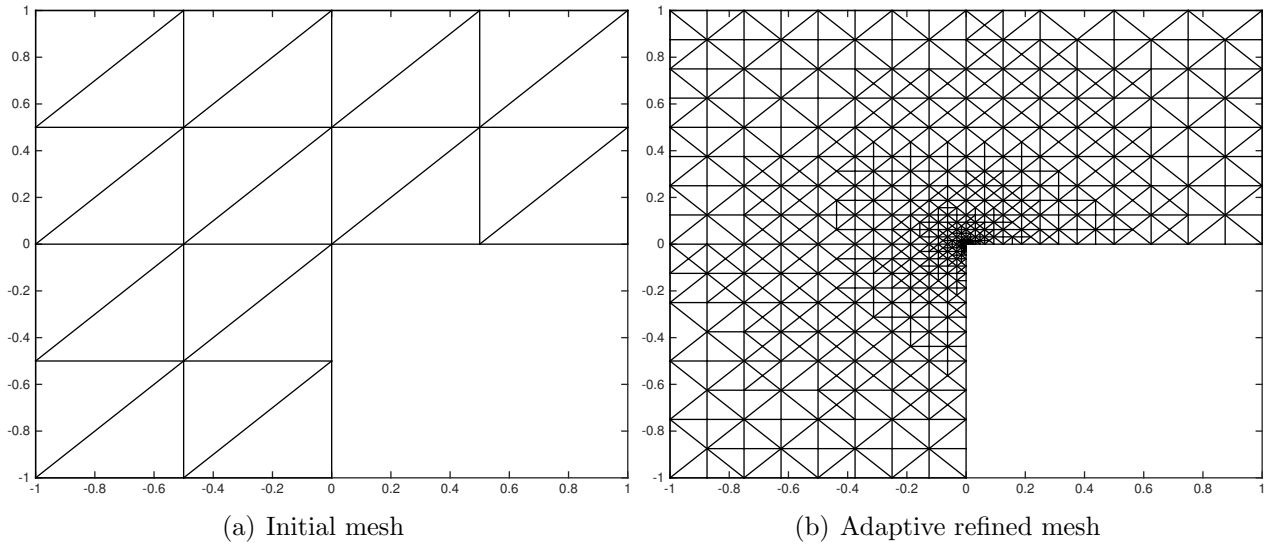


Figure 17: Meshes of Poisson equation on L-shaped domain

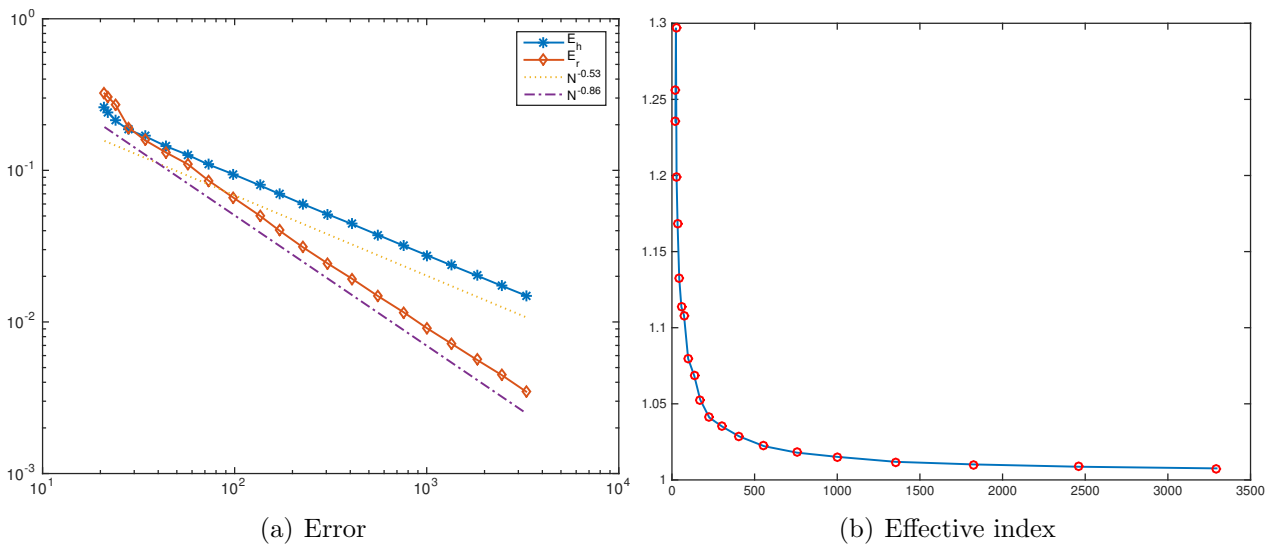


Figure 18: Numerical results of Poisson equation on L-shaped domain

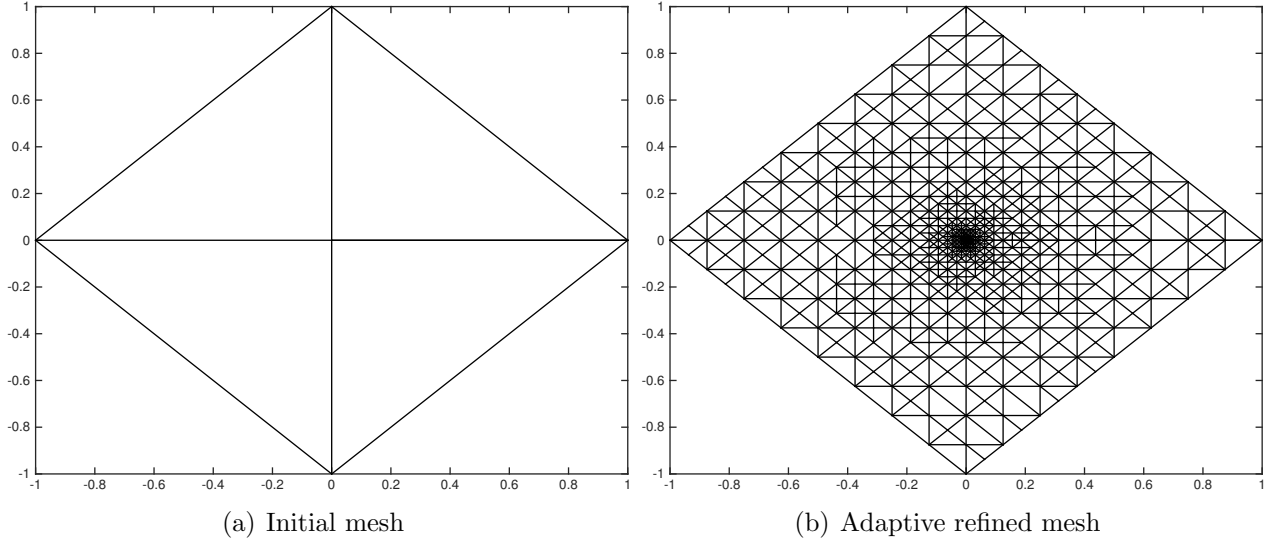


Figure 19: Meshes of Poisson equation on Crack domain

so that the true solution is  $r^{2/3} \sin(2\theta/3)$  in polar coordinates. The solution has a corner singularity at  $(0,0)$ . To obtain optimal convergence rate  $O(h)$ , we use adaptive finite element method. Fig 17(a) shows the initial mesh while Fig 17(b) plots the adaptive refined mesh. The mesh is locally refined at the singularity point. We numerically observed from Fig 18(a) that

$$\|u - \nabla u_h\|_{1,h,\Omega} \approx O(N^{-0.54}) \quad \|u - G_h u_h\|_{0,\Omega} \approx O(N^{-0.86}).$$

Notice that  $\|u - u_h\|_{1,h,\Omega}$  converges at optimal rate.

The effectivity index are plotted in Fig 18(b). We see that  $\kappa$  converges to 1 quickly after the first few iterations which indicates the posteriori error estimator (5.26) or (3.12) is asymptotically exact.

**Example 2.5.2.** The second benchmark problem is elliptic equation (2.1) with  $\mathcal{D} = I$  and  $c = 0$  on the crack domain  $\Omega = \{|x| + |y| \leq 1\} \setminus \{0 \leq x \leq 1, y = 0\}$ . The right hand side function is 1 and the exact solution is  $u = \sqrt{\frac{1}{2}(r-x)} - \frac{1}{4}r^2$  with  $r = \sqrt{x^2 + y^2}$ . The

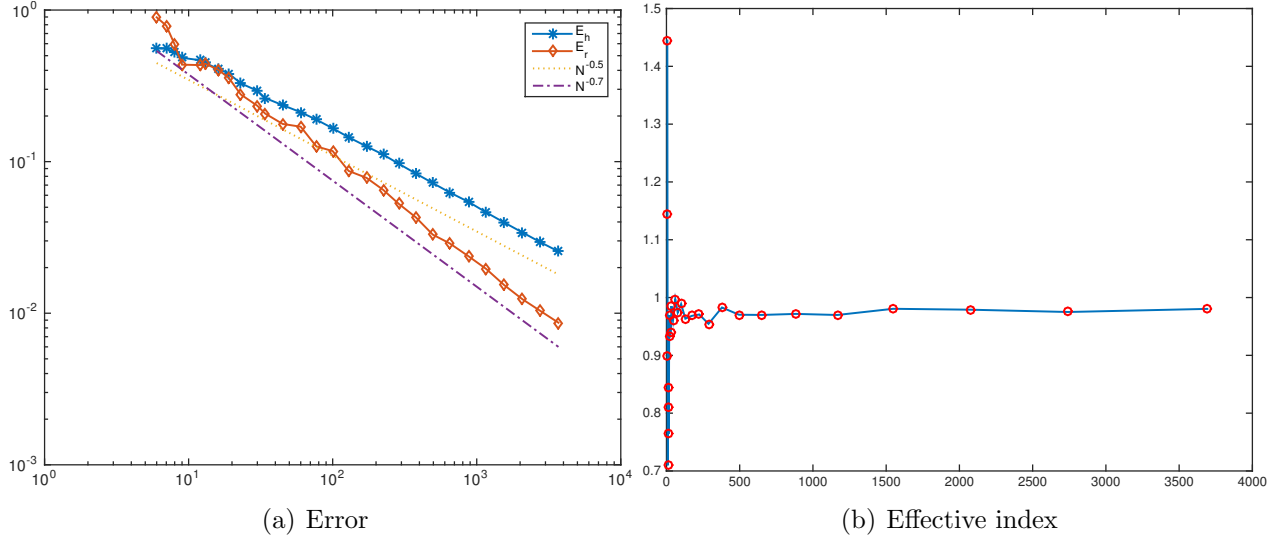


Figure 20: Numerical results of Poisson equation on Crack domain

initial mesh and the adaptive refined mesh is plotted in Fig 19. Fig 20(a) shows that  $H^1$  error of numerical solution is optimal and the recovery gradient error superconverges at rate of  $O(h^{1.4})$ . Similarly to previous example, the error estimator is asymptotically exact which is indicated by the effective index closing to 1, see Figure 20(b).

**Example 2.5.3.** Our third benchmark problem employs homogeneous boundary data and an oscillation source term  $f$  that matches the exact solution  $u(x, y) = x(x - 1)y(y - 1) \exp(-100(x - 1/2)^2 - 100(y - 117/1000)^2)$  on the square domain  $\Omega = [0, 1]^2$ , see[30, 74]. We use the initial mesh as in Figure 21(a) and the resulting adaptive refined mesh is in Figure 21(b) which is locally refined near the oscillation point. The numerical result is presented in Figure 22. Since the solution is smooth, we observe  $O(h^2)$  superconvergence for gradient recovery error. The estimator is asymptotically exact.

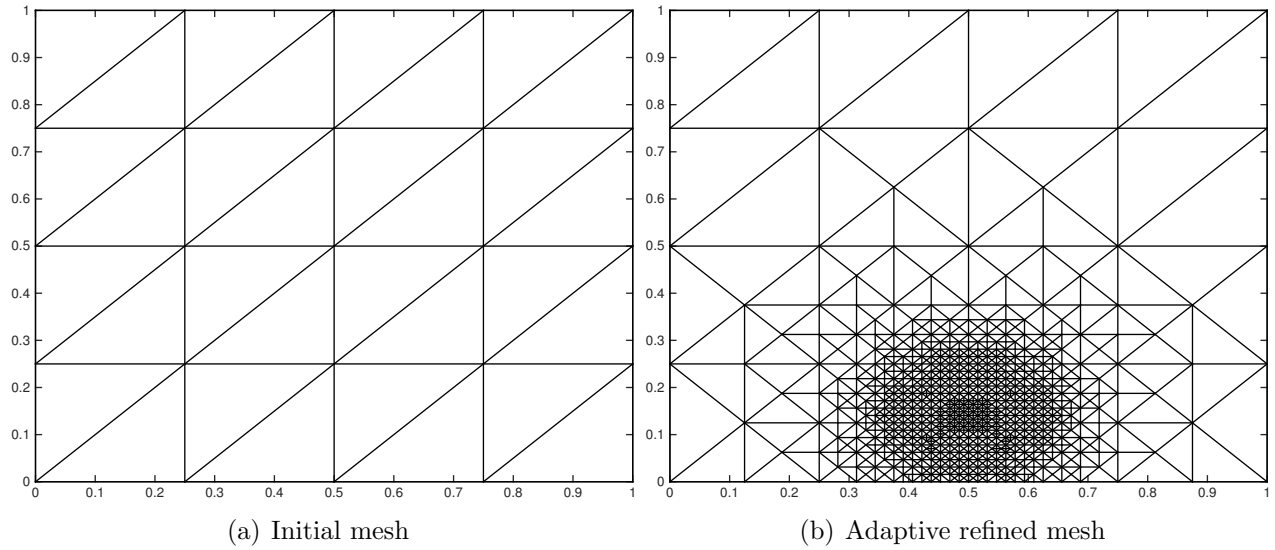


Figure 21: Meshes of Poisson equation on square domain with oscillations

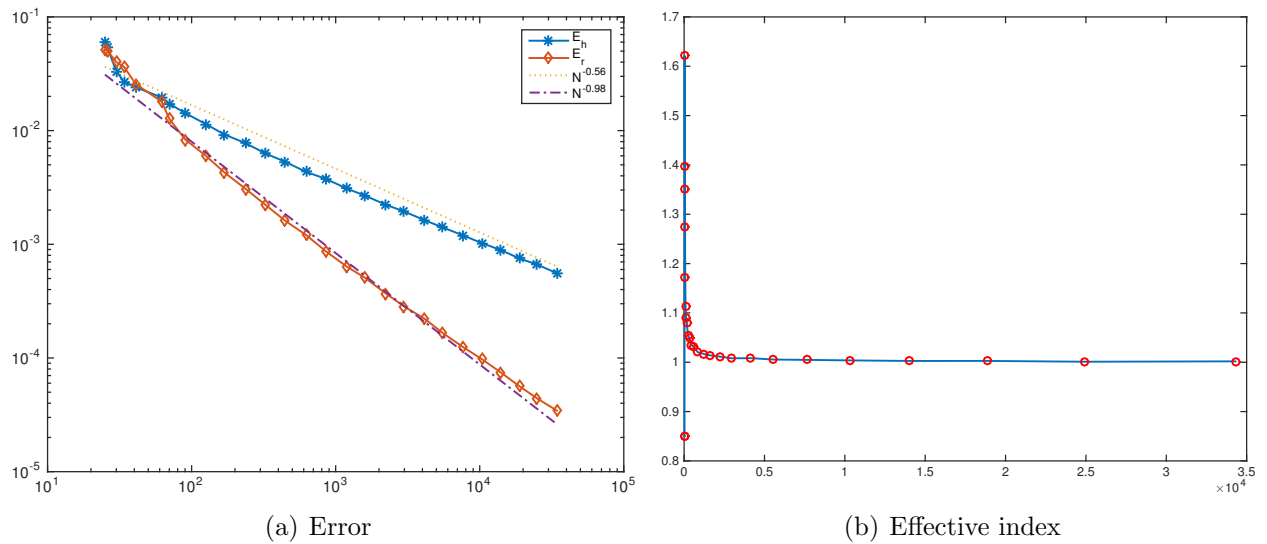


Figure 22: Numerical Results of Poisson equation on square domain with oscillations

### 3.6 Conclusion

We proposed a gradient recovery method for the Crouzeix-Raviart element. The proposed method fits a quadratic polynomial in least square sense at any edge center and then take derivative to get recovered gradient. We proved that it is polynomial preserving and is a bounded linear operator. Numerical experiment showed that it produces superconvergent recovered gradient and can serve as an asymptotically exact posteriori error estimator.

## CHAPTER 4 Hessian recovery for finite element

Hessian matrix has many applications in scientific computing [61, 65, 66, 83, 45]. which is typically unavailable in a numerical simulation. A widely-used approach to avoid this difficulty in practical computation is to replace the information by one recovered from the obtained numerical approximation [60, 89, 94]. However, there is no general theory guaranteeing convergence for existing Hessian recovery methods.

In this chapter, we study Hessian recovery for  $C_0$  finite element methods of arbitrary order. Our approach is to apply PPR twice to the primary computed data. We proved that the proposed Hessian recovery method preserves polynomials of degree  $r + 1$  on general unstructured meshes and superconvergence at a rate of  $o(h^r)$  on mildly structured meshes. In addition, the method is proved to be ultraconvergent (two order higher) for translation invariant finite element space of any order.

In Section 4.1, we introduce some notation. Since the building block of our Hessian recovery is polynomial preserving recovery (PPR), we describe the definition of PPR in Section 4.2. The formal definition of our Hessian recovery operator is provided in Section 4.3. This definition is illustrated with two examples on uniform meshes and followed by discussion of its properties. The main superconvergence analysis of the proposed Hessian recovery method is shown in Section 4.4. In Section 4.5, we use some numerical examples to verify our theoretical results. We end this chapter with some concluding remarks in Section 4.6.

## 4.1 Notation

To simplify notation, the Hessian operator  $H$  is denoted by

$$H = \begin{pmatrix} \partial_{xx} & \partial_{xy} \\ \partial_{yx} & \partial_{yy} \end{pmatrix}. \quad (4.1)$$

For any  $0 < h < \frac{1}{2}$ , let  $\mathcal{T}_h$  be a shape regular triangulation of  $\bar{\Omega}$  as defined in section 2.3.1. For any  $r \in \mathbb{N}$ , let  $S^{h,r}$  be the continuous finite element space with piecewise polynomial of degree  $r$ , see section 2.3.2. In this section, we suppose  $\mathcal{T}_h$  is quasi-uniform. Let  $\mathcal{N}_h$  denote the set of mesh nodes, i.e. the dual space of  $S^{h,r}$ .

Then we talk about a special finite element space, which is widely used for superconvergence analysis [97]. For  $\mathcal{A} \subset \Omega \subset \mathbb{R}^2$ , let  $S^{h,r}(\mathcal{A})$  denote the restrictions of functions in  $S^{h,r}$  to  $\mathcal{A}$  and let  $S_{00}^{h,r}(\mathcal{A})$  denote the set of those functions in  $S^{h,r}(\mathcal{A})$  with compact support in the interior of  $\mathcal{A}$  [97]. Let  $\Omega_0 \subset\subset \Omega_1 \subset\subset \Omega_2 \subset\subset \Omega$  be separated by  $d \geq c_0 h$  and  $\ell$  be a direction, i.e., a unit vector in  $\mathbb{R}^2$ . Let  $\tau$  be a parameter, which will typically be a multiple of  $h$ . Let  $T_\tau^\ell$  denote translation by  $\tau$  in the direction  $\ell$ , i.e.,

$$T_\tau^\ell v(z) = v(z + \tau\ell), \quad (4.2)$$

and for an integer  $\nu$

$$T_{\nu\tau}^\ell v(z) = v(z + \nu\tau\ell). \quad (4.3)$$

Following the definition of [97], the finite element space  $S^{h,r}$  is called translation invariant



by  $\tau$  in the direction  $\ell$  if

$$T_{\nu\tau}^\ell v \in S_{00}^{h,r}(\Omega), \quad \forall v \in S_{00}^{h,r}(\Omega_1), \quad (4.4)$$

for some integer  $\nu$  with  $|\nu| < M$ . Equivalently,  $\mathcal{T}_h$  is called a translation invariant mesh.

To clarify the matter, we consider five popular triangular mesh patterns: Regular, Chevron, Union-Jack, Criss-cross, and equilateral patterns, as shown in Figure 23.

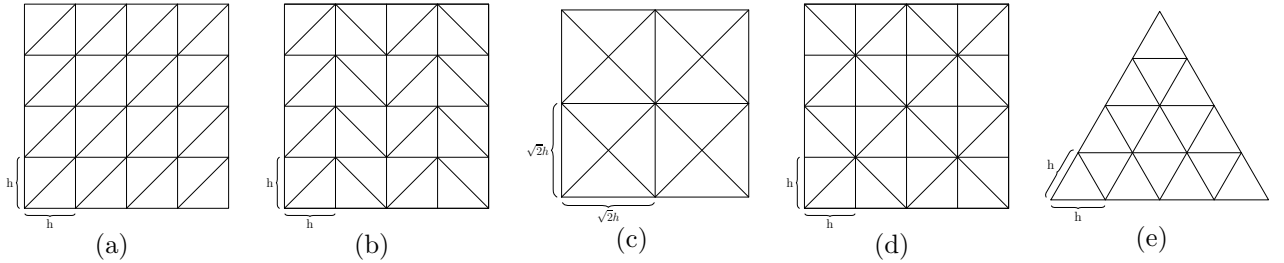


Figure 23: Five types of uniform meshes: (a) Regular pattern; (b) Chevron pattern; (c) Criss-cross pattern; (d) Union-Jack pattern; (e) Equilateral pattern

We see that:

- 1) Regular pattern is translation invariant by  $h$  in directions  $(1, 0)$  and  $(0, 1)$ , by  $2\sqrt{2}h$  in directions  $(\pm\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2})$ , and by  $\sqrt{5}h$  in directions  $(\frac{2\sqrt{5}}{5}, \pm\frac{\sqrt{5}}{5})$  and  $(\pm\frac{\sqrt{5}}{5}, \frac{2\sqrt{5}}{5})$ , .....
- 2) Chevron pattern is translation invariant by  $h$  in the direction  $(0, 1)$ , by  $2h$  in the direction  $(1, 0)$ , and by  $2\sqrt{2}h$  in directions  $(\pm\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2})$ , and by  $\sqrt{5}h$  in directions  $(\pm\frac{\sqrt{5}}{5}, \frac{2\sqrt{5}}{5})$ , .....
- 3) Criss-cross pattern is translation invariant by  $\sqrt{2}h$  in directions  $(1, 0)$  and  $(0, 1)$ , and by  $2h$  in directions  $(\pm\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2})$ , .....
- 4) Union-Jack pattern is translation invariant by  $2h$  in directions  $(1, 0)$  and  $(0, 1)$ , and by  $2\sqrt{2}h$  in directions  $(\pm\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2})$ , .....
- 5) Equilateral pattern is translation invariant by  $h$  in directions  $(1, 0)$  and  $(\pm\frac{1}{2}, \frac{\sqrt{3}}{2})$ , and

by  $\sqrt{3}h$  in directions  $(0, 1)$  and  $(\frac{\sqrt{3}}{2}, \pm\frac{1}{2})$ , .....

## 4.2 Polynomial preserving recovery

Let  $G_h : S^{h,r} \rightarrow \prod_{i=1}^2 S^{h,r}$  denote the PPR gradient recovery operator [79, 80, 107]. Given a function  $u_h \in S^{h,r}$ , it suffices to define  $(G_h u_h)(z)$  for all  $z \in \mathcal{N}_h$ . Let  $z \in \mathcal{N}_h$  be a vertex and  $\mathcal{K}_z$  be a patch of elements around  $z$  which is defined in [80, 107]. Select all nodes in  $\mathcal{N}_h \cap \mathcal{K}_z$  as sampling points and fit a polynomial  $p_z \in \mathbb{P}_{r+1}(\mathcal{K}_z)$  in the least squares sense at those sampling points, i.e.

$$p_z = \arg \min_{p \in \mathbb{P}_{r+1}(\mathcal{K}_z)} \sum_{\tilde{z} \in \mathcal{N}_h \cap \mathcal{K}_z} (u_h - p)^2(\tilde{z}). \quad (4.5)$$

We called  $p_z$  is the least-square polynomial approximation (LSPA) of  $u_h$  at  $z$ . The recovered gradient at  $z$  is defined as

$$(G_h u_h)(z) = \nabla p_z(z).$$

For linear element, all nodes in  $\mathcal{N}_h$  are vertices and hence  $G_h u_h$  is well defined. However,  $\mathcal{N}_h$  may contain edge nodes or interior nodes for higher order elements. If  $z$  is an edge node which lies on an edge between two vertices  $z_1$  and  $z_2$ , we define

$$(G_h u_h)(z) = \beta \nabla p_{z_1}(z) + (1 - \beta) \nabla p_{z_2}(z)$$

where  $\beta$  is determined by the ratio of distances of  $z$  to  $z_1$  and  $z_2$ . If  $z$  is an interior node which lies in a triangle formed by three vertices  $z_1$ ,  $z_2$ , and  $z_3$ , we define

$$(G_h u_h)(z) = \sum_{j=1}^3 \beta_j \nabla p_{z_j}(z),$$

where  $\beta_j$  is the barycentric coordinate of  $z$ .

Zhang and Naga proved the following properties of PPR in [79, 80, 107]:

1.  $G_h$  is linear.
2.  $G_h$  satisfied the consistency condition, i.e.

$$G_h(I_h p) = \nabla p, \forall p \in \mathbb{P}_{r+1}(\Omega). \quad (4.6)$$

3.  $G_h$  is bounded in the following sense:

$$\|G_h v\|_{0,T} \lesssim |v|_{1,\mathcal{K}_T}, \quad \forall T \in \mathcal{T}_h, \text{ and } \forall v \in S^{h,r} \quad (4.7)$$

where

$$\mathcal{K}_T = \bigcup \{\mathcal{K}_z : z \text{ is a vertice of } T\}$$

is the patch corresponding to  $T$ . This condition is called the boundedness condition.

**Remark.** It was proved in [79] that certain rank condition and geometric condition guarantee the uniqueness of  $p_z$  in (4.5).

**Remark.** In order to avoid numerical instability, a discrete least squares fitting process is carried out on a reference patch  $\omega_z$ .

### 4.3 Hessian recovery method

Given  $u \in S^{h,r}$ , let  $G_h u \in \prod_{i=1}^2 S_h$  be the recovered gradient using PPR as defined in previous section. We rewrite  $G_h u$  as

$$G_h u = \begin{pmatrix} G_h^x u \\ G_h^y u \end{pmatrix}. \quad (4.8)$$

In order to recover the Hessian matrix of  $u$ , we apply gradient recovery operator  $G_h$  to  $G_h^x u$  and  $G_h^y u$  one more time, respectively, and define the Hessian recovery operator  $H_h$  as follows

$$H_h u = \begin{pmatrix} G_h(G_h^x u), & G_h(G_h^y u) \end{pmatrix} = \begin{pmatrix} G_h^x(G_h^x u) & G_h^x(G_h^y u) \\ G_h^y(G_h^x u) & G_h^y(G_h^y u) \end{pmatrix}. \quad (4.9)$$

Just as PPR, we obtain  $H_h : S^{h,r} \rightarrow \prod_{i=1}^2 S^{h,r} \times \prod_{i=1}^2 S^{h,r}$  on the whole domain  $\Omega$  by interpolation after determining values of  $H_h u$  at all nodes in  $\mathcal{N}_h$ .

**Remark.** The two gradient recovery operators in definition (4.9) of  $H_h$  can be different.

Actually we can define the Hessian recovery operator  $H_h$  as following

$$H_h u = \begin{pmatrix} \tilde{G}_h(G_h^x u), & \tilde{G}_h(G_h^y u) \end{pmatrix}.$$

By choosing  $G_h$  and  $\tilde{G}_h$  as PPR or SPR operator, we obtain four different Hessian recovery operators, i.e., PPR-PPR, PPR-SPR, SPR-PPR, and SPR-SPR. However, numerical tests have shown that PPR-PPR is the best one.

In order to demonstrate our method, we shall discuss two examples in detail. For the sake of simplicity, only linear element on uniform mesh will be considered. In practice, the method can be applied to arbitrary mesh and higher order elements.

**Example 4.1.** Consider the regular pattern uniform mesh as in Figure 24(a). We want to recovery the Hessian matrix at  $z_0$ . As deduced in [107], the recovered gradient at  $z_0$  is given by

$$(G_h u)(z_0) = \frac{1}{6h} \left( \begin{pmatrix} 2 \\ 1 \end{pmatrix} u_1 + \begin{pmatrix} 1 \\ 2 \end{pmatrix} u_2 + \begin{pmatrix} -1 \\ 1 \end{pmatrix} u_3 - \begin{pmatrix} 2 \\ 1 \end{pmatrix} u_4 - \begin{pmatrix} 1 \\ 2 \end{pmatrix} u_5 + \begin{pmatrix} 1 \\ -1 \end{pmatrix} u_6 \right).$$

Here  $u_i = u(z_i)$ , ( $i = 0, 1, \dots, 18$ ) represents function value of  $u$  at node  $z_i$ . Thus, according to the definition (4.9) of the Hessian recovery operator  $H_h$ , we have

$$\begin{pmatrix} H_h^{xx} u \\ H_h^{xy} u \end{pmatrix} (z_0) = \frac{1}{6h} (2(G_h u)(z_1) + (G_h u)(z_2) - (G_h u)(z_3) - 2(G_h u)(z_4) - (G_h u)(z_5) + (G_h u)(z_6)), \quad (4.10)$$

and

$$\begin{pmatrix} H_h^{yx} u \\ H_h^{yy} u \end{pmatrix} (z_0) = \frac{1}{6h} ((G_h u)(z_1) + 2(G_h u)(z_2) + (G_h u)(z_3) - (G_h u)(z_4) - 2(G_h u)(z_5) - (G_h u)(z_6)), \quad (4.11)$$

where

$$(G_h u)(z_1) = \frac{1}{6h} \left( \begin{pmatrix} 2 \\ 1 \end{pmatrix} u_7 + \begin{pmatrix} 1 \\ 2 \end{pmatrix} u_8 + \begin{pmatrix} -1 \\ 1 \end{pmatrix} u_2 - \begin{pmatrix} 2 \\ 1 \end{pmatrix} u_0 - \begin{pmatrix} 1 \\ 2 \end{pmatrix} u_{18} + \begin{pmatrix} 1 \\ -1 \end{pmatrix} u_6 \right),$$

and  $(G_h u)(z_2), \dots, (G_h u)(z_6)$  follow the similar pattern. Direct calculation reveals that

$$(H_h^{xx} u)(z_0) = \frac{1}{36h^2} (-12u_0 + 2u_1 - 4u_2 - 4u_3 + 2u_4 - 4u_5 - 4u_6 + 4u_7 + 4u_8 + u_9 - 2u_{10} + u_{11} + 4u_{12} + 4u_{13} + 4u_{14} + u_{15} - 2u_{16} + u_{17} + 4u_{18}),$$

$$(H_h^{xy} u)(z_0) = \frac{1}{36h^2} (6u_0 - u_1 + 5u_2 - u_3 - u_4 + 5u_5 - u_6 - 2u_7 + u_8 + u_9 + u_{10} - 2u_{11} - 5u_{12} - 2u_{13} + u_{14} + u_{15} + u_{16} - 2u_{17} - 5u_{18}),$$

$$(H_h^{yx} u)(z_0) = \frac{1}{36h^2} (6u_0 - u_1 + 5u_2 - u_3 - u_4 + 5u_5 - u_6 - 2u_7 + u_8 + u_9 + u_{10} - 2u_{11} - 5u_{12} - 2u_{13} + u_{14} + u_{15} + u_{16} - 2u_{17} - 5u_{18}),$$

$$(H_h^{yy} u)(z_0) = \frac{1}{36h^2} (-12u_0 - 4u_1 - 4u_2 + 2u_3 - 4u_4 - 4u_5 + 2u_6 + u_7 - 2u_8 + u_9 + 4u_{10} + 4u_{11} + 4u_{12} + u_{13} - 2u_{14} + u_{15} + 4u_{16} + 4u_{17} + 4u_{18}).$$

It is observed that  $(H_h^{xy} u)(z_0) = (H_h^{yx} u)(z_0)$ , which means the recovered Hessian matrix is symmetric, a property of the exact Hessian we would like to maintain.

Using Taylor expansion, we can show that

$$\begin{aligned}
(H_h^{xx}u)(z_0) &= u_{xx}(z_0) + \frac{h^2}{3}(u_{xxxx}(z_0) + u_{xxxy}(z_0) + u_{xxyy}(z_0)) + O(h^4), \\
(H_h^{xy}u)(z_0) &= u_{xy}(z_0) + \frac{h^2}{3}(u_{xxxxy}(z_0) + u_{xxxyy}(z_0) + u_{xyyy}(z_0)) + O(h^4), \\
(H_h^{yx}u)(z_0) &= u_{yx}(z_0) + \frac{h^2}{3}(u_{xxxxy}(z_0) + u_{xxxyy}(z_0) + u_{xyyy}(z_0)) + O(h^4), \\
(H_h^{yy}u)(z_0) &= u_{yy}(z_0) + \frac{h^2}{3}(u_{xxyyy}(z_0) + u_{xyyy}(z_0) + u_{yyyy}(z_0)) + O(h^4),
\end{aligned}$$

which imply that  $H_h u$  provides a second order approximation of  $Hu$  at  $z_0$ .

**Example 4.2.** Consider the Chevron pattern uniform mesh as shown in Figure 24(b).

Repeating the procedure as in Example 4.1, we derive the recovered Hessian matrix at  $z_0$  as

$$\begin{aligned}
(H_h^{xx}u)(z_0) &= \frac{1}{144h^2}(-72u_0 + 36u_{13} + 36u_7), \\
(H_h^{xy}u)(z_0) &= \frac{1}{144h^2}(-12u_1 + 12u_3 + 24u_4 - 24u_6 + 6u_7 + \\
&\quad + 36u_9 - 36u_{11} - 6u_{13} + 6u_{14} - 6u_{18}), \\
(H_h^{yx}u)(z_0) &= \frac{1}{144h^2}(12u_1 - 12u_3 + 36u_4 - 36u_6 - 6u_7 + \\
&\quad + 6u_8 + 24u_9 - 24u_{11} - 6u_{12} + 6u_{13}), \\
(H_h^{yy}u)(z_0) &= \frac{1}{144h^2}(-48u_0 - 10u_1 - 22u_2 - 10u_3 - 10u_4 + 18u_5 - \\
&\quad - 10u_6 - 2u_7 + u_8 + 10u_9 + 36u_{10} + 10u_{11} + u_{12} - \\
&\quad - 2u_{13} + u_{14} + 10u_{15} + 16u_{16} + 10u_{17} + u_{18}).
\end{aligned}$$

In addition, we have the following Taylor expansion

$$\begin{aligned}
(H_h^{xx}u)(z_0) &= u_{xx}(z_0) + \frac{h^2}{3}u_{xxxx}(z_0) + \frac{2h^4}{45}u_{xxxxxx}(z_0) + O(h^5), \\
(H_h^{xy}u)(z_0) &= u_{xy}(z_0) + \frac{h^2}{12}(3u_{xxyy}(z_0) + 2u_{xyyy}(z_0)) - \frac{h^3}{24}u_{xxyyy}(z_0) + O(h^4), \\
(H_h^{yx}u)(z_0) &= u_{yx}(z_0) + \frac{h^2}{12}(3u_{xxyy}(z_0) + 2u_{xyyy}(z_0)) + \frac{h^3}{24}u_{xxyyy}(z_0) + O(h^4), \\
(H_h^{yy}u)(z_0) &= u_{yy}(z_0) + \frac{h^2}{6}(u_{xxyy}(z_0) + 2u_{yyyy}(z_0)) - \frac{5h^3}{72}u_{xxyyy}(z_0) + O(h^4).
\end{aligned}$$

We conclude that  $H_h u$  is a second order approximation to the Hessian matrix. It is worth pointing out that, though  $H_h^{xy} \neq H_h^{yx}$  for the Chevron pattern uniform mesh, they are both second order finite difference schemes at  $z_0$ .

**Remark.** PPR-PPR is the only one among the four Hessian recovery methods mentioned in Remark 4.3 that provides second order approximation for all five mesh patterns, especially the Chevron pattern.

Both Example 4.1 and 4.2 indicate that for linear element the PPR-PPR approach is equivalent to a finite difference scheme of second order accuracy at vertex  $z_0$ . In general, we can show that  $H_h$  preserves polynomials of degree up to  $k + 1$  for  $k$ th order element.

Consider  $P_k$ -element. Let  $u$  be a polynomial of degree  $k + 1$ . Since  $G_h$  preserves polynomials of degree  $k + 1$ , it follows that  $G_h u = \nabla u$  which is a polynomial of degree  $k$ . Therefore, we have

$$H_h u = (G_h(G_h^x u), G_h(G_h^y u)) = (G_h \frac{\partial u}{\partial x}, G_h \frac{\partial u}{\partial x}) = (\nabla \frac{\partial u}{\partial x}, \nabla \frac{\partial u}{\partial x}) = H u. \quad (4.12)$$

It means that  $H_h$  preserves polynomials of degree  $k + 1$  for arbitrary mesh.



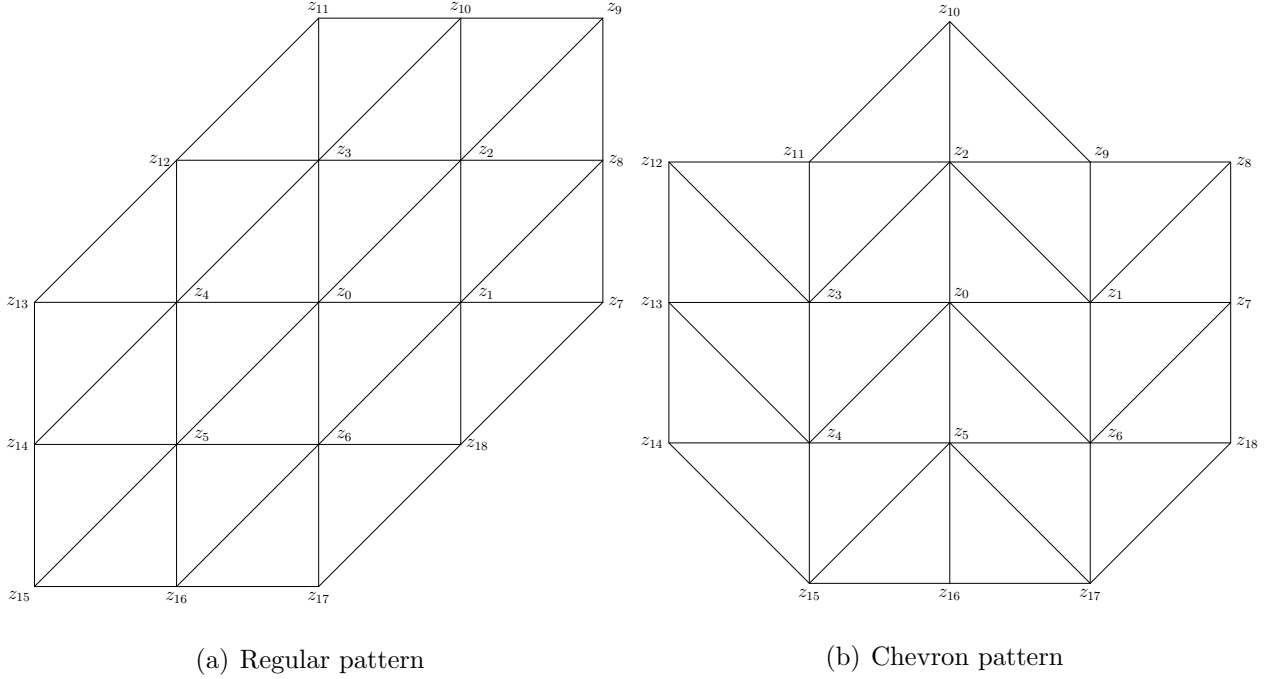


Figure 24: Illustration on Hessian recovery on uniform mesh

Now we proceed to translation invariant mesh. According to the polynomial preserving property (4.6), the recovered gradient is exact for polynomials of degree  $k + 1$ . Therefore

$$G_h^x u = D_x u + h^{k+1} \mathbf{a}^x \cdot D^{k+2} u + h^{k+2} \mathbf{b}^x \cdot D^{k+3} u + h^{k+3} \mathbf{c}^x \cdot D^{k+4} u + \dots ; \quad (4.13)$$

$$G_h^y u = D_y u + h^{k+1} \mathbf{a}^y \cdot D^{k+2} u + h^{k+2} \mathbf{b}^y \cdot D^{k+3} u + h^{k+3} \mathbf{c}^y \cdot D^{k+4} u + \dots . \quad (4.14)$$

Note that  $\mathbf{a}^x, \mathbf{a}^y, \mathbf{b}^x, \mathbf{b}^y, \mathbf{c}^x, \mathbf{c}^y, \dots$  are functions of  $(x, y)$  if  $z = (x, y)$  a nodal point of arbitrary mesh.

Let  $\mathbf{z} = (x, y)$  be any node on a translation invariant mesh. Notice that coefficients  $\mathbf{a}^x, \mathbf{a}^y, \mathbf{b}^x, \mathbf{b}^y, \dots$  depend only on the coordinates of nodes, since we recover gradient at nodes

only. Thus for translation invariant meshes,  $\mathbf{a}^x$ ,  $\mathbf{a}^y$ ,  $\mathbf{b}^x$ ,  $\mathbf{b}^y$ ,  $\dots$  are constants. Note that

$$\begin{aligned}
(H_h^{xy}u)(\mathbf{z}) &= (G_h^y(G_h^x u))(\mathbf{z}) \\
&= G_h^y[D_x u(\mathbf{z}) + h^{k+1}\mathbf{a}^x \cdot D^{k+2}u(\mathbf{z}) + h^{k+2}\mathbf{b}^x \cdot D^{k+3}u(\mathbf{z}) + \dots] \\
&= (G_h^y(D_x u))(\mathbf{z}) + h^{k+1}(\mathbf{a}^x \cdot G_h^y(D^{k+2}u))(\mathbf{z}) + h^{k+2}(\mathbf{b}^x \cdot G_h^y(D^{k+3}u))(\mathbf{z}) + \dots \\
&= (D_y D_x u)(\mathbf{z}) + h^{k+1}(\mathbf{a}^y \cdot D^{k+2}D_x u)(\mathbf{z}) + h^{k+2}(\mathbf{b}^y \cdot D^{k+3}D_x u)(\mathbf{z}) \\
&\quad + h^{k+1}(\mathbf{a}^x \cdot D_y(D^{k+2}u))(\mathbf{z}) + h^{k+2}(\mathbf{b}^x \cdot D_y(D^{k+3}u))(\mathbf{z}) + O(h^{k+3}) \\
&= (D_y D_x u)(\mathbf{z}) + h^{k+1}[\mathbf{a}^y \cdot D^{k+2}D_x u + \mathbf{a}^x \cdot D_y(D^{k+2}u)](\mathbf{z}) + \\
&\quad h^{k+2}[\mathbf{b}^y \cdot D^{k+3}D_x u + \mathbf{b}^x \cdot D_y(D^{k+3}u)](\mathbf{z}) + O(h^{k+3}).
\end{aligned} \tag{4.15}$$

Notice that (4.15) is valid only at nodal points. Similarly,

$$\begin{aligned}
(H_h^{yx}u)(\mathbf{z}) &= (D_x D_y u)(\mathbf{z}) + h^{k+1}[\mathbf{a}^x \cdot D^{k+2}D_y u + \mathbf{a}^y \cdot D_x(D^{k+2}u)](\mathbf{z}) + \\
&\quad h^{k+2}[\mathbf{b}^x \cdot D^{k+3}D_y u + \mathbf{b}^y \cdot D_x(D^{k+3}u)](\mathbf{z}) + O(h^{k+3});
\end{aligned} \tag{4.16}$$

$$\begin{aligned}
(H_h^{xx}u)(\mathbf{z}) &= (D_x D_x u)(\mathbf{z}) + h^{k+1}[\mathbf{a}^x \cdot D^{k+2}D_x u + \mathbf{a}^x \cdot D_x(D^{k+2}u)](\mathbf{z}) + \\
&\quad h^{k+2}[\mathbf{b}^x \cdot D^{k+3}D_x u + \mathbf{b}^x \cdot D_x(D^{k+3}u)](\mathbf{z}) + O(h^{k+3});
\end{aligned} \tag{4.17}$$

$$\begin{aligned}
(H_h^{yy}u)(\mathbf{z}) &= (D_y D_y u)(\mathbf{z}) + h^{k+1}[\mathbf{a}^y \cdot D^{k+2}D_y u + \mathbf{a}^y \cdot D_y(D^{k+2}u)](\mathbf{z}) + \\
&\quad h^{k+2}[\mathbf{b}^y \cdot D^{k+3}D_y u + \mathbf{b}^y \cdot D_y(D^{k+3}u)](\mathbf{z}) + O(h^{k+3}).
\end{aligned} \tag{4.18}$$

(4.15)–(4.18) imply that the Hessian recovery operator  $H_h$  is exact for polynomials of degree  $k + 2$  for translation invariant meshes. Also, we observe  $H_h^{xy} = H_h^{yx}$  from (4.15) and (4.16) if  $z$  is a local symmetric center.

It is worth pointing out that, except for the Chevron pattern, (4.15)–(4.18) are valid for

the other four patterns of uniform meshes, since the recovered gradient  $G_h u$  produces the same stencil at each node.

Next we consider even order ( $k = 2r$ ) element on translation invariant meshes and further assume that  $\mathbf{z}$  is a local symmetry center for all sampling points involved, in which case

$$\mathbf{a}^x(\mathbf{z}) = \mathbf{0}, \quad \mathbf{c}^x(\mathbf{z}) = \mathbf{0}, \quad \mathbf{a}^y(\mathbf{z}) = \mathbf{0}, \quad \mathbf{c}^y(\mathbf{z}) = \mathbf{0}; \quad (4.19)$$

$$D\mathbf{a}^x(\mathbf{z}) = \mathbf{0}, \quad D\mathbf{c}^x(\mathbf{z}) = \mathbf{0}, \quad D\mathbf{a}^y(\mathbf{z}) = \mathbf{0}, \quad D\mathbf{c}^y(\mathbf{z}) = \mathbf{0}. \quad (4.20)$$

and  $\mathbf{b}^x, \mathbf{b}^y, \dots$  are constants in (4.14). Here the symbol  $D$  is understood as taking all partial derivatives to each entry of the vector. Consequently,

$$(G_h^y u)(\mathbf{z}) = (D_y u)(\mathbf{z}) + h^{k+2}(\mathbf{b}^y \cdot D^{k+3} u)(\mathbf{z}) + O(h^{k+4}), \quad (4.21)$$

Also, (4.21) is valid only at nodal points. Plugging (4.13) into (4.21) yields

$$\begin{aligned} (H_h^{xy} u)(\mathbf{z}) &= (G_h^y G_h^x u)(\mathbf{z}) \\ &= (D_y G_h^x u)(\mathbf{z}) + h^{k+2}(\mathbf{b}^y \cdot D^{k+3} G_h^x u)(\mathbf{z}) + O(h^{k+4}) \\ &= D_y(D_x u + h^{k+1}\mathbf{a}^x \cdot D^{k+2} u + h^{k+2}\mathbf{b}^x \cdot D^{k+3} u + h^{k+3}\mathbf{c}^x \cdot D^{k+4} u \\ &\quad + \dots)(\mathbf{z}) + h^{k+2}(\mathbf{b}^y \cdot D^{k+3} D_x u)(\mathbf{z}) + O(h^{k+4}) \\ &= (D_y D_x u)(\mathbf{z}) + h^{k+2}(\mathbf{b}^x \cdot D_y D^{k+3} u + \mathbf{b}^y \cdot D^{k+3} D_x u)(\mathbf{z}) + O(h^{k+4}). \end{aligned}$$

In the last identity we have used (4.19) and (4.20).

The argument for the other three entries of recovered Hessian matrix are similar. We

conclude that the Hessian recovery operator  $H_h$  is exact for polynomials of degree up to  $k + 3$  when  $k$  is even and the mesh is translation invariant and symmetric with respect to  $x$  and  $y$ .

The above results can be summarized as the following theorem:

**Theorem 4.1.** *The Hessian recovery operator  $H_h$  preserves polynomials of degree  $k+1$  for an arbitrary mesh. If  $z$  is a node of a translation invariant mesh, then  $H_h$  preserves polynomials of degree  $k + 2$ . If we further suppose  $z$  is local symmetry center for all sampling points involved and  $k$  is a even number, then  $H_h$  preserves polynomials of degree  $k + 3$ . Moreover, if the sampling points are symmetric with respect to  $z$ , then  $H_h$  is symmetric.*

**Remark.** According to [94], the best Hessian recovery method in the literature preserves polynomial of degree 2 for linear element. Our method preserves polynomial of degree 2 on general unstructured meshes and preserves polynomials of degree 3 on translation invariant meshes for linear element.

**Theorem 4.2.** *Let  $u \in W_\infty^{k+2}(\mathcal{K}_z)$ ; then*

$$\|Hu - H_h u\|_{0,\infty,\mathcal{K}_z} \lesssim h^k |u|_{k+2,\infty,\mathcal{K}_z}.$$

*If  $z$  is a node of translation invariant mesh and  $u \in W_\infty^{k+3}(\mathcal{K}_z)$ , then*

$$|(Hu - H_h u)(z)| \lesssim h^{k+1} |u|_{k+3,\infty,\mathcal{K}_z}.$$

*Furthermore, if  $z$  is a symmetric node of translation invariant mesh and  $u \in W_\infty^{k+4}(\mathcal{K}_z)$  with*

$k$  an even number, then

$$|(Hu - H_h u)(z)| \lesssim h^{k+2} |u|_{k+4, \infty, \mathcal{X}_z}.$$

*Proof.* It is a direct result of Theorem 4.1 and application of the Hilbert-Bramble Lemma.  $\square$

## 4.4 Superconvergence analysis

In this section, we first use the supercloseness between the gradient of the finite element solution  $u_h$  and the gradient of the interpolation  $I_h u$  [12, 28, 56, 57, 98, 101], and properties of the PPR operator [107, 79] to establish the superconvergence property of our Hessian recovery operator on mildly structured mesh. Then we utilize the tool of superconvergence by difference quotients from [97] to prove the proposed Hessian recovery method is ultraconvergent for translation invariant finite element space of any order.

### 4.4.1 Linear element

Linear finite element space  $S^{h,1}$  on quasi-uniform mesh  $\mathcal{T}_h$  is considered in this subsection. In order to discuss superconvergent result, we need some condition on the mesh. We firstly talk about mesh condition.

**Definition 4.3.** *The triangulation  $\mathcal{T}_h$  is said to satisfy Condition  $(\sigma, \alpha)$  if there exist a partition  $\mathcal{T}_{h,1} \cup \mathcal{T}_{h,2}$  of  $\mathcal{T}_h$  and positive constants  $\alpha$  and  $\sigma$  such that every two adjacent triangles in  $\mathcal{T}_{h,1}$  form an  $O(h^{1+\alpha})$  parallelogram and*

$$\sum_{T \in \mathcal{T}_{h,2}} |T| = O(h^\sigma).$$

An  $O(h^{1+\alpha})$  parallelogram is a quadrilateral shifted from a parallelogram by  $O(h^{1+\alpha})$ .

For general  $\alpha$  and  $\sigma$ , Xu and Zhang [101] proved the following theorem.

**Theorem 4.4.** *Let  $u$  be the solution of (2.2),  $u_h \in S^{h,1}$  be the finite element solution of (2.8), and  $I_h u \in S^{h,1}$  be the linear interpolation of  $u$ . If the triangulation  $\mathcal{T}_h$  satisfies Condition  $(\sigma, \alpha)$  and  $u \in H^3(\Omega) \cap W_\infty^2(\Omega)$ , then*

$$|u_h - I_h u|_{1,\Omega} \lesssim h^{1+\rho}(|u|_{3,\Omega} + |u|_{2,\infty,\Omega}),$$

where  $\rho = \min(\alpha, \sigma/2, 1/2)$ .

Using the above result, we are able to obtain a convergent result for our Hessian recovery operator.

**Theorem 4.5.** *Under the same condition as Theorem 4.4, we have*

$$\|Hu - H_h u_h\|_{0,\Omega} \leq h^\rho \|u\|_{3,\infty,\Omega}.$$

*Proof.* We decompose  $Hu - H_h u_h$  as  $(Hu - H_h u) + H_h(I_h u - u_h)$ , since  $H_h u = H_h(I_h u)$ .

Using the triangle inequality and the definition of  $H_h$ , we obtain

$$\begin{aligned} \|Hu - H_h u_h\|_{0,\Omega} &\leq \|Hu - H_h u\|_{0,\Omega} + \|H_h(I_h u - u_h)\|_{0,\Omega} \\ &= \|Hu - H_h u\|_{0,\Omega} + \|G_h(G_h(I_h u - u_h))\|_{0,\Omega}. \end{aligned}$$

The first term in the above expression is bounded by  $h|u|_{3,\infty,\Omega}$  according to Theorem 4.2.

Since  $G_h$  is a bounded linear operator [80], it follows that

$$\|H_h(I_h u - u_h)\|_{0,\Omega} \lesssim \|\nabla(G_h(I_h u - u_h))\|_{0,\Omega}$$

Notice that  $G_h(I_h u - u_h)$  is a function in  $S_h$  and hence the inverse estimate [35, 25] can be applied. Thus,

$$\|H_h(I_h u - u_h)\|_{0,\Omega} \lesssim h^{-1} \|G_h(I_h u - u_h)\|_{0,\Omega} \lesssim h^{-1} \|I_h u - u_h\|_{1,\Omega}$$

and hence Theorem 4.4 implies that

$$\|H_h(I_h u - u_h)\|_{0,\Omega} \lesssim h^\rho \|u\|_{3,\infty,\Omega}.$$

Combining the above two estimates completes our proof. □

#### 4.4.2 Quadratic element

We proceed to quadratic finite element space  $S^{h,2}$ . According to [57], a triangulation  $\mathcal{T}_h$  is strongly regular if any two adjacent triangles in  $\mathcal{T}_h$  form an  $O(h^2)$  approximate parallelogram.

Huang and Xu proved the following superconvergence results in [57].

**Theorem 4.6.** *et  $u$  be the solution of (2.2),  $u_h \in S^{h,2}$  be the finite element solution of (2.8), and  $I_h u \in S^{h,2}$  be the quadratic interpolation of  $u$ . If the triangulation  $\mathcal{T}_h$  is uniform or strongly regular and  $u \in H^4(\Omega)$ , then*

$$|u_h - I_h u|_{1,\Omega} \lesssim h^3 |u|_{4,\Omega}.$$

Based on the above theorem, we obtain the following superconvergent result.

**Theorem 4.7.** *Under the same assumption as in Theorem 4.6, we have*

$$\|Hu - H_h u_h\|_{0,\Omega} \leq h^2 \|u\|_{4,\Omega}.$$

*Proof.* The proof is similar to the proof of Theorem 4.5 by using Theorem 4.6 and the inverse estimate. □

**Remark.** Theorem 4.7 can be generalized to mildly structured meshes as in [57].

#### 4.4.3 Translation invariant element of any order

First, we observe that the Hessian recovery operator results in a difference quotient. It is due to the fact that  $G_h$  is a difference quotient [107] and the composition of two difference quotients is still a difference quotient. Let us take linear element on uniform triangular mesh of the regular pattern as an example, see Figure 24(a). The recovered second order derivative at a nodal point  $z$  is

$$\begin{aligned} (H_h^{xx} u_h)(z) = \frac{1}{36h^2} & (-12u_0 + 2u_1 - 4u_2 - 4u_3 + 2u_4 - 4u_5 - 4u_6 + 4u_7 + 4u_8 + u_9 \\ & - 2u_{10} + u_{11} + 4u_{12} + 4u_{13} + 4u_{14} + u_{15} - 2u_{16} + u_{17} + 4u_{18}). \end{aligned}$$



Let  $\phi_j$  be the nodal shape functions. Since  $\phi_z(z') = \delta_{zz'}$ , it follows that

$$\begin{aligned}
& (H_h^{xx} u_h) \phi_0(x, y) \\
&= \frac{1}{36h^2} [-12u_0 \phi_0(x, y) + 2u_1 \phi_1(x + h, y) - 4u_2 \phi_2(x + h, y + h) \\
&\quad - 4u_3 \phi_3(x, y + h) + 2u_4 \phi_4(x - h, y) - 4u_5 \phi_5(x - h, y - h) \\
&\quad - 4u_6 \phi_6(x, y - h) + 4u_7 \phi_7(x + 2h, y) + 4u_8 \phi_8(x + 2h, y + h) \\
&\quad + u_9 \phi_9(x + 2h, y + 2h) - 2u_{10} \phi_{10}(x + h, y + 2h) + u_{11} \phi_{11}(x, y + 2h) \\
&\quad + 4u_{12} \phi_{12}(x - h, y + h) + 4u_{13} \phi_{13}(x - 2h, y) + 4u_{14} \phi_{14}(x - 2h, y - h) \\
&\quad + u_{15} \phi_{15}(x - 2h, y - 2h) - 2u_{16} \phi_{16}(x - h, y - 2h) + u_{17} \phi_{17}(x, y - 2h) \\
&\quad + 4u_{18} \phi_{18}(x + h, y - h)].
\end{aligned}$$

The translations are in the directions of  $\ell_1 = (1, 0)$ ,  $\ell_2 = (0, 1)$ ,  $\ell_3 = (\frac{\sqrt{2}}{2}, \frac{\sqrt{2}}{2})$ ,  $\ell_4 = (\frac{\sqrt{2}}{2}, -\frac{\sqrt{2}}{2})$ ,  $\ell_5 = (\frac{\sqrt{5}}{5}, \frac{2\sqrt{5}}{5})$ , and  $\ell_6 = (\frac{2\sqrt{5}}{5}, \frac{\sqrt{5}}{5})$ . Therefore, we can express the recovered second order derivative as

$$(H_h^{xx} u_h)(z) = \sum_{|\nu| \leq M} \sum_{i=1}^6 C_{\nu, h}^i u_h(z + \nu h \ell_i), \quad (4.22)$$

for some integer  $M$ .

Since  $u$  and  $u_h$  are the solution of variational problem (2.2) and (2.8), respectively. Then for any  $v \in S_0^{h, r}$ , we deduce that

$$\mathcal{B}(u - u_h, v) = 0. \quad (4.23)$$

In particular, (4.23) holds for any  $v \in S_{00}^{h, r}$ .

Let all coefficients in the bilinear form  $\mathcal{B}(\cdot, \cdot)$  be constant. Then the orthogonal property

(4.23) implies

$$\mathcal{B}(T_{\nu\tau}^\ell(u - u_h), v) = \mathcal{B}(u - u_h, T_{-\nu\tau}^\ell v) = \mathcal{B}(u - u_h, (T_{\nu\tau}^\ell)^* v) = 0.$$

for any  $v \in S_{00}^{h,r}$ .

Therefore, Theorem 5.5.2 of [97] (with  $F \equiv 0$ ) implies that

$$\begin{aligned} \|H_h^{xx}(u - u_h)\|_{0,\infty,\Omega_0} &\lesssim \left(\ln \frac{d}{h}\right)^{\bar{r}} \min_{v \in S_h} \|H_h^{xx}u - v\|_{0,\infty,\Omega_1} \\ &\quad + d^{-s-\frac{2}{q}} \|H_h^{xx}(u - u_h)\|_{-s,q,\Omega_1}. \end{aligned} \quad (4.24)$$

Here  $\bar{r} = 1$  for linear element and  $\bar{r} = 0$  for higher order element. Note that  $H_h^{xx}u \in S^{h,r}$  and hence the first term on the right hand side of (4.24) can be estimated by standard approximation theory under the assumption that the finite element space includes piecewise polynomial of degree  $k$ :

$$\min_{v \in S_h} \|H_h^{xx}u - v\|_{0,\infty,\Omega_1} \lesssim h^{k+1} |u|_{k+3,\infty,\Omega_1}, \quad (4.25)$$

provided  $u \in W_\infty^{k+3}(\Omega)$ , see [25, 35]. It remains to attack the second term on the right hand side of (4.24). Note that

$$\|H_h^{xx}(u - u_h)\|_{-s,q,\Omega_1} = \sup_{\phi \in C_0^\infty(\Omega_1), \|\phi\|_{s,q',\Omega_1}=1} (H_h^{xx}(u - u_h), \phi). \quad (4.26)$$

Here  $\frac{1}{q} + \frac{1}{q'} = 1$  and

$$\begin{aligned}
(H_h^{xx}(u - u_h), \phi) &= (u - u_h, (H_h^{xx})^* \phi) \\
&\lesssim \|u - u_h\|_{0,\infty,\Omega_2} \|(H_h^{xx})^* \phi\|_{0,1,\Omega_2} \\
&\lesssim \|u - u_h\|_{0,\infty,\Omega_2},
\end{aligned} \tag{4.27}$$

where we use the fact that  $\|(H_h^{xx})^* \phi\|_{0,1,\Omega_2}$  is bounded uniformly with respect to  $h$  when  $s \geq 1$ . We now once again apply Theorem 5.5.1 from [97] to  $\|u - u_h\|_{0,\infty,\Omega_2}$  with  $\Omega_2 \subset\subset \Omega$  separated by  $d$ , then

$$\begin{aligned}
\|u - u_h\|_{0,\infty,\Omega_2} &\lesssim \left(\ln \frac{d}{h}\right)^{\bar{r}} \min_{v \in S_h} \|u - v\|_{0,\infty,\Omega} \\
&\quad + d^{-s-\frac{2}{q}} \|u - u_h\|_{-s,q,\Omega}.
\end{aligned} \tag{4.28}$$

If the separation parameter  $d = O(1)$ , then we combine (4.24), (4.25) and (5.13) to obtain

$$\|H_h^{xx}(u - u_h)\|_{0,\infty,\Omega_0} \lesssim \left(\ln \frac{1}{h}\right)^{\bar{r}} h^{k+1} \|u\|_{k+3,\infty,\Omega} + \|u - u_h\|_{-s,q,\Omega}. \tag{4.29}$$

Following the same argument, we can establish the same result for  $H_h^{xy}$ ,  $H_h^{yx}$ , and  $H_h^{yy}$ .

Therefore, (4.29) is satisfied by replacing  $H_h^{xx}$  with  $H_h$ :

$$\|H_h(u - u_h)\|_{0,\infty,\Omega_0} \lesssim \left(\ln \frac{1}{h}\right)^{\bar{r}} h^{k+1} \|u\|_{k+3,\infty,\Omega} + \|u - u_h\|_{-s,q,\Omega}. \tag{4.30}$$

Now we are in a perfect position to prove our main result for translation invariant finite element space of any order.

**Theorem 4.8.** *Let all the coefficients in the bilinear operator  $B(\cdot, \cdot)$  be constant; let  $\Omega_0 \subset\subset \Omega_2 \subset\subset \Omega$  be separated by  $d = O(1)$ ; let the finite element space  $S^{h,r}$ , which includes piecewise polynomials of degree  $r$ , be translation invariant in the directions required by the Hessian recovery operator  $H_h$  on  $\Omega_2$ ; and let  $u \in W_\infty^{k+3}(\Omega)$ . Assume that Theorem 5.2.2 from [97] is applicable. Then*

$$\|Hu - H_h u_h\|_{0,\infty,\Omega_0} \lesssim \left(\ln \frac{1}{h}\right)^{\bar{r}} h^{k+1} \|u\|_{k+3,\infty,\Omega} + \|u - u_h\|_{-s,q,\Omega}. \quad (4.31)$$

for some  $s \geq 0$  and  $q \geq 1$ .

*Proof.* We decompose

$$Hu - H_h u_h = (Hu - I_h(Hu)) + (I_h(Hu) - H_h u) + H_h(u - u_h), \quad (4.32)$$

where  $I_h(Hu) \in \prod_{i=1}^2 S^{h,r} \times \prod_{i=1}^2 S^{h,r}$  is the standard Lagrange interpolation of  $Hu$  in the finite element space  $S^{h,r}$ . By the standard approximation theory, we obtain

$$\|Hu - I_h(Hu)\|_{0,\infty,\Omega} \lesssim h^{k+1} |Hu|_{k+1,\infty,\Omega} \lesssim h^{k+1} |u|_{k+3,\infty,\Omega}. \quad (4.33)$$

For the second term, using Theorem 4.2, we have

$$\begin{aligned} \|I_h(Hu) - H_h u_h\|_{0,\infty,\Omega_0} &= \left\| \sum_{z \in \mathcal{N}_h} ((Hu)(z) - (H_h u)(z)) \phi_z \right\|_{0,\infty,\Omega_0} \\ &\lesssim \max_{z \in \mathcal{N}_h \cap \Omega_0} |(Hu)(z) - (H_h u)(z)| \\ &\lesssim h^{k+1} |u|_{k+3,\infty,\Omega}. \end{aligned} \quad (4.34)$$

The last term in (4.32) is bounded by (4.30). The conclusion follows by combining (4.30), (4.33) and (4.34).  $\square$

**Remark.** Theorem 4.8 is a ultraconvergence result under the condition

$$\|u - u_h\|_{-s,q,\Omega} \lesssim h^{k+\sigma}, \quad \sigma > 0.$$

The reader is referred to [85] for negative norm estimates.

## 4.5 Numerical tests

In this section, two numerical examples are provided to illustrate our Hessian recovery method. The first one is designed to demonstrate the polynomial preserving property of the proposed Hessian recovery method. The second one is devoted to a comparison of our method with some existing Hessian recovery methods in the literature on both uniform and unstructured meshes.

In order to evaluate the performance of Hessian recovery methods, we split mesh nodes  $\mathcal{N}_h$  into  $\mathcal{N}_{h,1}$  and  $\mathcal{N}_{h,2}$ , where  $\mathcal{N}_{h,2} = \{z \in \mathcal{N}_h : \text{dist}(z, \partial\Omega) \leq L\}$  denotes the set of nodes near boundary and  $\mathcal{N}_{h,1} = \mathcal{N}_h \setminus \mathcal{N}_{h,2}$  denotes rest interior nodes. Now, we can define

$$\Omega_{h,1} = \bigcup \{\tau \in \mathcal{T}_h : \tau \text{ has all of its vertices in } \mathcal{N}_{h,1}\},$$

and  $\Omega_{h,2} = \Omega \setminus \Omega_{h,1}$ . In the following examples we choose  $L = 0.1$ .

Let  $\tilde{G}_h$  be the weighted average recovery operator. Then we define

$$H_h^{ZZ}u_h = \left( \tilde{G}_h(\tilde{G}_h^x u_h), \tilde{G}_h(\tilde{G}_h^y u_h) \right),$$

and

$$H_h^{LS}u_h = \left( \tilde{G}_h(G_h^{x_1} u_h), \tilde{G}_h(G_h^{x_2} u_h) \right).$$

For any nodal point  $z$ , fit a quadratic polynomial  $p_z$  at  $z$  as PPR. Then  $H_h^{QF}$  is defined as

$$H_h^{QF}u_h(z) = \begin{pmatrix} \frac{\partial^2 p_z}{\partial x_1^2}(0,0) & \frac{\partial^2 p_z}{\partial x_1 \partial x_2}(0,0) \\ \frac{\partial^2 p_z}{\partial x_2 \partial x_1}(0,0) & \frac{\partial^2 p_z}{\partial x_2^2}(0,0) \end{pmatrix}.$$

$H_h^{ZZ}$ ,  $H_h^{LS}$ , and  $H_h^{QF}$  are the first three Hessian recovery methods in [89]. To compare them, define

$$\begin{aligned} De &= \|H_h u_h - Hu\|_{L^2(\Omega_{1,h})}, & De^{ZZ} &= \|H_h^{ZZ}u_h - Hu\|_{L^2(\Omega_{1,h})}, \\ De^{LS} &= \|H_h^{LS}u_h - Hu\|_{L^2(\Omega_{1,h})}, & De^{QF} &= \|H_h^{QF}u_h - Hu\|_{L^2(\Omega_{1,h})}. \end{aligned}$$

where  $u_h$  is the finite element solution.

**Example 4.3.** Consider the following function

$$u(x, y) = \sin(\pi x)\sin(\pi y), \quad (x, y) \in \Omega = (0, 1) \times (0, 1). \quad (4.35)$$

Let  $I_h u$  be the standard Lagrangian interpolation of  $u$  in the finite element space. To validate Theorem 4.2, we apply the Hessian recovery operator  $H_h$  to  $I_h u$  and consider the discrete

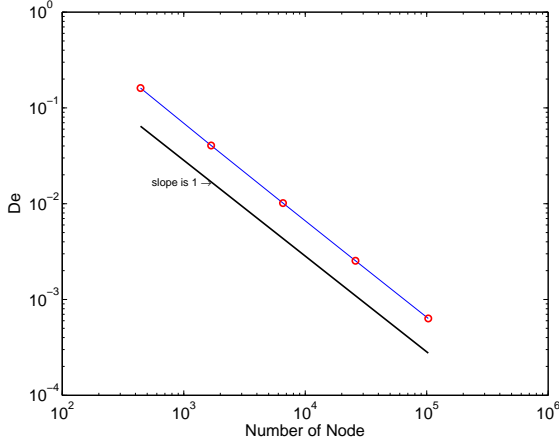
maximum error of  $H_h(I_h u) - Hu$  at all vertices in  $N_{1,h}$ . First, linear element on uniform meshes are taken into account. Figure 25 display the numerical results on the uniform meshes. The numerical errors decrease at a rate of  $O(h^2)$  for four different pattern uniform meshes. It means the proposed Hessian recovery method preserves polynomial of degree 3 for linear element on uniform meshes.

Next, we consider unstructured meshes. We start from an initial mesh generated by EasyMesh[84] as shown in Figure 26(a), followed by four levels of refinement using bisection. Figure 26(b) shows that the recovered Hessian  $H_h(I_h u)$  converges to the exact Hessian at rate  $O(h)$ . This coincides with the result in Theorem 4.1 that  $H_h$  only preserves polynomials of degree 2 on general unstructured meshes

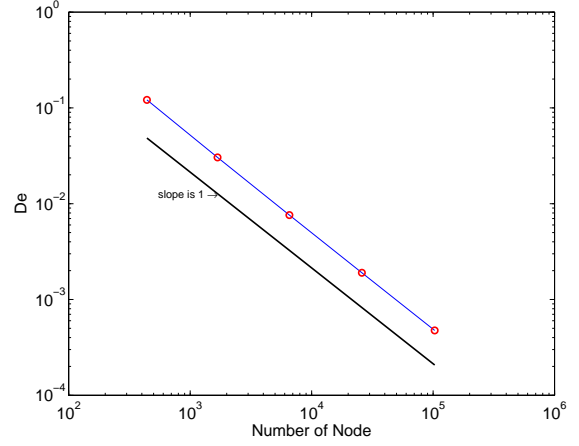
Then we turn to quadratic element. We test the discrete error of recovered Hessian  $H_h(I_h u)$  and the exact Hessian  $Hu$  using uniform meshes of regular pattern and the same Delaunay meshes. Similarly, we define  $\|\cdot\|_{\infty,h}$  as a discrete maximum norm at all vertices and edge centers in an interior region  $\Omega_{1,h}$ . The result of uniform mesh of regular pattern is reported in Figure 27(a). As predicted by Theorem 4.2,  $H_h u_I$  converges to  $Hu$  at rate of  $O(h^4)$  which implies  $H_h$  preserves polynomials of degree 5 for quadratic element on uniform triangulation. For unstructured mesh, we observe that  $H_h u_I$  approximates  $Hu$  at a rate of  $O(h^2)$  from Figure 27(b).

**Example 3.5.2.** We consider the following elliptic equation

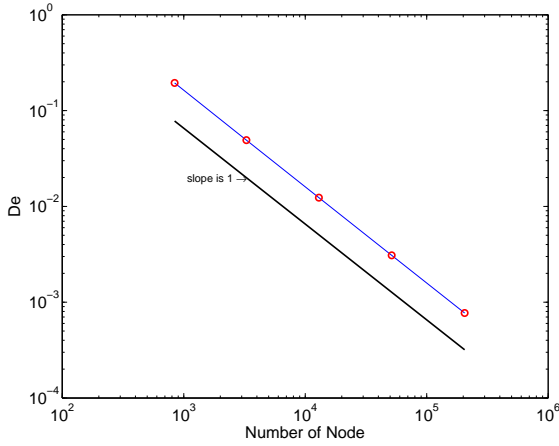
$$\begin{cases} -\Delta u = 2\pi^2 \sin \pi x \sin \pi y, & \text{in } \Omega = [0, 1] \times [0, 1], \\ u = 0, & \text{on } \partial\Omega. \end{cases} \quad (4.36)$$



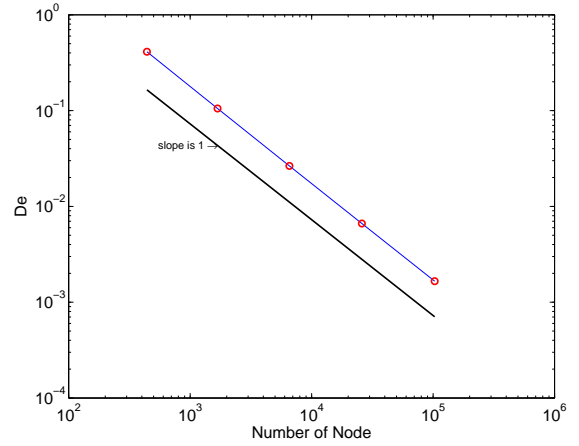
(a) Regular pattern



(b) Chevron pattern



(c) Crisscross pattern



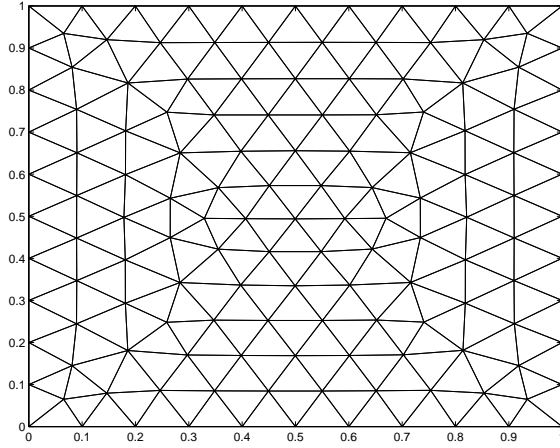
(d) Unionjack pattern

Figure 25: Polynomial preserving property of Hessian recovery for linear element on 2D uniform mesh

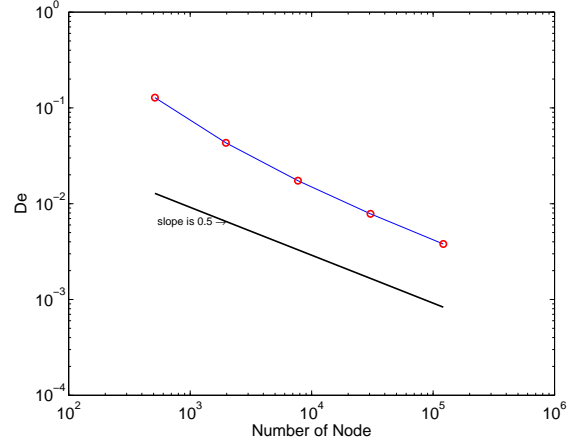
The exact solution is  $u(x, y) = \sin(\pi x) \sin(\pi y)$ . First, linear element is considered. In Table 1, we report the numerical results for regular pattern meshes. All four methods ultraconverge at a rate of  $O(h^2)$  in the interior subdomain. It is not a surprise that  $H_h^{LS}$  and  $H_h^{ZZ}$  perform as good as  $H_h$  since it is well known that the polynomial preserving recovery is the same as weighted average for uniform triangular mesh of the regular pattern.

The results of the Chevron pattern is shown in Table 2.  $H_h u_h$  approximates  $Hu$  at rate  $O(h^2)$  while  $H_h^{LS} u_h$ ,  $H_h^{ZZ} u_h$  and  $H_h^{QF} u_h$  approximate  $Hu$  at rate  $O(h)$ . It is observed that



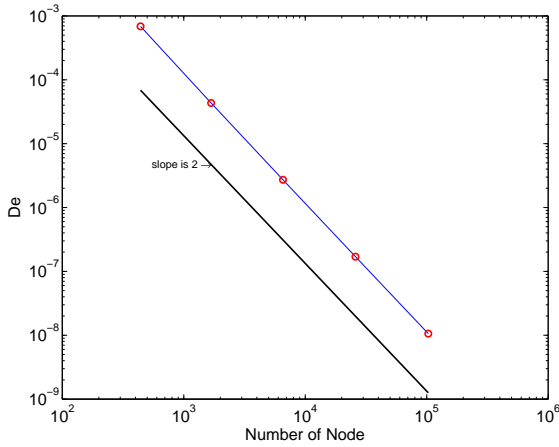


(a) Delaunay mesh

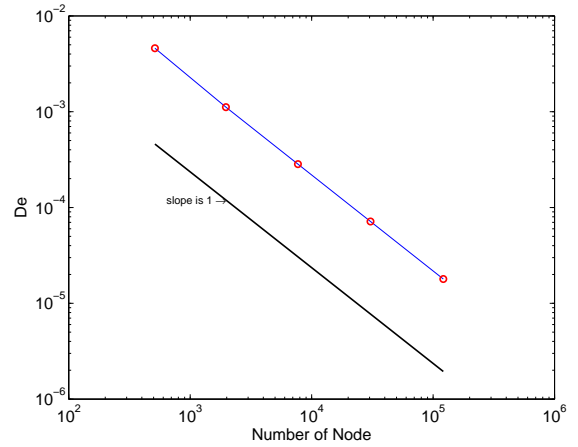


(b) Numerical Result

Figure 26: Polynomial preserving property of Hessian recovery for linear element on 2D unstructured mesh



(a) Uniform mesh



(b) Delaunay mesh

Figure 27: Polynomial preserving property of Hessian recovery for 2D quadratic element our method out-performs other three Hessian recovery methods on the Chevron pattern uniform meshes. To the best of our knowledge, the proposed PPR-PPR Hessian recovery is the only method to achieve  $O(h^2)$  superconvergence for linear element under the Chevron pattern triangular mesh.

Then the Criss-cross pattern mesh is considered and results are displayed in Table 3. An  $O(h^2)$  convergence rate is observed for our recovery method,  $H_h^{LS}$  and  $H_h^{ZZ}$  while no

Table 1: Comparative results for linear element on 2D regular pattern mesh

Dof	$De$	order	$De^{ZZ}e$	order	$De^{LS}$	order	$De^{QF}$	order
121	7.93e-001		9.73e-001		7.93e-001		4.01e-001	
441	2.02e-001	1.06	2.02e-001	1.22	2.02e-001	1.06	1.03e-001	1.05
1681	5.10e-002	1.03	5.10e-002	1.03	5.10e-002	1.03	2.61e-002	1.03
6561	1.28e-002	1.02	1.28e-002	1.02	1.28e-002	1.02	6.53e-003	1.02
25921	3.20e-003	1.01	3.20e-003	1.01	3.20e-003	1.01	1.63e-003	1.01
103041	8.00e-004	1.00	8.00e-004	1.00	8.00e-004	1.00	4.08e-004	1.00

Table 2: Comparative results for linear element on 2D chevron pattern mesh

Dof	$De$	order	$De^{ZZ}e$	order	$De^{LS}$	order	$De^{QF}$	order
121	6.51e-001		7.98e-001		7.82e-001		9.03e-001	
441	1.34e-001	1.22	2.12e-001	1.03	2.34e-001	0.93	4.30e-001	0.57
1681	3.38e-002	1.03	7.96e-002	0.73	9.87e-002	0.64	2.11e-001	0.53
6561	8.46e-003	1.02	3.57e-002	0.59	4.68e-002	0.55	1.05e-001	0.51
25921	2.11e-003	1.01	1.73e-002	0.53	2.30e-002	0.52	5.23e-002	0.51
103041	5.29e-004	1.00	8.57e-003	0.51	1.15e-002	0.50	2.62e-002	0.50

Table 3: Comparative results for linear element on 2D crisscross pattern mesh

Dof	$De$	order	$De^{ZZ}e$	order	$De^{LS}$	order	$De^{QF}$	order
221	5.49e-001		3.57e-001		4.40e-001		7.14e-001	
841	1.28e-001	1.09	8.03e-002	1.12	1.04e-001	1.08	6.17e-001	0.11
3281	3.22e-002	1.01	2.01e-002	1.02	2.62e-002	1.01	5.95e-001	0.03
12961	8.06e-003	1.01	5.04e-003	1.01	6.55e-003	1.01	5.90e-001	0.01
51521	2.02e-003	1.00	1.26e-003	1.00	1.64e-003	1.00	5.89e-001	0.00
205441	5.04e-004	1.00	3.15e-004	1.00	4.09e-004	1.00	5.88e-001	0.00

Table 4: Comparative results for linear element on 2D unionjack pattern mesh

Dof	$De$	order	$De^{ZZ}e$	order	$De^{LS}$	order	$De^{QF}$	order
121	1.25e+000		8.40e-001		9.87e-001		1.05e+000	
441	3.16e-001	1.06	1.77e-001	1.20	2.48e-001	1.07	6.95e-001	0.32
1681	7.96e-002	1.03	4.46e-002	1.03	6.24e-002	1.03	6.14e-001	0.09
6561	2.00e-002	1.02	1.12e-002	1.02	1.56e-002	1.02	5.95e-001	0.02
25921	5.00e-003	1.01	2.80e-003	1.01	3.91e-003	1.01	5.90e-001	0.01
103041	1.25e-003	1.00	6.99e-004	1.00	9.78e-004	1.00	5.89e-001	0.00

Table 5: Comparative results for linear element on 2D unstructured mesh

Dof	$De$	order	$De^{ZZ}e$	order	$De^{LS}$	order	$De^{QF}$	order
139	4.31e-001		4.38e-001		4.40e-001		3.26e-001	
513	1.38e-001	0.87	2.20e-001	0.53	1.49e-001	0.83	1.79e-001	0.46
1969	5.39e-002	0.70	2.36e-001	-0.05	5.85e-002	0.69	8.88e-002	0.52
7713	2.38e-002	0.60	1.62e-001	0.28	2.55e-002	0.61	4.35e-002	0.52
30529	1.14e-002	0.54	1.13e-001	0.26	1.19e-002	0.56	2.15e-002	0.51
121473	5.59e-003	0.51	7.97e-002	0.25	5.73e-003	0.53	1.07e-002	0.51

convergence rate is observed for  $H_h^{QF}$ . The results for the Union-Jack pattern mesh is very similar to the Criss-cross pattern mesh except that our recovery method superconverges at

rate  $O(h^2)$  as shown in Table 4.

Now, we turn to unstructured mesh generated by EasyMesh [84] as in the previous examples. Numerical data are listed in Table 5.  $H_h$ ,  $H_h^{LS}$  and  $H_h^{QF}$  converge at a rate of  $O(h^2)$  while  $H_h^{ZZ}$  only converges at a rate of  $O(h)$ .

The results above indicate clearly that our Hessian recovery method converges at rate  $O(h)$  on general Delaunay meshes, which is predicted by Theorem 4.5. On uniform meshes, we can obtain  $O(h^2)$  ultraconvergence on an interior sub-domain as predicted by Theorem 4.8.

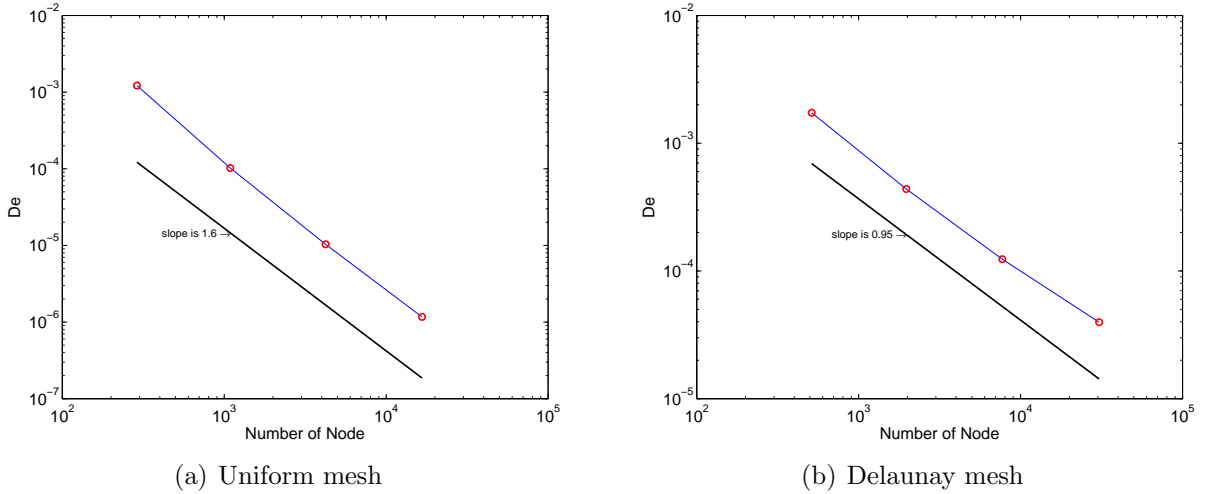


Figure 28: Numerical result of Hessian recovery for 2D quadratic element

Now, we consider quadratic element. Note that our Hessian recovery method is well defined for arbitrary order elements. However, the extension of the other three methods to quadratic element is not straightforward or even impossible and hence only our method is implemented here. We report the numerical results in Figure 28(a) for regular pattern uniform mesh. About  $O(h^{3.2})$  order convergence is observed, which is a bit better than the theoretical result predicted by Theorem 4.8. Figure 28(b) shows the result for Delaunay mesh

generated by EasyMesh [84]. About  $O(h^{1.9})$  superconvergence is observed.

## 4.6 Conclusion

In this chapter, we introduced a Hessian recovery method for arbitrary order Lagrange finite elements. Theoretically, we proved that the PPR-PPR Hessian recovery operator  $H_h$  preserves polynomials of degree  $k + 1$  on general unstructured meshes and preserves polynomials of degree  $k + 2$  on translation invariant meshes. This polynomial preserving property, combined with the supercloseness property of the finite element method, enables us to prove convergence and superconvergence results for our Hessian recovery method on mildly structured meshes. Moreover, we proved the ultraconvergence result for translation invariant finite element space of any order by using the argument of superconvergence by difference quotient from [97].

## CHAPTER 5 Superconvergent two-grid schemes for elliptic eigenvalue problems

A tremendous variety of science and engineering applications, e.g. the buckling of columns and shells and the vibration of elastic bodies, contain models of eigenvalue problems of partial differential equations. A recent survey article [46] of SIAM Review listed 515 references on theory and application of the Laplacian eigenvalue problem.

In this chapter, we apply PPR gradient recovery to efficient computation of eigenvalue. We combine ideas of the two-grid method[102, 34, 62, 72, 108], two-space method[5, 91], shifted-inverse power method[55, 103], and PPR recovery enhancement [82, 78, 81] to design our new algorithms. The first purpose is to introduce two superconvergent two-grid methods for eigenvalue problems. The new proposed methods enjoy all advantages of the above methods : low computational cost and superconvergence.

In addition, we apply PPR gradient recovery for adaptive finite element method of eigenvalue problems. In the context of adaptive finite element method for elliptic eigenvalue problems, residual type a posteriori error estimators are analyzed in [40, 53, 67] and recovery type a posteriori error estimators are investigated by [75, 99, 73]. For all adaptive methods mentioned above, an algebraic eigenvalue problem has to be solved during every iteration, which is very time consuming. This cost dominates the computational cost of AFEM and usually is ignored. To reduce computational cost, Mehrmann and Miedlar [77] introduced a new adaptive method which only requires an inexact solution of algebraic eigenvalue equation on each iteration by only performing a few iterations of Krylov subspace solver. Recently, Li and Yang [68] proposed an adaptive finite element method based on multi-scale discretization

for eigenvalue problems and Xie [100] introduced a type of adaptive finite element method based on the multilevel correction scheme. Both methods only solve an eigenvalue problem on the coarsest mesh and solve boundary value problems on adaptive refined meshes.

The second purpose of this chapter is to propose two multilevel adaptive methods. Using our methods, solving an eigenvalue problem by AFEM will not be more difficult than solving a boundary value problem by AFEM. The most important feature which makes them distinguishing from the methods in [68, 100] is that superconvergence of eigenfunction approximation and ultraconvergence (two order higher) of eigenvalue approximation can be numerically observed.

In Section 5.1, we introduce the model eigenvalue problem and its conforming finite element discretization. Section 5.2 is devoted to presenting two superconvergent two-grid methods and their error estimates. In Section 5.3, we propose two multilevel adaptive methods. Section 5.4 gives some numerical examples to demonstrate efficiency of our new methods and finally some conclusions are drawn in Section 5.5.

## 5.1 A PDE eigenvalue problem and its conforming finite element discretization

Consider the following second order self adjoint elliptic eigenvalue problem:

$$\begin{cases} -\nabla(\mathcal{D}\nabla u) + cu = \lambda u, & \text{in } \Omega, \\ u = 0, & \text{on } \partial\Omega; \end{cases} \quad (5.1)$$

where  $\mathcal{D}$  is a  $2 \times 2$  symmetric positive definite matrix and  $c \in L^\infty(\Omega)$ . Define a bilinear form

$\mathcal{B}(\cdot, \cdot) : H^1(\Omega) \times H^1(\Omega) \rightarrow \mathbb{R}$  by

$$\mathcal{B}(u, v) = \int_{\Omega} (\mathcal{D}\nabla u \cdot \nabla v + cuv) dx.$$

Without loss of generality, we may assume that  $c \geq 0$ . It is easy to see that

$$\mathcal{B}(u, v) \leq C\|u\|_{1,\Omega}\|v\|_{1,\Omega}, \quad \forall u, v \in H_0^1(\Omega),$$

and

$$\mathcal{B}(u, u) \geq \alpha\|u\|_{1,\Omega}^2, \quad \forall u \in H_0^1(\Omega).$$

Define  $\|\cdot\|_{\Omega} = \sqrt{\mathcal{B}(\cdot, \cdot)}$ . Then  $\|\cdot\|_{\Omega}$  and  $\|\cdot\|_{1,\Omega}$  are two equivalent norms in  $H_0^1(\Omega)$ .

The variational formulation of (5.1) reads as: Find  $(\lambda, u) \in \mathbb{R} \times H_0^1(\Omega)$  with  $u \neq 0$  such that

$$\mathcal{B}(u, v) = \lambda(u, v), \quad \forall v \in H_0^1(\Omega). \quad (5.2)$$

It is well known that (5.2) has a countable sequence of real eigenvalues  $0 < \lambda_1 \leq \lambda_2 \leq \lambda_3 \leq \dots \rightarrow \infty$  and corresponding eigenfunctions  $u_1, u_2, u_3, \dots$  which can be assumed to satisfy  $\mathcal{B}(u_i, u_j) = \lambda_i(u_i, u_j) = \delta_{ij}$ . In the sequence  $\{\lambda_j\}$ , the  $\lambda_i$  are repeated according to geometric multiplicity.

The finite element discretization of (5.1) is : Find  $(\lambda_h, u_h) \in \mathbb{R} \times S_0^{h,r}$  with  $u_h \neq 0$  such that

$$\mathcal{B}(u_h, v_h) = \lambda_h(u_h, v_h), \quad \forall v_h \in S_0^{h,r}. \quad (5.3)$$

Similarly, (5.3) has a finite sequence of eigenvalues  $0 < \lambda_{1,h} \leq \lambda_{2,h} \leq \dots \leq \lambda_{n_h,h}$  and corresponding eigenfunctions  $u_{1,h}, u_{2,h}, \dots, u_{n_h,h}$  which can be chosen to satisfy  $\mathcal{B}(u_{i,h}, u_{j,h}) = \lambda_{i,h}(u_{i,h}, u_{j,h}) = \delta_{ij}$  with  $i, j = 1, 2, \dots, n_h$  and  $n_h = \dim S_0^{h,r}$ .

Suppose that the algebraic multiplicity of  $\lambda_i$  is equal to  $q$ , i.e.  $\lambda_i = \lambda_{i+1} = \dots = \lambda_{i+q-1}$ . Let  $M(\lambda_i)$  be the space spanned by all eigenfunctions corresponding to  $\lambda_i$ . Also, let  $M_h(\lambda_h)$  be the direct sum of eigenspaces corresponding to all eigenvalue  $\lambda_{i,h}$  that converges to  $\lambda_i$ .

For the above conforming finite element discretization, the following result has been established by many authors [7, 8, 31, 92, 102, 103].

**Theorem 5.1.** *Suppose  $M(\lambda_i) \subset H_0^1(\Omega) \cap H^{r+1}(\Omega)$ . Let  $\lambda_{i,h}$  and  $\lambda_i$  be the  $i$ th eigenvalue of (5.3) and (5.2), respectively. Then*

$$\lambda_i \leq \lambda_{i,h} \leq \lambda_i + Ch^{2r}. \quad (5.4)$$

For any eigenfunction  $u_{i,h}$  corresponding to  $\lambda_{i,h}$  satisfying  $\|u_{i,h}\|_{a,\Omega} = 1$ , there exists  $u_i \in M(\lambda_i)$  such that

$$\|u_i - u_{i,h}\|_{\Omega} \leq Ch^r. \quad (5.5)$$

Before ending this subsection, we present an important identity [8] of eigenvalue and eigenfunction approximation.

**Lemma 5.2.** *Let  $(\lambda, u)$  be the solution of (5.2). Then for any  $w \in H_0^1(\Omega) \setminus \{0\}$ , there holds*

$$\frac{\mathcal{B}(w, w)}{(w, w)} - \lambda = \frac{\mathcal{B}(w - u, w - u)}{(w, w)} - \lambda \frac{(w - u, w - u)}{(w, w)}. \quad (5.6)$$



This identity will play an important role in our superconvergence analysis.

According to [82], two adjacent triangles (sharing a common edge) form an  $O(h^{1+\alpha})$  ( $\alpha > 0$ ) approximate parallelogram if the lengths of any two opposite edges differ by only  $O(h^{1+\alpha})$ .

**Definition 5.3.** *The triangulation  $\mathcal{T}_h$  is said to satisfy Condition  $\alpha$  if any two adjacent triangles form an  $O(h^{1+\alpha})$  parallelogram.*

Let  $G_h$  be the PPR gradient recovery operator as defined in section 4.2. Using the same methods [107, 93], we can prove the following superconvergence result:

**Theorem 5.4.** *Suppose  $M(\lambda_i) \subset H_0^1(\Omega) \cap W^{3,\infty}(\Omega)$ . Further, suppose  $\mathcal{T}_h$  satisfies Condition  $\alpha$ . Let  $G_h$  be the polynomial preserving recovery operator and  $r = 1$ . Then for any eigenfunction of (5.3) corresponding to  $\lambda_{i,h}$ , there exists an eigenfunction  $u_i \in M(\lambda_i)$  corresponding to  $\lambda_i$  such that*

$$\|\mathcal{D}^{\frac{1}{2}}\nabla u_i - \mathcal{D}^{\frac{1}{2}}G_h u_{i,h}\|_{0,\Omega} \lesssim h^{1+\rho} \|u_i\|_{3,\infty,\Omega}, \quad (5.7)$$

where  $\rho = \min(\alpha, 1)$ .

As pointed out in [82],  $\alpha = \infty$  if  $\mathcal{T}_h$  is generated using regular refinement. Fortunately, the fine grid  $\mathcal{T}_h$  is always a regular refinement of some coarse grid  $\mathcal{T}_H$  for two-grid method. When we introduce two-grid methods in Section 5.2, we only perform gradient recovery on fine grid  $\mathcal{T}_h$ . Thus we assume  $\alpha = \infty$  and hence  $\rho = 1$  in section 5.2.

## 5.2 Superconvergent two-grid methods

In the literature, two-grid methods [102, 103, 55, 108] were proposed to reduce the cost of eigenvalue computations. To further improve the accuracy, two different approaches: gradient recovery enhancement [82, 93, 78, 99] and two-space methods [5, 91] can be used. Individually, those tools are useful in certain circumstances. Combined them properly, we are able to design much effective and superconvergence algorithms, which we shall describe below.

### 5.2.1 Gradient recovery enhanced shifted inverse power two-grid scheme

In this scheme, we first use the shifted inverse power based two-grid scheme [103, 55] and then apply the gradient recovery enhancing technique [82].

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#### Algorithm 1

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1. Solve the eigenvalue problem on a coarse grid  $\mathcal{T}_H$ : Find  $(\lambda_{i,H}, u_{i,H}) \in \mathbb{R} \times S_0^{H,1}$  and  $\|u_{i,H}\|_a = 1$  satisfying

$$\mathcal{B}(u_{i,H}, v_H) = \lambda_{i,H}(u_{i,H}, v_H), \quad \forall v_H \in S_0^{H,1}. \quad (5.8)$$

2. Solve a source problem on the fine grid  $\mathcal{T}_h$ : Find  $u_h^i \in S_0^{h,1}$  such that

$$\mathcal{B}(u_h^i, v_h) - \lambda_{i,H}(u_h^i, v_h) = (u_{i,H}, v_h), \quad \forall v_h \in S_0^{h,1}, \quad (5.9)$$

and set  $u^{i,h} = \frac{u_h^i}{\|u_h^i\|_a}$ .

3. Apply the gradient recovery operator  $G_h$  on  $u^{i,h}$  to get  $G_h u^{i,h}$ .

4. Set

$$\lambda^{i,h} = \frac{\mathcal{B}(u^{i,h}, u^{i,h})}{(u^{i,h}, u^{i,h})} - \frac{\|\mathcal{D}^{\frac{1}{2}} \nabla u^{i,h} - \mathcal{D}^{\frac{1}{2}} G_h u^{i,h}\|_{0,\Omega}^2}{(u^{i,h}, u^{i,h})}. \quad (5.10)$$


---

To prove our main superconvergence result, we need the following Lemma, which was proved in [103, Theorem 4.1].

**Lemma 5.5.** *Suppose that  $M(\lambda_i) \subset H_0^1(\Omega) \cap W^{3,\infty}(\Omega)$ . Let  $(\lambda^{i,h}, u^{i,h})$  be an approximate eigenpair of (5.2) obtained by Algorithm 1 and let  $H$  be properly small. Then*

$$\text{dist}(u^{i,h}, M_h(\lambda_i)) \lesssim H^4 + h^2, \quad (5.11)$$

where  $\text{dist}(u^{i,h}, M_h(\lambda_i)) = \inf_{v \in M_h(\lambda_i)} \|u^{i,h} - v\|_{\Omega}$ .

Based on the above Lemma, we can establish the superconvergence result for eigenfunctions.

**Theorem 5.6.** *Suppose that  $M(\lambda_i) \subset H_0^1(\Omega) \cap W^{3,\infty}(\Omega)$ . Let  $(\lambda^{i,h}, u^{i,h})$  be an approximate eigenpair of (5.2) obtained by Algorithm 1 and let  $H$  be properly small. Then there exists  $u_i \in M(\lambda_i)$  such that*

$$\|\mathcal{D}^{\frac{1}{2}} G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}} \nabla u_i\|_{0,\Omega} \lesssim (H^4 + h^2). \quad (5.12)$$

*Proof.* Let the eigenfunctions  $\{u_{j,h}\}_{j=i}^{i+q-1}$  be an orthonormal basis of  $M_h(\lambda_i)$ . Note that

$$\text{dist}(u^{i,h}, M_h(\lambda_i)) = \|u^{i,h} - \sum_{j=i}^{j=i+q-1} \mathcal{B}(u^{i,h}, u_{j,h}) u_{j,h}\|_{\Omega}.$$

Let  $\tilde{u}_h = \sum_{j=i}^{j=i+q-1} \mathcal{B}(u^{i,h}, u_{j,h}) u_{j,h}$ . According to Theorem 5.4, there exist  $\{\tilde{u}_j\}_{j=i}^{i+q-1} \subset M(\lambda_i)$  such that

$$\|\mathcal{D}^{\frac{1}{2}} G_h u^{j,h} - \mathcal{D}^{\frac{1}{2}} \nabla \tilde{u}_j\|_{0,\Omega} \lesssim h^2. \quad (5.13)$$

Let  $u_i = \sum_{j=i}^{j=i+q-1} \mathcal{B}(u^{i,h}, u_{j,h}) \tilde{u}_j$ ; then  $u_i \in M(\lambda_i)$ . Using (5.13), we can derive that

$$\begin{aligned}
& \|\mathcal{D}^{\frac{1}{2}} G_h \tilde{u}^h - \mathcal{D}^{\frac{1}{2}} \nabla u_i\|_{0,\Omega} \\
&= \left\| \sum_{j=i}^{j=i+q-1} \mathcal{B}(u^{i,h}, u_{j,h}) (\mathcal{D}^{\frac{1}{2}} G_h u_{j,h} - \mathcal{D}^{\frac{1}{2}} \nabla \tilde{u}_j) \right\|_{0,\Omega} \\
&\lesssim \left( \sum_{j=i}^{j=i+q-1} \|(\mathcal{D}^{\frac{1}{2}} G_h u_{j,h} - \mathcal{D}^{\frac{1}{2}} \nabla \tilde{u}_j)\|_{0,\Omega}^2 \right)^{\frac{1}{2}} \\
&\lesssim h^2.
\end{aligned}$$

Thus, we have

$$\begin{aligned}
& \|\mathcal{D}^{\frac{1}{2}} G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}} \nabla u_i\|_{0,\Omega} \\
&\leq \|\mathcal{D}^{\frac{1}{2}} G_h (u^{i,h} - \tilde{u}_h)\|_{0,\Omega} + \|\mathcal{D}^{\frac{1}{2}} G_h \tilde{u}_h - \mathcal{D}^{\frac{1}{2}} \nabla u_i\|_{0,\Omega} \\
&\lesssim \|G_h (u^{i,h} - \tilde{u}_h)\|_{0,\Omega} + h^2 \\
&\lesssim \|\nabla (u^{i,h} - \tilde{u}_h)\|_{0,\Omega} + h^2 \\
&\lesssim \|u^{i,h} - \tilde{u}_h\|_{\Omega} + h^2 \\
&\lesssim (H^4 + h^2) + h^2 \\
&\lesssim H^4 + h^2;
\end{aligned}$$

where we use Lemma 5.5 to bound  $\|u^{i,h} - \tilde{u}_h\|_{a,\Omega}$ .

The following Lemma is needed in the proof of a superconvergence property of our eigenvalue approximation.

**Lemma 5.7.** *Suppose that  $M(\lambda_i) \subset H_0^1(\Omega) \cap W^{3,\infty}(\Omega)$ . Let  $(\lambda^{i,h}, u^{i,h})$  be an approximate*

eigenpair of (5.2) obtained by Algorithm 1 and let  $H$  be properly small. Then

$$\|\mathcal{D}^{\frac{1}{2}}G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}\nabla u^{i,h}\|_{0,\Omega} \lesssim (H^2 + h). \quad (5.14)$$

*Proof.* Let  $\tilde{u}_h$  be defined as in Theorem 5.6. Then we have

$$\begin{aligned} & \|\mathcal{D}^{\frac{1}{2}}G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}\nabla u^{i,h}\|_{0,\Omega} \\ & \leq \|\mathcal{D}^{\frac{1}{2}}G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}G_h \tilde{u}_h\|_{0,\Omega} + \|\mathcal{D}^{\frac{1}{2}}G_h \tilde{u}_h - \mathcal{D}^{\frac{1}{2}}\nabla \tilde{u}_h\|_{0,\Omega} + \|\mathcal{D}^{\frac{1}{2}}\nabla \tilde{u}_h - \mathcal{D}^{\frac{1}{2}}\nabla u^{i,h}\|_{0,\Omega} \\ & \lesssim \|G_h u^{i,h} - G_h \tilde{u}_h\|_{0,\Omega} + \|\mathcal{D}^{\frac{1}{2}}G_h \tilde{u}_h - \mathcal{D}^{\frac{1}{2}}\nabla \tilde{u}_h\|_{0,\Omega} + \|\mathcal{D}^{\frac{1}{2}}\nabla \tilde{u}_h - \mathcal{D}^{\frac{1}{2}}\nabla u^{i,h}\|_{0,\Omega} \\ & \lesssim \|\nabla u_{i,h} - \nabla \tilde{u}_h\|_{0,\Omega} + \|\mathcal{D}^{\frac{1}{2}}G_h \tilde{u}_h - \mathcal{D}^{\frac{1}{2}}\nabla \tilde{u}_h\|_{0,\Omega} \\ & \lesssim \|u_{i,h} - \tilde{u}_h\|_{\Omega} + \|\mathcal{D}^{\frac{1}{2}}G_h \tilde{u}_h - \mathcal{D}^{\frac{1}{2}}\nabla \tilde{u}_h\|_{0,\Omega} \\ & \lesssim (H^4 + h^2) + h \\ & \lesssim (H^2 + h). \end{aligned}$$

Here we use the fact that  $\|\cdot\|_{\Omega}$  and  $\|\cdot\|_{1,\Omega}$  are two equivalent norms on  $H_0^1(\Omega)$ .

Now we are in a perfect position to prove our main superconvergence result for eigenvalue approximation.

**Theorem 5.8.** *Suppose that  $M(\lambda_i) \subset H_0^1(\Omega) \cap W^{3,\infty}(\Omega)$ . Let  $(\lambda^{i,h}, u^{i,h})$  be an approximate eigenpair of (5.2) obtained by Algorithm 1 and let  $H$  be properly small.*

$$|\lambda^{i,h} - \lambda_i| \lesssim H^6 + h^3. \quad (5.15)$$

*Proof.* It follows from (5.6) and (5.10) that

$$\begin{aligned}
& \lambda^{i,h} - \lambda_i \\
&= \frac{\mathcal{B}(u^{i,h}, u^{i,h})}{(u^{i,h}, u^{i,h})} - \frac{\|\mathcal{D}^{\frac{1}{2}}\nabla u^{i,h} - \mathcal{D}^{\frac{1}{2}}G_h u^{i,h}\|_{0,\Omega}^2}{(u^{i,h}, u^{i,h})} - \lambda_i \\
&= \frac{\mathcal{B}(u^{i,h} - u_i, u^{i,h} - u_i)}{(u^{i,h}, u^{i,h})} - \frac{\|\mathcal{D}^{\frac{1}{2}}\nabla u^{i,h} - \mathcal{D}^{\frac{1}{2}}G_h u^{i,h}\|_{0,\Omega}^2}{(u^{i,h}, u^{i,h})} - \frac{\lambda_i(u^{i,h} - u_i, u^{i,h} - u_i)}{(u^{i,h}, u^{i,h})} \\
&= \frac{(\mathcal{D}^{\frac{1}{2}}(u^{i,h} - u_i), \mathcal{D}^{\frac{1}{2}}(u^{i,h} - u_i))}{(u^{i,h}, u^{i,h})} - \frac{\|\mathcal{D}^{\frac{1}{2}}\nabla u^{i,h} - \mathcal{D}^{\frac{1}{2}}G_h u^{i,h}\|_{0,\Omega}^2}{(u^{i,h}, u^{i,h})} + \\
&\quad \frac{(c(u^{i,h} - u_i), u^{i,h} - u_i) - \lambda_i(u^{i,h} - u_i, u^{i,h} - u_i)}{(u^{i,h}, u^{i,h})} \\
&= \frac{\|\mathcal{D}^{\frac{1}{2}}\nabla G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}\nabla u_i\|_{0,\Omega}^2}{(u^{i,h}, u^{i,h})} + \frac{2(\mathcal{D}^{\frac{1}{2}}G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}\nabla u_i, \mathcal{D}^{\frac{1}{2}}\nabla u^{i,h} - \mathcal{D}^{\frac{1}{2}}G_h u^{i,h})}{(u^{i,h}, u^{i,h})} + \\
&\quad \frac{(c(u^{i,h} - u_i), u^{i,h} - u_i) - \lambda_i(u^{i,h} - u_i, u^{i,h} - u_i)}{(u^{i,h}, u^{i,h})}.
\end{aligned}$$

From Theorem 4.1 in [103], we know that  $\|u^{i,h} - u_i\|_{0,\Omega} \lesssim (H^4 + h^2)$  and hence the last term in the above equation is bounded by  $O((H^4 + h^2)^2)$ . Theorem 5.6 implies that the first term is also bounded by  $O((H^4 + h^2)^2)$ . Using the Hölder inequality, we obtain

$$\begin{aligned}
& |(\mathcal{D}^{\frac{1}{2}}G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}\nabla u_i, \mathcal{D}^{\frac{1}{2}}\nabla u^{i,h} - \mathcal{D}^{\frac{1}{2}}G_h u^{i,h})| \\
&\leq \|\mathcal{D}^{\frac{1}{2}}G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}\nabla u_i\|_{0,\Omega} \|\mathcal{D}^{\frac{1}{2}}\nabla u^{i,h} - \mathcal{D}^{\frac{1}{2}}G_h u^{i,h}\|_{0,\Omega} \\
&\lesssim (H^4 + h^2)(H^2 + h) \lesssim H^6 + h^3
\end{aligned} \tag{5.16}$$

and hence

$$|\lambda^{i,h} - \lambda_i| \lesssim H^6 + h^3.$$

This completes our proof.

Taking  $H = O(\sqrt{h})$ , Theorem 5.6 and 5.8 implies that we can get  $O(h^2)$  superconvergence and  $O(h^3)$  superconvergence for eigenfunction and eigenvalue approximation, respectively.

**Remark.** Using the Hölder inequality to estimate (5.16) does not take into account the cancellation in the integral. Similar as [82], numerical experiments show that the actual bound is

$$|(\mathcal{D}^{\frac{1}{2}}G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}\nabla u_i, \mathcal{D}^{\frac{1}{2}}\nabla u^{i,h} - \mathcal{D}^{\frac{1}{2}}G_h u^{i,h})| \lesssim (H^4 + h^2)^2,$$

which says that we have “double”-order gain by applying recovery.

**Remark.** Algorithm 1 is a combination of the shifted inverse power two-grid method [103, 55] and gradient recovery enhancement [82]. It inherits all excellent properties of both methods: low computational cost and superconvergence. We will demonstrate in our numerical tests that Algorithm 1 outperforms shifted inverse power two-grid method in [103, 55].

**Remark.** If we firstly use classical two-grid methods as in [102] and then apply gradient recovery, we can prove  $\|\mathcal{D}^{\frac{1}{2}}G_h u^{i,h} - \mathcal{D}^{\frac{1}{2}}\nabla u_i\|_{0,\Omega} \lesssim (H^2 + h^2)$  and  $|\lambda^{i,h} - \lambda_i| \lesssim H^3 + h^3$ . It means we can only get optimal convergence rate instead of superconvergent convergence rate when  $H = O(\sqrt{h})$ .

### 5.2.2 Higher order space based superconvergent two-grid scheme

Our second scheme can be viewed as a combination of the two-grid scheme proposed by Yang and Bi [103] or Hu and Cheng [55] and the two-space method introduced by Racheva and Andreev [91].

Note that we use linear finite element space  $S_0^{H,1}$  on coarse grid  $\mathcal{T}_H$  and quadratic finite element space  $S_0^{h,2}$  on fine grid  $\mathcal{T}_h$ . Compared with the two-grid scheme [103, 55], the main

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**Algorithm 2**


---

1. Solve an eigenvalue problem on a coarse grid  $\mathcal{T}_H$ : Find  $(\lambda_{i,H}, u_{i,H}) \in \mathbb{R} \times S_0^{H,1}$  and  $\|u_{i,H}\|_a = 1$  satisfying

$$\mathcal{B}(u_{i,H}, v_H) = \lambda_{i,H}(u_{i,H}, v_H), \quad \forall v_H \in S_0^{H,1}. \quad (5.17)$$

2. Solve a source problem on the fine grid  $\mathcal{T}_h$ : Find  $u_h^i \in S_0^{h,2}$  such that

$$\mathcal{B}(u^{i,h}, v_h) - \lambda_{i,H}(u^{i,h}, v_h) = (u_{i,H}, v_h), \quad \forall v_h \in S_0^{h,2}. \quad (5.18)$$

3. Compute the Rayleigh quotient

$$\lambda^{i,h} = \frac{\mathcal{B}(u^{i,h}, u^{i,h})}{(u^{i,h}, u^{i,h})}. \quad (5.19)$$


---

difference is that Algorithm 2 uses linear element on coarse grid  $\mathcal{T}_H$  and quadratic element on fine grid  $\mathcal{T}_h$  while the two-grid uses linear element on both coarse grid  $\mathcal{T}_H$  and  $\mathcal{T}_h$ . Compared with the two-space method [91], the main difference is that Algorithm 2 uses a coarse grid  $\mathcal{T}_H$  and a fine grid  $\mathcal{T}_h$  whereas the two-space method only uses a grid  $\mathcal{T}_h$ . Algorithm 2 shares the advantages of both methods: low computational cost and high accuracy. Thus, we would expect Algorithm 2 performs much better than both methods.

For Algorithm 2, we have the following Theorem:

**Theorem 5.9.** *Suppose that  $M(\lambda_i) \subset H_0^1(\Omega) \cap H^3(\Omega)$ . Let  $(\lambda^{i,h}, u^{i,h})$  be an approximate eigenpair of (5.2) by Algorithm 1 and let  $H$  be properly small. Then there exists  $u_i \in M(\lambda_i)$  such that*

$$\| \|u^{i,h} - u_i\| \|_{\Omega} \lesssim (H^4 + h^2); \quad (5.20)$$

$$\lambda^{i,h} - \lambda_i \lesssim (H^8 + h^4). \quad (5.21)$$



*Proof.* By Theorem 4.1 in [103], we have

$$\|u^{i,h} - u_i\|_{\Omega} \lesssim \eta_a(H)\delta_H^3(\lambda_i) + \delta_h(\lambda_i); \quad (5.22)$$

and

$$\lambda^{i,h} - \lambda_i \lesssim \eta_a^2(H)\delta_H^6(\lambda_i) + \delta_h^2(\lambda_i). \quad (5.23)$$

Since we use linear element on  $\mathcal{T}_H$  and quadratic element on  $\mathcal{T}_h$ , it follows from the interpolation error estimate [25, 35] that

$$\eta_a(H) \lesssim H, \quad \delta_H(\lambda_i) \lesssim H, \quad \delta_h(\lambda_i) \lesssim h^2.$$

Substituting the above three estimate into (5.22) and (5.23), we get (5.20) and (5.21).

Comparing Algorithm 1 and 2, the main difference is that Algorithm 1 solves a source problem on fine grid  $\mathcal{T}_h$  using linear element and hence perform gradient recovery while Algorithm 2 solves a source problem on fine grid  $\mathcal{T}_h$  using quadratic element. Both Algorithm 1 and 2 lead to  $O(h^2)$  superconvergence for eigenfunction approximation and  $O(h^4)$  ultraconvergence for eigenvalue approximation by taking  $H = O(\sqrt{h})$ . The message we would like to deliver here is that polynomial preserving recovery plays a similar role as quadratic element, but with much lower computational cost.

**Remark.** In order to get higher order convergence, we require higher regularity such as  $M(\lambda_i) \subset H_0^1(\Omega) \cap W^{3,\infty}(\Omega)$  for Algorithm 1 and  $M(\lambda_i) \subset H_0^1(\Omega) \cap H^3(\Omega)$  for Algorithm 2, in the proof. However, we can use Algorithm 1 and 2 to get high accuracy approximation even with low regularity.

### 5.3 Multilevel adaptive methods

---

**Algorithm 3** Given a tolerance  $\epsilon > 0$  and a parameter  $0 \leq \theta < 1$ .

---

1. Generate an initial mesh  $\mathcal{T}_{h_0}$ .
2. Solve (5.2) on  $\mathcal{T}_{h_0}$  to get a discrete eigenpair  $(\bar{\lambda}^{h_0}, u^{h_0})$ .
3. Set  $\ell = 0$ .
4. Compute  $\eta(u^{h_\ell}, T)$  and  $\eta(u^{h_\ell}, \Omega)$ , then let

$$\lambda^{h_\ell} = \bar{\lambda}^{h_\ell} - \eta(u^{h_\ell}, \Omega)^2.$$

5. If  $\eta(u^{h_\ell}, \Omega)^2 < \epsilon$ , stop; else go to 6.
6. Choose a minimal subset of elements  $\widehat{\mathcal{T}}_{h_\ell} \subset \mathcal{T}_{h_\ell}$  such that

$$\sum_{T \in \widehat{\mathcal{T}}_{h_\ell}} \eta^2(u_h, T) \geq \theta \eta^2(u_h, \Omega);$$

then refine the elements in  $\widehat{\mathcal{T}}_{h_\ell}$  and necessary elements to get a new conforming mesh  $\mathcal{T}_{h_{\ell+1}}$ .

7. Find  $u \in S_0^{h_{\ell+1}, 1}$  such that

$$\mathcal{B}(u, v) = \lambda_{h_\ell}(u^{h_\ell}, v), \quad v \in S_0^{h_{\ell+1}, 1},$$

and set  $u^{h_{\ell+1}} = \frac{u}{\|u\|_{0, \Omega}}$ . Define

$$\bar{\lambda}^{h_{\ell+1}} = \frac{\mathcal{B}(u^{h_{\ell+1}}, u^{h_{\ell+1}})}{(u^{h_{\ell+1}}, u^{h_{\ell+1}})}. \quad (5.24)$$

8. Let  $\ell = \ell + 1$  and go to 4.
- 

In this section, we incorporate two-grid methods and gradient recovery enhancing technique into the framework of adaptive finite element method and propose two multilevel adaptive methods. Both methods only need to solve an eigenvalue problem on initial mesh and solve an associated boundary value problem on adaptive refined mesh during every iteration.

---

**Algorithm 4** Given a tolerance  $\epsilon > 0$  and a parameter  $0 \leq \theta < 1$ .

---

1. Generate an initial mesh  $\mathcal{T}_{h_0}$ .
2. Solve (5.2) on  $\mathcal{T}_{h_0}$  to get a discrete eigenpair  $(\bar{\lambda}^{h_0}, u^{h_0})$ .
3. Set  $\ell = 0$ .
4. Compute  $\eta(u^{h_\ell}, T)$  and  $\eta(u^{h_\ell}, \Omega)$ , then let

$$\lambda^{h_\ell} = \bar{\lambda}^{h_\ell} - \eta(u^{h_\ell}, \Omega)^2.$$

5. If  $\eta(u^{h_\ell}, \Omega)^2 < \epsilon$ , stop; else go to 6.
6. Choose a minimal subset of elements  $\widehat{\mathcal{T}}_{h_\ell} \subset \mathcal{T}_{h_\ell}$  such that

$$\sum_{T \in \widehat{\mathcal{T}}_{h_\ell}} \eta^2(u_h, T) \geq \theta \eta^2(u_h, \Omega);$$

then refine the elements in  $\widehat{\mathcal{T}}_{h_\ell}$  and necessary elements to get a new conforming mesh  $\mathcal{T}_{h_{\ell+1}}$ .

7. Find  $u \in S_0^{h_{\ell+1}, 1}$  such that

$$\mathcal{B}(u, v) - \lambda_{h_\ell}(u, v) = (u^{h_\ell}, v), \quad v \in S_0^{h_{\ell+1}, 1}, \quad (5.25)$$

and set  $u^{h_{\ell+1}} = \frac{u}{\|u\|_{0, \Omega}}$ . Define

$$\bar{\lambda}^{h_{\ell+1}} = \frac{\mathcal{B}(u^{h_{\ell+1}}, u^{h_{\ell+1}})}{(u^{h_{\ell+1}}, u^{h_{\ell+1}})}.$$

8. Let  $\ell = \ell + 1$  and go to 4.
-

Let  $u_h$  be a finite element solution in  $S^{h,1}$  and  $G_h$  be PPR recovery operator. Define a local a posteriori error estimator on the element  $T$  as

$$\eta(u_h, T) = \|\mathcal{D}^{\frac{1}{2}}G_h u_h - \mathcal{D}^{\frac{1}{2}}\nabla u_h\|_{0,T}, \quad (5.26)$$

and a global error estimator as

$$\eta(u_h, \Omega) = \left( \sum_{T \in \mathcal{T}_h} \eta(u_h, T)^2 \right)^{\frac{1}{2}}. \quad (5.27)$$

Given a tolerance  $\epsilon$  and a parameter  $\theta$ , we describe our multilevel adaptive methods in Algorithm 3 and 4. Here we use Dörfler marking strategy [38] in step 6.

Note that the only difference between Algorithm 3 and 4 is that they solve different boundary value problems on step 7. Algorithm 3 solves boundary value problem (5.24) like two-grid scheme in [102] while Algorithm 4 solves boundary value problem (5.25) similar to two-grid scheme in [103, 55]. Boundary value problem (5.25) would lead to a near singular linear system. Although there are many efficient iterative methods, like multigrid methods, as pointed out in [55], the computational cost of solving (5.24) should be higher than (5.25). Numerical results of both methods are almost the same as indicated by examples in next section. Thus, Algorithm 3 is highly recommended.

Compared to methods in [68, 100], Algorithm 3 and 4 use recovery based a posteriori error estimator. The propose of gradient recovery in the above two algorithms is twofold. The first one is to provide an asymptotically exact a posteriori error estimator. The other is to greatly improve the accuracy of eigenvalue and eigenfunction approximations. Superconvergence

result  $O(N^{-1})$  and ultraconvergence  $O(N^{-2})$  are numerically observed for eigenfunction and eigenvalue approximation respectively. However, methods in [68, 100] can only numerically give asymptotically optimal results. We would like to emphasize that the new algorithms can get superconvergence or ultraconvergence results with no more or even less computational cost compared to the methods proposed in [68, 100].

## 5.4 Numerical Experiment

In this section, we present several numerical examples to demonstrate the effectiveness and superconvergence of the proposed algorithms and validity our theoretical results.

The first example is designed to demonstrate superconvergence property of Algorithm 1 and 2 and make some comparison with the two-grid scheme in [103, 55]. Let the  $i$ th eigenpairs obtained by Algorithm 1 and 2 be denoted by  $(\lambda^{i,A1}, u^{i,A1})$  and  $(\lambda^{i,A2}, u^{i,A2})$ . Also, let  $(\lambda^{i,TG}, u^{i,TG})$  be the  $i$ th eigenpair produced by the shift inverse based two-grid scheme in [103, 55].

The presentation of other examples is used to illustrate the effectiveness and superconvergence of Algorithm 3 and 4. In these examples, we focus on the first eigenpair. Let  $\bar{\lambda}_{A3}$  and  $\lambda_{A3}$  be the eigenvalue generated by Algorithm 3 without and with gradient recovery enhancing, respectively. Define  $\bar{\lambda}_{A4}$ ,  $\lambda_{A4}$ ,  $u_{A3}$ , and  $u_{A4}$  in a similar way.

**Example 5.1.** Consider the following Laplace eigenvalue problem

$$\begin{cases} -\Delta u = \lambda u, & \text{in } \Omega, \\ u = 0, & \text{on } \partial\Omega, \end{cases} \quad (5.28)$$

where  $\Omega = (0, 1) \times (0, 1)$ . The eigenvalue of (5.28) are  $\lambda_{k,l} = (k^2 + l^2)\pi^2$  and the corresponding eigenfunctions are  $u_{k,l} = \sin(k\pi x) \sin(l\pi y)$  with  $k, l = 1, 2, \dots$ . It is easy to see the first three eigenvalues are  $\lambda_1 = 2\pi^2$  and  $\lambda_2 = \lambda_3 = 5\pi^2$ .

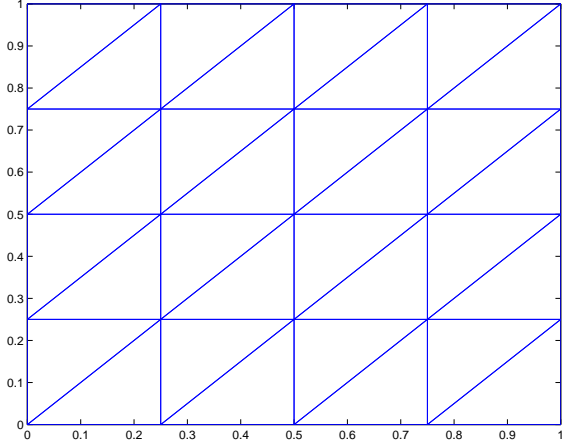


Figure 29: Uniform Mesh

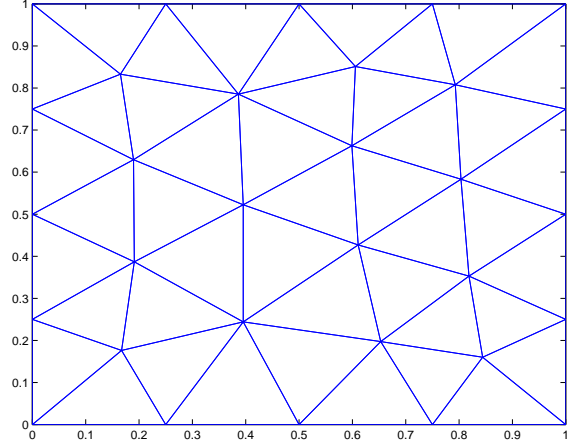


Figure 30: Delaunay Mesh

Table 6: Eigenpair errors of Algorithm 1 for Example 1 on Uniform Mesh

i	H	h	$\lambda^{i,A1}$	$\lambda^{i,A1} - \lambda_i$	Order	$\ G_h u^{i,A1} - \nabla u_i\ _{0,\Omega}$	Order
1	1/4	1/16	19.733813512912	-5.40e-03		7.059395e-02	
1	1/8	1/64	19.739186935311	-2.19e-05	3.97	4.387700e-03	2.00
1	1/16	1/256	19.739208716241	-8.59e-08	4.00	2.734342e-04	2.00
1	1/32	1/1024	19.739208801843	-3.36e-10	4.00	1.707544e-05	2.00
2	1/4	1/16	49.311524605286	-3.65e-02	0.00	————	
2	1/8	1/64	49.347897768530	-1.24e-04	4.10	————	
2	1/16	1/256	49.348021565420	-4.40e-07	4.07	————	
2	1/32	1/1024	49.348022003783	-1.66e-09	4.02	————	
3	1/4	1/16	49.311750580349	-3.63e-02	0.00	————	
3	1/8	1/64	49.347802761238	-2.19e-04	3.69	————	
3	1/16	1/256	49.348021182216	-8.23e-07	4.03	————	
3	1/32	1/1024	49.348022002296	-3.15e-09	4.01	————	

First, uniform mesh as in Fig 29 is considered. The fine meshes  $\mathcal{T}_h$  are of sizes  $h = 2^{-j}$  ( $j = 4, 6, 8, 10$ ) and the corresponding coarse meshes  $\mathcal{T}_H$  of size  $H = \sqrt{h}$ . Table 6 lists the numerical results for Algorithm 1.  $\|G_h u^{i,A1} - \nabla u_i\|_{0,\Omega}$  ( $i = 1$ ) superconverges at rate of  $O(h^2)$  which consists with our theoretical analysis. However,  $|\lambda^{i,A1} - \lambda_i|$  ( $i = 1, 2, 3$ ) ultraconverges at rate of  $O(h^4)$  which is better than the results predicted by Theorem 5.8. In particular, it verifies the statement in Remark 5.2.1. Since  $\lambda_2$  and  $\lambda_3$  are multiples eigenvalues, the error of

Table 7: Eigenpair errors of Algorithm 2 for Example 1 on Uniform Mesh

i	H	h	$\lambda^{i,A2}$	$\lambda^{i,A2} - \lambda_i$	Order	$\ \nabla u^{i,A2} - \nabla u_i\ _{0,\Omega}$	Order
1	1/4	1/16	19.740140941323	9.32e-04		3.344371e-02	
1	1/8	1/64	19.739212357340	3.56e-06	4.02	2.076378e-03	2.00
1	1/16	1/256	19.739208816236	1.41e-08	3.99	1.308168e-04	1.99
1	1/32	1/1024	19.739208802235	5.59e-11	3.99	8.198527e-06	2.00
2	1/4	1/16	49.399143348018	5.11e-02	0.00	-----	
2	1/8	1/64	49.348217238157	1.95e-04	4.02	-----	
2	1/16	1/256	49.348022827362	8.22e-07	3.95	-----	
2	1/32	1/1024	49.348022008741	3.29e-09	3.98	-----	
3	1/4	1/16	49.573605264596	2.26e-01	0.00	-----	
3	1/8	1/64	49.348559514553	5.38e-04	4.36	-----	
3	1/16	1/256	49.348024046492	2.04e-06	4.02	-----	
3	1/32	1/1024	49.348022013418	7.97e-09	4.00	-----	

Table 8: Eigenpair errors of shift-inverse Two-grid scheme for Example 1 on Uniform Mesh

i	H	h	$\lambda^{i,TG}$	$\lambda^{i,TG} - \lambda_i$	Order	$\ \nabla u^{i,TG} - \nabla u_i\ _{0,\Omega}$	Order
1	1/4	1/16	19.930259632276	1.91e-01		4.375101e-01	
1	1/8	1/64	19.751103117985	1.19e-02	2.00	1.090672e-01	1.00
1	1/16	1/256	19.739951989101	7.43e-04	2.00	2.726155e-02	1.00
1	1/32	1/1024	19.739255250511	4.64e-05	2.00	6.815303e-03	1.00
2	1/4	1/16	50.199210624678	8.51e-01	0.00	-----	
2	1/8	1/64	49.399315353599	5.13e-02	2.03	-----	
2	1/16	1/256	49.351217793553	3.20e-03	2.00	-----	
2	1/32	1/1024	49.348221696982	2.00e-04	2.00	-----	
3	1/4	1/16	50.779973345337	1.43e+00	0.00	-----	
3	1/8	1/64	49.428220994371	8.02e-02	2.08	-----	
3	1/16	1/256	49.353003975409	4.98e-03	2.00	-----	
3	1/32	1/1024	49.348333256327	3.11e-04	2.00	-----	

eigenfunctions approximation are not available and it is represented by “—” in Tables 6-12.

One important thing we want to point out is that we observe numerically that  $\lambda_{A1}$  obtained by Algorithm 1 approximates the exact eigenvalue from below; see column 4 in Table 6.

Similar phenomenon was observed in [44] where they use a local high-order interpolation recovery. We want to remark that lower bound of eigenvalue is very important in practice and

there are many efforts for obtaining eigenvalue approximation from below. The readers are

referred to [6, 70, 104, 105] for other ways to approximate eigenvalue from below. In Table 7,

we report the numerical result of Algorithm 2. As expected,  $O(h^4)$  convergence of eigenvalue

approximation and  $O(h^2)$  convergence of eigenfunction approximation are observed which

validate our Theorem 5.9. The shift-inverse power method based two-grid scheme in [103, 55]

is then considered, the result being displayed in Table 8.  $\lambda^{i,\text{TG}}$  approximates  $\lambda_i$  ( $i = 1, 2, 3$ ) at a rate  $O(h^2)$  and  $\|u^{i,\text{TG}} - u_i\|_{a,\Omega}$  ( $i=1$ ) converges at a rate of  $O(h)$ .

Comparing Tables 6 to 8, huge advantages of Algorithm 1 and 2 are demonstrated. For instance, on the fine grid with size  $h = 1/1024$  and corresponding coarse grid with size  $H = 1/32$ , the approximate first eigenvalues produced by Algorithm 1 and 2 are exact up to 10 digits while one can only trust the first five digits of the first eigenvalue generated by the two-grid scheme in [103, 55].

Table 9: Comparison of Three Algorithms for Example 1 on Uniform mesh

i	H	h	$i, \lambda^{A1}$	$\lambda^{i,A1} - \lambda_i$	$\lambda^{i,A2}$	$\lambda^{i,A2} - \lambda_i$	$\lambda^{i,\text{TG}}$	$\lambda^{i,\text{TG}} - \lambda_i$
1	1/2	1/16	20.1083669	3.69e-01	20.2080796	4.69e-01	20.3504780	6.11e-01
1	1/4	1/256	19.7398503	6.41e-04	19.7398588	6.50e-04	19.7406011	1.39e-03

Then we consider the case  $H = O(\sqrt[4]{h})$  for the first eigenvalue. We use the fine meshes of mesh size  $h = 2^{-j}$  with  $j = 4, 8$  and corresponding coarse meshes satisfying  $H = \sqrt[4]{h}$ . The numerical results are shown in Table 9. We can see that the two proposed Algorithms give better approximate eigenvalues. Thus Algorithm 1 and 2 outperforms the two-grid scheme even in the case  $H = \sqrt[4]{h}$ . One interesting thing that we would like to mention is that  $\lambda^{i,A1}$  approximates  $\lambda_i$  from above in this case, see column 4 in Table 9.

Table 10: Eigenpair errors of Algorithm 1 for Example 1 on Delaunay Mesh

i	H	h	$\lambda^{i,A1}$	$\lambda^{i,A1} - \lambda_i$	Order	$\ G_h u^{i,A1} - \nabla u_i\ _{0,\Omega}$	Order
1	31	385	19.735647110619	-3.56e-03		5.338236e-02	
1	105	5761	19.739198229599	-1.06e-05	2.15	2.835582e-03	1.08
1	385	90625	19.739208765246	-3.69e-08	2.05	1.686396e-04	1.02
1	1473	1443841	19.739208802041	-1.38e-10	2.02	1.049196e-05	1.00
2	31	385	49.307472112236	-4.05e-02	0.00	—————	
2	105	5761	49.347888708818	-1.33e-04	2.11	—————	
2	385	90625	49.348021524994	-4.80e-07	2.04	—————	
2	1473	1443841	49.348022003630	-1.82e-09	2.01	—————	
3	31	385	49.301142920140	-4.69e-02	0.00	—————	
3	105	5761	49.347856273486	-1.66e-04	2.09	—————	
3	385	90625	49.348021393237	-6.12e-07	2.03	—————	
3	1473	1443841	49.348022003123	-2.32e-09	2.01	—————	



Table 11: Eigenpair errors of Algorithm 2 for Example 1 on Delaunay Mesh

i	H	h	$\lambda^{i,A2}$	$\lambda^{i,A2} - \lambda_i$	Order	$\ \nabla u^{i,A2} - \nabla u_i\ _{0,\Omega}$	Order
1	31	385	19.739293668773	8.49e-05		9.258930e-03	
1	105	5761	19.739209125443	3.23e-07	2.06	5.705799e-04	1.03
1	385	90625	19.739208803434	1.26e-09	2.01	3.555028e-05	1.01
1	1473	1443841	19.739208802184	5.33e-12	1.97	2.220103e-06	1.00
2	31	385	49.350648806465	2.63e-03	0.00	—————	
2	105	5761	49.348029138391	7.13e-06	2.18	—————	
2	385	90625	49.348022031328	2.59e-08	2.04	—————	
2	1473	1443841	49.348022005547	1.00e-10	2.01	—————	
3	31	385	49.351570779092	3.55e-03	0.00	—————	
3	105	5761	49.348029733509	7.73e-06	2.27	—————	
3	385	90625	49.348022033250	2.78e-08	2.04	—————	
3	1473	1443841	49.348022005554	1.07e-10	2.01	—————	

Table 12: Eigenpair errors of shift-inverse Two-grid scheme for Example 1 on Delaunay Mesh

i	H	h	$\lambda^{i,TG}$	$\lambda^{i,TG} - \lambda_i$	Order	$\ \nabla u^{i,TG} - \nabla u_i\ _{0,\Omega}$	Order
1	31	385	19.821235920927	8.20e-02		2.865766e-01	
1	105	5761	19.744334806708	5.13e-03	1.02	7.159881e-02	0.51
1	385	90625	19.739529185236	3.20e-04	1.01	1.789929e-02	0.50
1	1473	1443841	19.739228826191	2.00e-05	1.00	4.474820e-03	0.50
2	31	385	49.828430094852	4.80e-01	0.00	—————	
2	105	5761	49.377951127988	2.99e-02	1.03	—————	
2	385	90625	49.349892261888	1.87e-03	1.01	—————	
2	1473	1443841	49.348138895061	1.17e-04	1.00	—————	
3	31	385	49.893495693695	5.45e-01	0.00	—————	
3	105	5761	49.381970792689	3.39e-02	1.03	—————	
3	385	90625	49.350143791388	2.12e-03	1.01	—————	
3	1473	1443841	49.348154618353	1.33e-04	1.00	—————	

Now, we turn to unstructured meshes. First we generate a coarse mesh  $\mathcal{T}_H$  and repeat regular refinement on  $\mathcal{T}_H$  until  $H = O(\sqrt{h})$  to get the corresponding fine mesh  $\mathcal{T}_h$ . The first level coarse mesh is generated by EasyMesh [84] and the other three level coarse mesh are generated by regular refinement. The numerical results are provided in Tables 10 to 12. Note that  $N_H$  and  $N_h$  denote the number of vertices on coarse mesh  $\mathcal{T}_H$  and fine mesh  $\mathcal{T}_h$ , respectively. Concerning the convergence of eigenvalue, Algorithm 1 and 2 ultraconverge at rate  $O(h^4)$  while the two-grid scheme converges at rate  $O(h^2)$ . Note that in Tables 5.5–5.7,  $N_H \approx H^{-2}$  and  $N_h \approx h^{-2}$ . Therefore, convergent rates for  $H$  and  $h$  “double” the rates for  $N_H$  and  $N_h$ , respectively. As for eigenfunction,  $\|G_h u^{i,A1} - \nabla u_i\|_{0,\Omega}$  and  $\|\nabla u^{i,A2} - \nabla u_i\|_{0,\Omega}$  are about  $O(h^2)$  while  $\|\nabla u^{i,TG} - \nabla u_i\|_{0,\Omega} \approx O(h)$ .

**Example 5.2.** In the previous example, the eigenfunctions  $u$  are analytic. Here we consider Laplace eigenvalue value problem on the L-shaped domain  $\Omega = (-1, 1) \times (-1, 1) / [0, 1) \times (-1, 0]$ . The first eigenfunction has a singularity at the origin. To capture this singularity, multilevel adaptive algorithms 3 and 4 are used with  $\theta = 0.4$ . Since the first exact eigenvalue is not available, we choose an approximation  $\lambda = 9.6397238440219$  obtained by Betcke and Trefethen in [14], which is correct up to 14 digits.

Fig 31 shows the initial uniform mesh while Fig 32 is the mesh after 18 adaptive iterations. Fig 33 reports numerical results of the first eigenvalue approximation. It indicates clearly  $\bar{\lambda}_{A3}$  and  $\bar{\lambda}_{A4}$  approximate  $\lambda$  at a rate of  $O(N^{-1})$  while  $\lambda_{A3}$  and  $\lambda_{A4}$  approximate  $\lambda$  at a rate of  $O(N^{-2})$ . The numerical results for Algorithms 3 and 4 are almost the same. Furthermore, we notice that  $\lambda_{A3}$  and  $\lambda_{A4}$  approximate the exact eigenvalue from below. It is well known that  $\bar{\lambda}_{A3}$  and  $\bar{\lambda}_{A4}$  are upper bounds for the exact eigenvalue. In actual computation, we use  $\bar{\lambda}_{A3} - \lambda_{A3} \leq \epsilon$  as stop criteria for adaptive Algorithm 3 where  $\epsilon$  is the given tolerance. A similar procedure is applied to Algorithm 4.

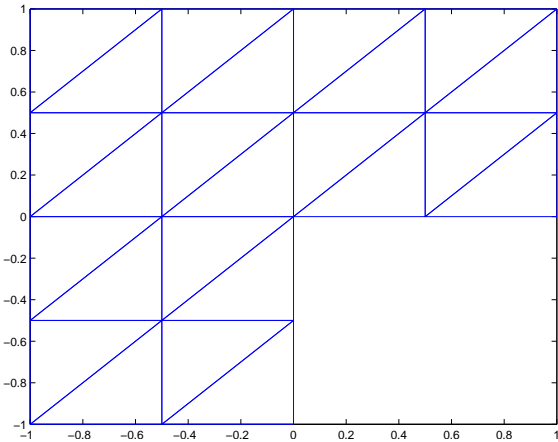


Figure 31: Initial Mesh for Example 5.2

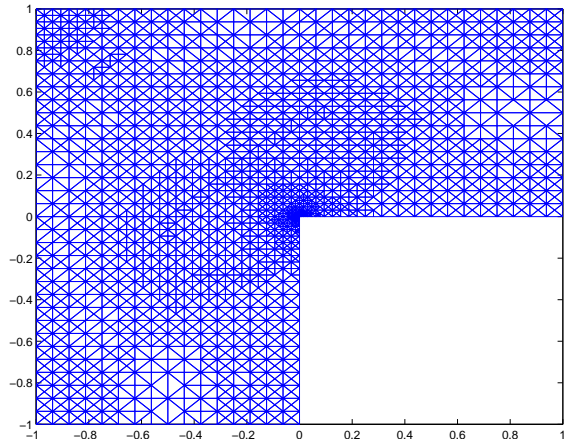


Figure 32: Adaptive Mesh for Example 5.2

In the context of adaptive finite element method for boundary value problems, the ef-

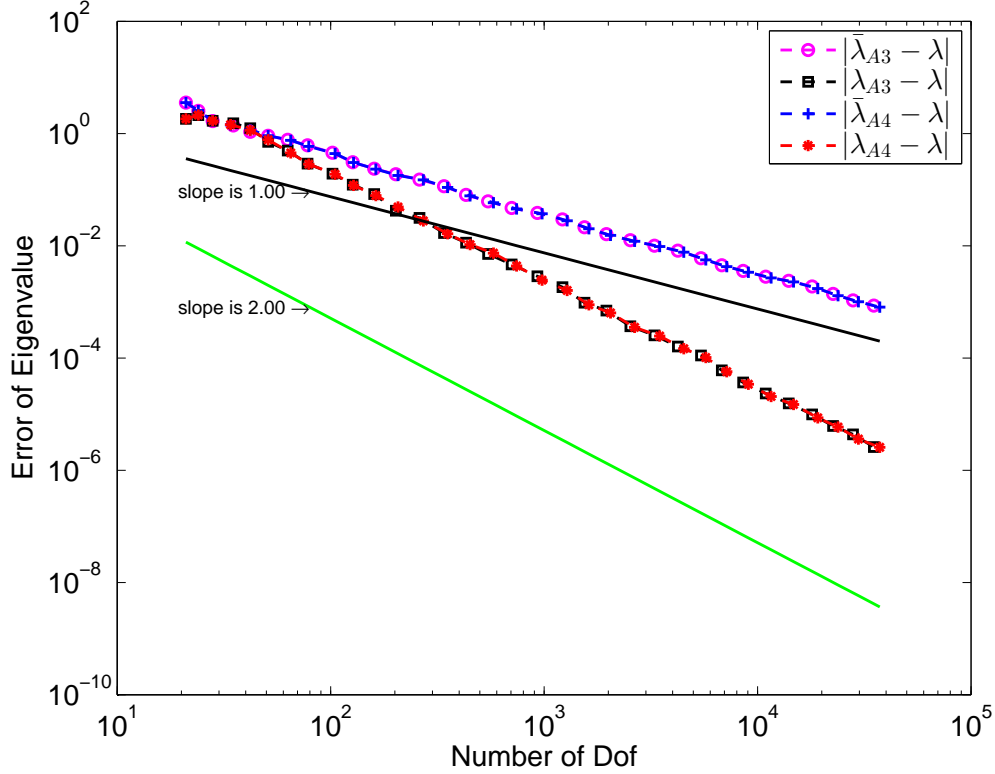


Figure 33: Eigenvalue Approximation Error for Example 5.2

fectivity index  $\kappa$  is used to measure the quality of an error estimator [4, 11]. For eigenvalue problem, it is better to consider eigenvalue effectivity index instead of traditional effectivity index in [4, 11]. In the article, we consider a similar eigenvalue effective index as in [48]

$$\kappa = \frac{\|\mathcal{D}^{\frac{1}{2}}G_h u_h - \mathcal{D}^{\frac{1}{2}}\nabla u_h\|_{0,\Omega}^2}{|\lambda - \lambda_h|}, \quad (5.29)$$

where  $u_h$  is either  $u_{A3}$  or  $u_{A4}$  and  $\lambda_h$  is either  $\lambda_{A3}$  or  $\lambda_{A4}$ . The effectivity index for the two proposed multilevel adaptive algorithms are reported in Figs 34 and 35. We see that  $\kappa$  converges to 1 quickly after the first few iterations, which indicates that the posteriori error estimator (5.26) or (5.27) is asymptotically exact.

**Example 5.3.** Consider the following harmonic oscillator equation [47], which is a simple

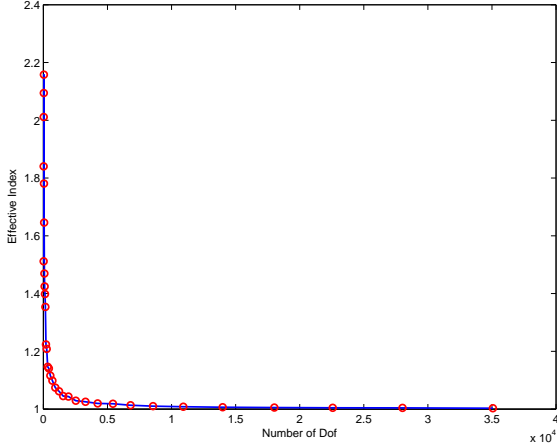


Figure 34: Effective index of Algorithm 3 for Example 5.2

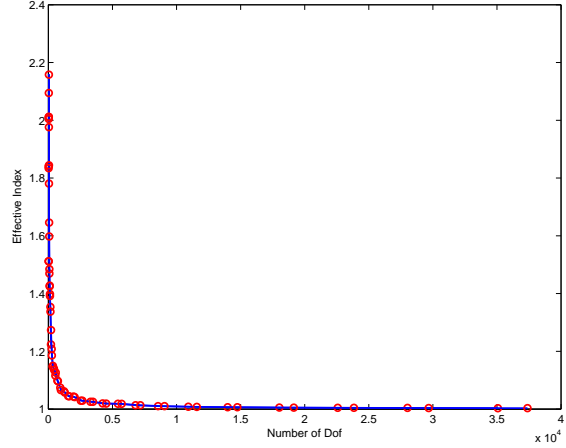


Figure 35: Effective index of Algorithm 4 for Example 5.2

model in quantum mechanics,

$$-\frac{1}{2}\Delta u + \frac{1}{2}|x|^2 u = \lambda u, \quad \text{in } \mathbb{R}^2, \quad (5.30)$$

where  $|x| = \sqrt{|x_1|^2 + |x_2|^2}$ . The first eigenvalue of (5.30) is  $\lambda = 1$  and the corresponding eigenfunction is  $u = \gamma e^{-|x|^2/2}$  with any nonzero constant  $\gamma$ .

We solve this eigenvalue problem with  $\Omega = (-5, 5) \times (-5, 5)$  and zero boundary condition as in [100]. The initial mesh is shown in Fig 36 and the adaptive mesh after 20 iterations is displayed in Fig 37. The parameter  $\theta$  is chosen as 0.4. Numerical results are presented in Figs 38 and 39. For eigenvalue approximation,  $O(N^{-1})$  convergence rate is observed for  $|\bar{\lambda}_{A3} - \lambda|$  while  $O(N^{-2})$  ultraconvergence rate is observed for  $|\lambda_{A3} - \lambda|$ . For eigenfunction approximation,  $\|\mathcal{D}^{\frac{1}{2}}\nabla u_{A3} - \mathcal{D}^{\frac{1}{2}}\nabla u\|_{0,\Omega} \approx O(N^{-0.5})$  and  $\|\mathcal{D}^{\frac{1}{2}}G_h u_{A3} - \mathcal{D}^{\frac{1}{2}}\nabla u\|_{0,\Omega} \approx O(N^{-1})$ . The numerical result of Algorithm 4 is similar.

Figs 40 and 41 graph the eigenvalue effectivity index for the two proposed multilevel adaptive algorithms. They also indicate that the posteriori error estimator (5.26) or (5.27)

is asymptotically exact for problem (5.30).

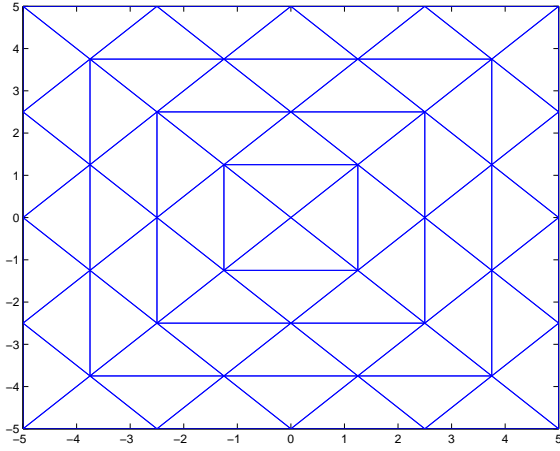


Figure 36: Initial Mesh for Example 5.3

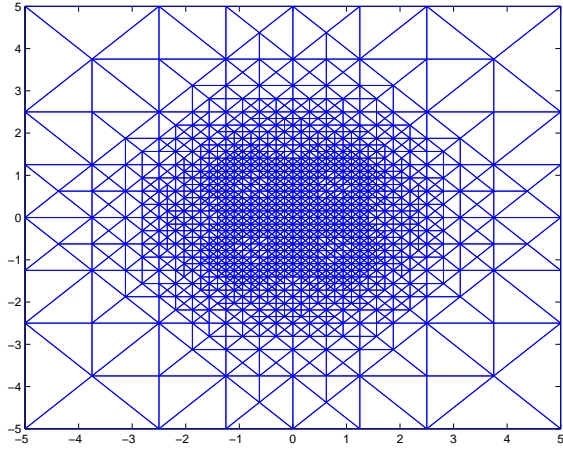


Figure 37: Adaptive Mesh for Example 5.3

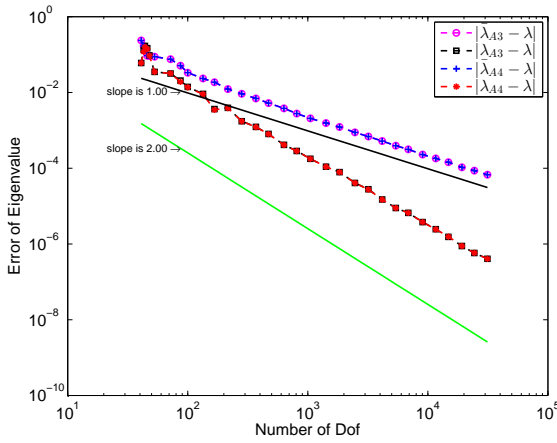


Figure 38: Eigenvalue approximation Error for Example 5.3

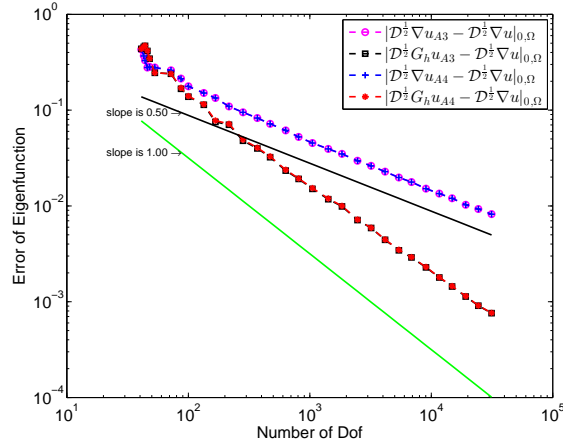


Figure 39: Eigenfunction approximation Error for Example 5.3

## 5.5 Conclusion

When eigenfunctions are relatively smooth, two-space method (using higher-order elements in the second stage) is superior to two-grid methods (using the same element at finer grids in the second stage). They have the comparable accuracy. However, at the last stage, the degrees of freedom of the two-space method is much smaller than that of the two-grid method.

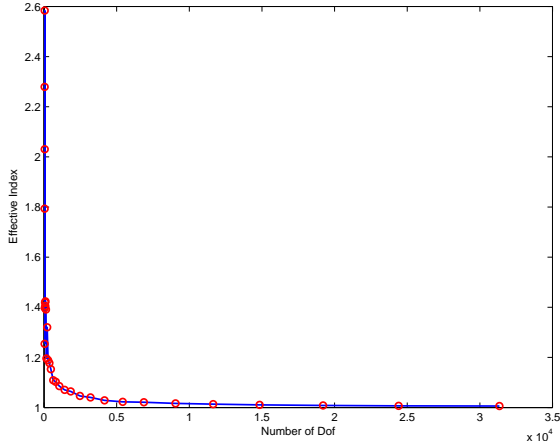


Figure 40: Effective index of Algorithm 3 for Example 5.3

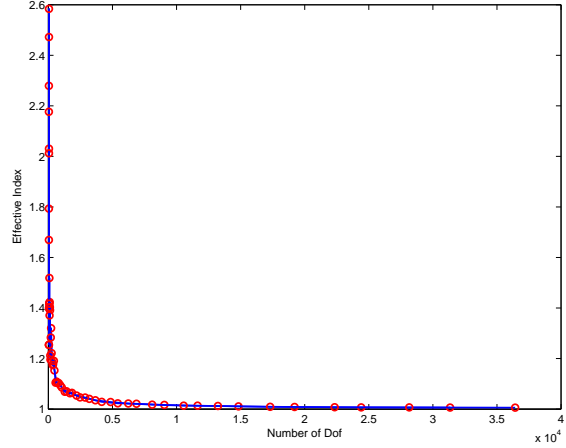


Figure 41: Effective index of Algorithm 4 for Example 5.3

For linear element on structured meshes, using gradient recovery at the last stage achieves similar accuracy as the quadratic element on the same mesh. Therefore, with much reduced cost, the gradient recovery is comparable with the two-stage method on the same mesh.

Algorithms 3 and 4 use recovery type error estimators to adapt the mesh, and have two advantages comparing with the residual based adaptive algorithms. 1) Cost effective. In fact, the recovery based error estimator plays two roles: one is to measure the error, and another is to enhance the eigenvalue approximation. 2) Higher accuracy. Indeed, after recovery enhancement, the approximation error is further reduced.

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**ABSTRACT****RECOVERY TECHNIQUES FOR FINITE ELEMENT METHODS  
AND THEIR APPLICATIONS**

by

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**Major:** Mathematics  
**Degree:** Doctor of Philosophy

Recovery techniques are important post-processing methods to obtain improved approximate solutions from primary data with reasonable cost. The practical usage of recovery techniques is not only to improve the quality of approximation, but also to provide an asymptotically exact posteriori error estimators for adaptive methods. This dissertation presents recovery techniques for nonconforming finite element methods and high order derivative as well as applications of gradient recovery.

Our first target is to develop a systematic gradient recovery technique for Crouzeix-Raviart element. The proposed method uses finite element solution to build a better approximation of the exact gradient based on local least square fittings. Due to polynomial preserving property of least square fitting, it is easy to show that the new proposed method preserves quadratic polynomials. In addition, the proposed gradient recovery is linearly bounded. Numerical tests indicate the recovered gradient is superconvergent to the exact gradient for both second order elliptic equation and Stokes equation. The gradient recovery technique can be used in a posteriori error estimates for Crouzeix-Raviart element, which is relatively simple to implement and problem independent.

Our second target is to propose and analyze a new effective Hessian recovery for continuous finite element of arbitrary order. The proposed Hessian recovery is based on polynomial preserving recovery. The proposed method preserves polynomials of degree  $(k + 1)$  on general unstructured meshes and polynomials of degree  $(k + 2)$  on translation invariant meshes. Based on its polynomial preserving property, we can be able to prove superconvergence of the proposed method on mildly structured meshes. In addition, we establish the ultraconvergence result for the new Hessian recovery technique on translation invariant finite element space of arbitrary order.

Our third target is to demonstrate application of gradient recovery in eigenvalue computation. We propose two superconvergent two-grid methods for elliptic eigenvalue problems by taking advantage of two-grid method, two-space method, shifted-inverse power method, and gradient recovery enhancement. Theoretical and numerical results reveal that the proposed methods provide superconvergent eigenfunction approximation and ultraconvergent eigenvalue approximation. In addition, two multilevel adaptive methods based recovery type a posteriori error estimate are proposed.

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## Education

- Ph.D. in Mathematics, 2015  
Wayne State University, Detroit, Michigan, USA
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Wayne State University, Detroit, Michigan, USA
- M.S. in Computational Mathematics, 2010  
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1. The Maurich J. Zelonka Endowed Mathematics Scholarship, Department of Mathematics, Wayne State University, 2013.
2. Dr. Paul A. Catlin Award for outstanding achievement in the Master's Program, Wayne State University, 2014.
3. The Karl W. and Helen L. Folley Endowed Mathematics Scholarship, Department of Mathematics, Wayne State University, 2015.
4. Graduate Student Traveling Award, Department of Mathematics, Wayne State University, 2015.
5. The Alfred L. Nelson Award for outstanding achievement in the Ph.D. Program, Department of Mathematics, Wayne State University, 2015.

## Selected List of Publications

1. C. Huang, H. Guo, and Z. Zhang, A spectral collocation method for eigenvalue problems of compact integral operators, *J. Integral Equations Appl.*, 25(2013), 79-101.
2. H. Guo, C. Huang, and Z. Zhang, Superconvergence of conforming finite element for fourth order singularly perturbed problems of reaction diffusion type in 1D, *Numer. Methods Partial Differential Equations*, 30(2014), 550-566.
3. H. Guo and Z. Zhang, Gradient Recovery for the Crouzeix-Raviart Element, accepted by *J. Sci. Comput.*
4. H. Guo, Z. Zhang, and R. Zhao, Hessian Recovery for Finite Element Methods, arXiv:1406.3108, submitted.
5. H. Guo, Z. Zhang and R. Zhao, Superconvergent two-grid schemes for Elliptic eigenvalue problems, arXiv:1405.4641, submitted.