

XMM–Newton discovery of the X-ray transient XMMU J181227.8–181234 in the Galactic plane

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ABSTRACT

We report the discovery of an X-ray transient, observed in outburst with *XMM–Newton* on 2003 March 20, and with position $\alpha = 18^{\text{h}}12^{\text{m}}27^{\text{s}}.8$, $\delta = -18^{\circ}12'34''$ (J2000, approximate positional error 2 arcsec). No known source is present at this position and the source was not detected during published *ROSAT* or *ASCA* observations of that region. However, the source may be associated with 1H1812–182 detected by *HEAO 1*, although the error bars on the *HEAO 1* position are very large and the two sources could also be unrelated. Therefore, we name the source XMMU J181227.8–181234. Initially, the source was not detected using the All-Sky Monitor (ASM) on-board the *Rossi X-ray Timing Explorer*, however, reprocessing of the ASM data shows that the source was in fact detected and it was active for about 50 d. The X-ray spectrum of this transient is fitted equally well by an absorbed power law (with a spectral index of 2.5) or multicolour disc blackbody model (with $kT \sim 2$ keV), where we find that the source is highly absorbed. We detect an unabsorbed 0.5–10 keV flux in the range $(2\text{--}5) \times 10^{-9}$ erg cm⁻² s⁻¹, which at a distance of 8 kpc corresponds to a 0.5–10 keV luminosity of $(1\text{--}4) \times 10^{37}$ erg s⁻¹. No pulsations were detected by timing analysis. A colour–colour diagram from ASM data of different accreting objects suggests that the transient is a high-mass X-ray binary, as is also suggested by the high absorption compared to the average interstellar value in the direction of the source. However, the power-law spectral index is far more typical of a low-mass X-ray binary. Thus, we are unable to conclusively identify the nature of the transient. We also report on three sources first detected by the *ASCA* Galactic Plane Survey that are close to this transient.

Key words: X-rays: binaries – X-rays: individual: XMMU J181227.8–181234 – X-rays: individual: AX J1811.2–1828 – X-rays: individual: AX J1812.1–1835 – X-rays: individual: AX J1812.2–1842.

1 INTRODUCTION

X-ray binaries are accreting neutron stars and black holes and can be divided into persistent and transient sources. The transient sources exhibit many orders of magnitude variation in their X-ray luminosity due to similar large fluctuations in the mass accretion rate on to the compact objects. Because of instrument limitations, most of our understanding of these systems is based on the properties and behaviour of the brightest systems ($L_x > 10^{36}\text{--}10^{37}$ erg s⁻¹). Although quite a few lower luminosity systems (with L_x between 10^{34} and 10^{36} erg s⁻¹) are known, they remain enigmatic as they have not been studied in the same systematic way as their brighter counterparts.

We have initiated several programs to systematically investigate the lower luminosity accreting neutron stars and black holes. One of these programs is to search and study very faint X-ray transients (with peak luminosities of 10^{34} and 10^{36} erg s⁻¹) near the Galactic Centre (see Wijnands et al. 2006).

Another program we are performing is to identify the many X-ray sources which were found during several different previous X-ray surveys of the Galactic Centre and the Galactic plane (e.g. Sidoli, Belloni & Mereghetti 2001; Sugizaki et al. 2001; Sakano et al. 2002). In particular, the *ASCA* Galactic Plane Survey (Sugizaki et al. 2001) detected 163 sources in the Galactic plane, and those that are at a distance of 8 kpc approximately fall into the $L_x = 10^{34}\text{--}10^{36}$ erg s⁻¹ range. The majority of the sources remain unidentified due to the poor spatial resolution of the *ASCA* observations. We have started a detailed investigation of the archival *XMM–Newton* and *Chandra*

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data for observations of these *ASCA* sources. Such observations will lead to better positional accuracy which will be used to follow-up the X-ray sources at optical and infrared (IR) wavelengths to try to constrain the nature of the sources. We will also study the X-ray spectral and timing properties which will provide additional insight into the nature of these sources. Whilst studying the archival *XMM-Newton* observations, we discovered a rather bright X-ray transient. In this paper, we report on this discovery, as well as discussing three sources first detected by the *ASCA* Galactic Plane Survey. In future work, we will discuss in detail the results of our archival searches for the *ASCA* and other low-luminosity sources.

2 *XMM-NEWTON* OBSERVATION

XMM-Newton (Jansen et al. 2001) observed a region of the Galactic plane in the direction of $\alpha = 18^{\text{h}}12^{\text{m}}6^{\text{s}}.3$, $\delta = -18^{\circ}22'30''$ on 2003 March 8 (see Fig. 1, ObsID 0152833701). The observation exposure times are ~ 8.7 ks for the EPIC-MOS and ~ 7 ks for the EPIC-pn, no background flares are present. The EPIC instrument was operated in full-frame mode with the medium filter. The observation data files were processed to produce calibrated event lists using the *XMM-Newton* Science Analysis Software (SAS, version 6.5.0).

In the field of view, we detect a bright source with position $\alpha = 18^{\text{h}}12^{\text{m}}27^{\text{s}}.8$, $\delta = -18^{\circ}12'34''$ (J2000). The statistical error in this source position is small (0.05 arcsec) as the source is bright, and so the positional error is dominated by the absolute astrometric accuracy of the telescope, which is estimated to be ~ 2 arcsec (Kirsch 2006). An X-ray source at this position is not identified in catalogues of previous *ROSAT* and *ASCA* surveys of the same region, identifying this source as being transient. However, this source lies close (~ 40 arcsec) to the 95 per cent confidence error box for the source 1H1812–182 detected during the Naval Research Laboratory (NRL) Large Area Sky Survey Experiment with

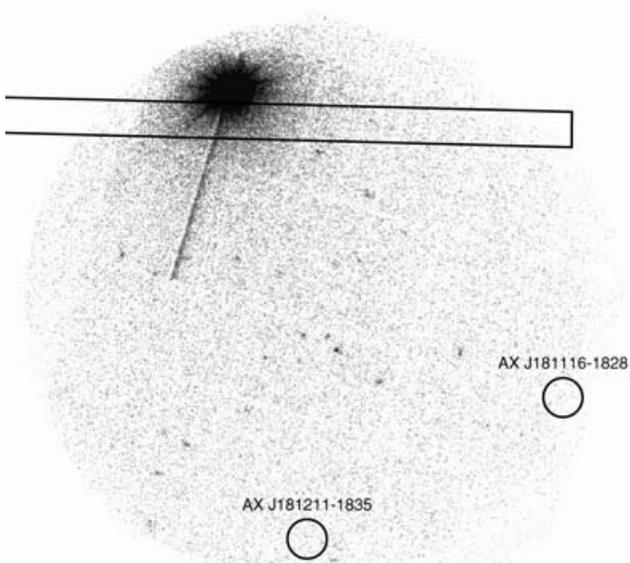


Figure 1. Combined EPIC image of the observed region. The transient is clearly seen in the top-left of the image. The rectangular error box for source 1H1812–182 observed with the *HEAO 1* satellite is shown (Wood et al. 1984). The error circles for sources AX J1811.2–1828 and AX J1812.1–1835 are shown which were detected during the *ASCA* Galactic Plane Survey (Sugizaki et al. 2001).

the *HEAO 1* satellite (Wood et al. 1984). Although the given position of 1H1812–182 ($\alpha = 18^{\text{h}}15^{\text{m}}27^{\text{s}}$, $\delta = -18^{\circ}12'55''$, J2000) is $\sim 0.7^{\circ}$ from the position of the *XMM-Newton* source, it has a large extended error box with an approximate length of 2.1° and width of 0.03° ($=1.8$ arcmin). The *HEAO 1* source is therefore a possible X-ray detection of this *XMM-Newton* transient during a previous outburst, though we note that the large error box of 1H1812–182 does not rule out the possibility that it could well be another X-ray source that is unrelated. 1H1812–182 was detected with a flux of 0.1470 ± 0.0025 counts $\text{cm}^{-2} \text{s}^{-1}$ in the 0.5–25 keV band which corresponds to a 2–10 keV flux of $7.0 \pm 0.1 \times 10^{-10}$ erg $\text{cm}^{-2} \text{s}^{-1}$ assuming a Crab-like spectrum (Wood et al. 1984).

As this new *XMM-Newton* source has not been previously detected, we name it XMMU J181227.8–181234. Initially, the source was not detected using the All-Sky Monitor (ASM) on-board the *Rossi X-ray Timing Explorer (RXTE)*, however, reprocessing of the ASM data showed that the source was in fact detected (see Section 4 for further details).

There is no optical counterpart identified in a Digitized Sky Survey (DSS) image and no counterpart in the Two Micron All-Sky Survey (2MASS) All-Sky Catalog of Point Sources (Cutri et al. 2003) or USNO-B1.0 (Monet et al. 2003) catalogue within 6 arcsec of the source position. The large absorption in this direction makes finding an optical counterpart particularly hard.

3 TIMING AND SPECTRAL ANALYSIS

Extracting a 0.2–10 keV light curve from the EPIC-pn and MOS detectors, we detect no X-ray bursts (see Fig. 2). We use a periodogram analysis (Lomb 1976; Scargle 1982) of the combined EPIC-pn and MOS 0.2–10 keV light curve, with a time bin of 10 s, to search for any pulsations which would help in identifying the nature of the source, but none were found.

This source has a count rate high enough to cause pile up. Therefore, to perform spectral analysis of this data the pile up needs to be corrected for. We extract the spectrum using two different methods. First, from an annulus around the source thus excluding the source core which is most significantly effected by pile up, and secondly, from the readout streak present in the EPIC-pn data. This second method provides a check as to whether the annulus spectrum is significantly affected by the energy dependence of the point spread function.

In the annulus method, we use an annulus of inner radius 15 arcsec and outer radius 45 arcsec. The background was extracted from a

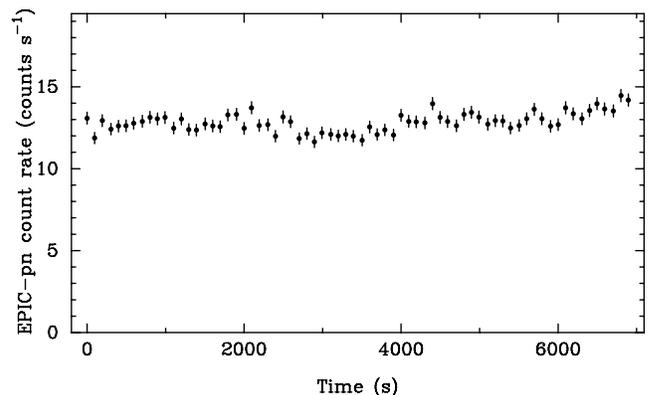


Figure 2. EPIC-pn 0.2–10 keV background subtracted light curve for XMMU J181227.8–181234. Time binning is 100 s.

Table 1. Model fits to the X-ray spectrum of XMMU J181227.8–181234 when fitting the EPIC-pn, MOS1 and MOS2 data simultaneously to a spectrum extracted from the annulus, and when fitting the EPIC-pn spectrum extracted from the readout streak. 1σ errors on the parameters are given.

Parameter	Annulus		Readout streak	
	Power law	diskbb	Power law	diskbb
N_{H} ($\times 10^{22}$ cm $^{-2}$)	12.8 ± 0.3	9.7 ± 0.2	$14.3^{+1.6}_{-1.4}$	$10.8^{+1.0}_{-0.9}$
Spectral index	2.47 ± 0.05		$2.34^{+0.23}_{-0.22}$	
kT (keV)		2.14 ± 0.04		$2.38^{+0.25}_{-0.22}$
0.5–10 keV absorbed flux ($\times 10^{-10}$ erg cm $^{-2}$ s $^{-1}$)	7.7 ± 0.1	7.6 ± 0.1		
0.5–10 keV unabsorbed flux ($\times 10^{-9}$ erg cm $^{-2}$ s $^{-1}$)	5.3 ± 0.4	1.81 ± 0.02		
χ^2_{ν} (d.o.f.)	1.4 (1459)	1.3 (1458)	0.99 (124)	0.90 (124)

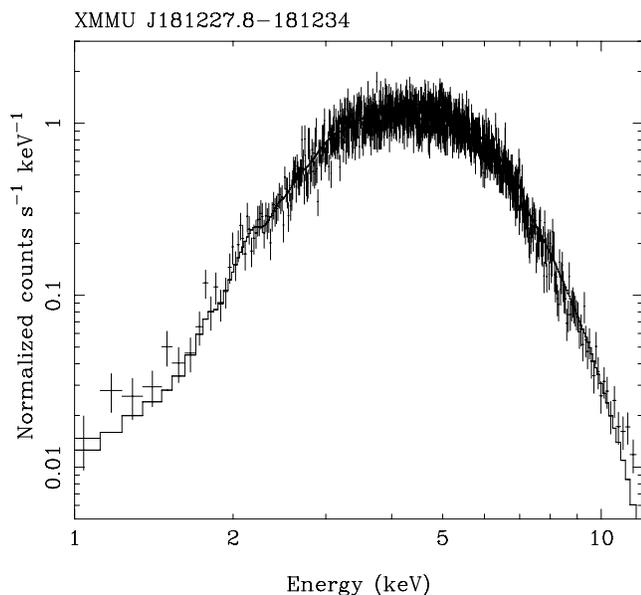


Figure 3. EPIC-pn spectrum for XMMU J181227.8–181234 with best-fitting absorbed multicolour disc blackbody model. The spectrum was extracted from an annulus around source to reduce the effects of pile up.

3-arcmin circular region close to the transient and containing no detected sources. For the MOS cameras, only single-pixel events (pattern = 0) were selected as recommended for dealing with pile up by Molendi & Sembay (2003). We use the SAS tool EPATPLOT to make sure that no pile up was present. Response matrices and ancillary files were created using RMFGEN and ARFGEN.

The data were binned to have a minimum of 25 counts per bin. We use XSPEC (v.11) (Arnaud 1996) to fit the EPIC-pn, MOS1 and MOS2 data simultaneously. Two different models were fit to the data – an absorbed power-law model and an absorbed multicolour disc blackbody model diskbb. Table 1 gives the best-fitting parameters for both models and the best-fitting diskbb model is shown in Fig. 3. We find that both models give reasonably good fits. From the spectrum and model fits it is apparent that the source is highly absorbed suggesting that it is not a nearby lower luminosity source. The power-law fit gives a spectral index, $\Gamma = 2.47 \pm 0.05$, where as the diskbb fit gives quite a high characteristic temperature of $kT = 2.14 \pm 0.04$ keV. These model fits lead to an unabsorbed 0.5–10 keV flux of $5.3 \pm 0.4 \times 10^{-9}$ and $1.81 \pm 0.02 \times 10^{-9}$ erg cm $^{-2}$ s $^{-1}$ for the power law and diskbb models, respectively. Assuming that this source is at a

distance of 8 kpc, this corresponds to a 0.5–10 keV luminosity of $4.1 \pm 0.3 \times 10^{37}$ erg cm $^{-2}$ s $^{-1}$ for the power-law model and $1.4 \pm 0.02 \times 10^{37}$ erg cm $^{-2}$ s $^{-1}$ for the diskbb model, quite typical of other X-ray transients.

The readout streak is only significantly present in the pn image, thus we only extract the readout streak spectrum from the pn data. Binning the data to have a minimum of 25 counts per bin, we again fit in XSPEC using both an absorbed multicolour disc blackbody model and an absorbed power-law model. The best-fitting parameters are given in Table 1, and are seen to be consistent with the parameters from the annulus spectrum. We do not give fluxes for the readout streak spectrum as the photons only hit the CCD in the streak region when the CCD is being readout, and so the exposure time is incorrect.

4 RXTE ALL SKY MONITOR OBSERVATIONS

Initially, the detection of this transient with the RXTE ASM was missed, due to the mask that is put around bright sources when looking at deep sky maps. In this case the bright persistent low-mass X-ray binary GX 13+1 (4U 1811–17) is $\sim 1.2^\circ$ away from the transient. Given the position and time of the outburst, we detect the transient with the RXTE ASM. We have also derived the source position from first principles, given the outburst time interval, that is consistent with the XMM–Newton position. Fig. 4 shows the RXTE ASM light curve for this source with the outburst clearly visible. The outburst lasts for approximately 50 d, at least, and there is only one outburst detected during the lifetime of RXTE. The XMM–Newton observation occurs after the peak of the outburst. The ASM count rate at the time of the XMM–Newton observation is approximately 3 counts s $^{-1}$. Using PIMMS, and assuming this count rate and the best-fitting power-law spectral model from Table 1, this corresponds to an unabsorbed 0.5–10 keV flux of 6×10^{-9} erg cm $^{-2}$ s $^{-1}$, similar to that detected with XMM–Newton.

The coarse shape of the spectrum in X-ray binary systems may help to reveal the type of the source. To investigate this question, we use the hardness ratios derived from the source count rates in the three standard ASM energy bands: A (1.5–3 keV), B (3–5 keV) and C (5–12 keV). The hardness ratios are then defined as $\text{HR1} = B/A$ and $\text{HR2} = C/B$. We expect that HR1 is sensitive to both the spectral shape of the X-ray source and the amount of interstellar absorption, while HR2 is more sensitive to the spectral shape.

The ‘dwell-by-dwell’ data for a particular source in the ASM archive provides independent measurements from each 90-s exposure by one of the three ASM cameras. We select cameras 2 and 3

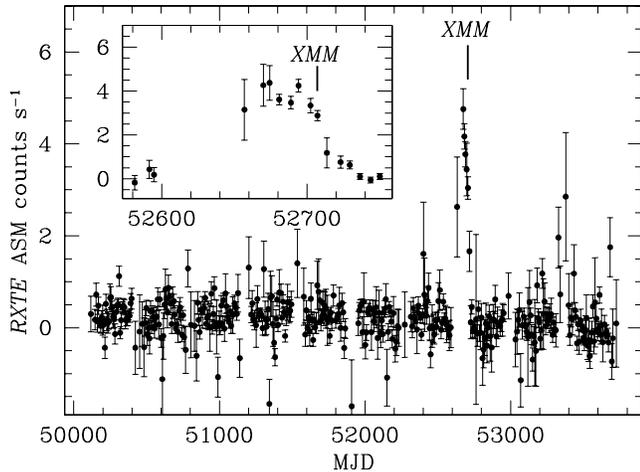


Figure 4. 7-d averaged *RXTE* ASM light curve for XMMU J181227.8–181234, with the time of the *XMM–Newton* observation marked. The inset shows the outburst in more detail, again with the time of the *XMM–Newton* observation marked. There is clearly only one outburst detected during the lifetime of *RXTE*.

only for this exercise. Camera 1 is excluded because it has a secular gain drift, caused by a pinhole Xe leak discovered before *RXTE* launch, and it is rigidly calibrated to the other cameras for the specific spectral shape of the Crab Nebula. For other spectral shapes, the hardness ratios from camera 1 exhibit a time dependence that would add systematic noise in the present study.

For persistent sources [six atolls, six Z sources, four classical X-ray pulsars and the two Large Magellanic Cloud (LMC) black holes], we compute the ASM count rates (weighted mean) in each energy band over a long time interval, arbitrarily chosen to be MJD 51000–52000. We also sample the persistent black hole binary Cyg X–1 in each of two X-ray states: soft (MJD 52250–52450) and hard (MJD 50720–50920). Count rates are determined for X-ray transients (eight atolls, six classical X-ray pulsars and seven black holes) for the 7-d time interval that captures the peak of the outburst. This is evaluated by finding the maximum flux in the weighted mean, while sliding a 7-d window along the light curve sampled in 1-d bins.

From this analysis, we find that a colour–colour diagram using these two ASM hardness ratios can give a good indication of source type (see Fig. 5). We note that HR1 is quite uncertain for this transient due to the low number of counts in the 1.5–3 keV band caused by the high absorption towards this object. The HR2 ratio, however, is less susceptible to interstellar absorption. We find that HR2 = 2.45, suggesting that this source is a high-mass X-ray binary pulsar.

5 DISCUSSION

5.1 The nature of XMMU J181227.8–181234

From our spectral fits to XMMU J181227.8–181234, we find a high absorption towards this source. Using the *FTOOL* nH (Dickey & Lockman 1990), we get $N_{\text{H}} \sim 2 \times 10^{22} \text{ cm}^{-2}$ in the direction of this source – a factor of at least 5 different from the value from spectral fitting. This may suggest that the source has significant intrinsic absorption which would be indicative of a high-mass X-ray binary (HMXB). *INTEGRAL* observations of the Galactic plane have discovered sources with high intrinsic absorption that

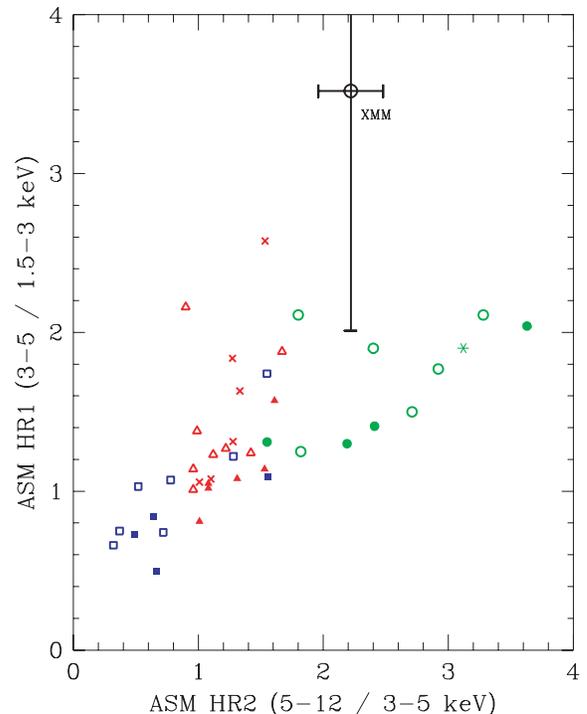


Figure 5. *RXTE* ASM colour–colour diagram. The shape (and also the colour) denote source types, an open symbol is a transient; filled symbols are persistent sources. The squares (blue) are black hole binaries with two states for Cyg X–1. Triangles (red) denote the bursters. Crosses (red) denote the Z sources (all of these are persistent). Circles (green) denote the pulsars. The green asterisk is the high-mass X-ray binary 4U 1700–377.

are likely HMXBs (e.g. Patel et al. 2004; Walter et al. 2004; Lutovinov et al. 2005a,b; Bodaghee et al. 2006). The intrinsic absorption is due to the accretion process in the majority of HMXBs in which the compact objects accrete matter from the stellar wind of a young companion star. The stellar wind can cause significant absorption. In low-mass X-ray binaries (LMXBs), however, accretion occurs via Roche lobe overflow through the inner Lagrangian point, thus in the majority of cases no material is available to cause significant absorption. Only in high inclination systems significant internal absorption might be present. However, for this transient we do not find any of the characteristic count rate variations (i.e. dips and eclipses) which would indicate a high inclination.

HMXB pulsars typically show significant slow pulsations, however, there are notable exceptions, for instance, the well-known HMXB, 4U 1700–377 (marked in Fig. 5 as an asterisk), has no detectable pulsations. In addition, HMXBs can typically be fitted by an absorbed power law with spectral index, $\Gamma \sim 1$. There are some lower luminosity HMXBs that have spectra that are softer (with $\Gamma \sim 2$ –2.5, see e.g. Reig & Roche 1999), though these sources are persistent and have $L_x \sim 10^{34} \text{ erg s}^{-1}$, where as XMMU J181227.8–181234 is transient in nature and is significantly brighter (with L_x a few times $10^{37} \text{ erg s}^{-1}$). XMMU J181227.8–181234 has no pulsations detected in the light curve and it is fit by a power-law spectral index of $\Gamma = 2.5$ which is far more typical of LMXBs. Thus, although the ASM colours and possible intrinsic absorption suggest that this source may be a HMXB, the power-law index is more suggestive of a LMXB. We are therefore unable to conclusively determine whether this transient is either a HMXB or a LMXB, and an IR follow-up is needed to search for a counterpart at this wavelength. As no optical or IR counterpart was found with USNO-B1.0 or 2MASS this

suggests that the source is not a nearby lower luminosity object which is supported by the high measured N_{H} .

5.2 ASCA Galactic Plane Survey sources

The ASCA Galactic Plane Survey (Sugizaki et al. 2001) detected 163 discrete sources in the 0.7–10 keV energy band within the central region of the Galactic plane with a spatial resolution of ~ 3 arcmin and a positional error of ~ 1 arcmin. The limiting flux of the survey in the 0.7–10 keV band was $\sim 3 \times 10^{-13}$ erg cm $^{-2}$ s $^{-1}$ which corresponds to a luminosity of $\sim 2 \times 10^{33}$ erg s $^{-1}$ assuming a distance of 8 kpc. We use PIMMS to predict the EPIC-pn on-axis count rate assuming such a flux, a column density $N_{\text{H}} = 2 \times 10^{22}$ cm $^{-2}$, and with both a power law with spectral index, $\Gamma = 0$ and spectral index, $\Gamma = 3$. We calculate the EPIC-pn on-axis count rate for this to be approximately 0.014 counts s $^{-1}$ when $\Gamma = 0$ and 0.041 counts s $^{-1}$ when $\Gamma = 3$. Thus, for an exposure time of 7 ks, the lower limit of ASCA Galactic Plane Survey corresponds roughly to between 100 and 290 counts on-axis (0.5–10 keV) for typical spectra. We note that off-axis the effective exposure time drops rapidly due to vignetting. By around 15 arcmin off-axis, the vignetting factor is about 0.3 (see fig. 13 of the *XMM–Newton* User’s Handbook¹), thus the lower limit drops to between 30 and 90 counts. Thus we expect to detect any persistent point sources from the ASCA Galactic Plane Survey with this *XMM–Newton* observation.

Two of the sources (AX J1811.2–1828 and AX J1812.1–1835) from this survey lie within the field of view of this *XMM–Newton* observation and another of the sources (AX J1812.2–1842) is in an *XMM–Newton* observation of a nearby region. We discuss the *XMM–Newton* observation of these sources here.

5.3 AX J1811.2–1828

AX J1811.2–1828 was detected for the first time with ASCA (Sugizaki et al. 2001). No spectral model was fit to data from AX J1811.2–1828, though these authors quote a 2–10 keV count rate of 6.3 counts ks $^{-1}$ GIS $^{-1}$. Depending on the spectral model, this converts to a 2–10 keV flux between 4×10^{-13} and 11.5×10^{-13} erg cm $^{-2}$ s $^{-1}$.

The majority of the error circle for this source falls off the edge of the MOS1 and MOS2 detectors, and so we can only use the pn detector to search for a detection in this *XMM–Newton* observation. No bright source is detected within the error circle, however, we tentatively detect a possible faint source on the edge of the error circle for this source with 40 ± 9 net counts in the 0.2–12 keV range for the pn detector. To check whether this source is in fact real, we analyse two additional *XMM–Newton* archival observations (ObsID 0152834201 and 0152835401) in the region of this source. In both observations we detect a source with a position $\alpha = 18^{\text{h}}11^{\text{m}}13^{\text{s}}.3$, $\delta = -18^{\circ}27'57''$ (J2000), and a positional error of 2.2 arcsec (including both statistical and boresight errors). The position of the potential source in the original data is consistent with that of the source in the two extra data sets and thus the source is detected in all three *XMM–Newton* observations. In Fig. 6, the improved position of this source can be seen in a combined image from all three *XMM–Newton* observations. In ObsID 0152835401 the source is on the very edge of the field-of-view and so we only analyse the spectrum from the ObsID 0152834201 data, where the source is bright enough to bin the spectrum with enough counts to validly use χ^2 statistics in the spectral fitting.

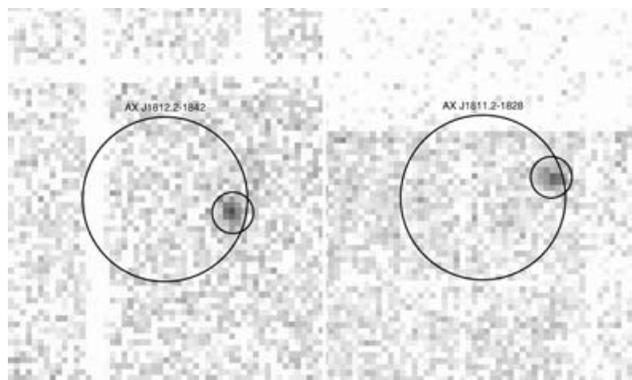


Figure 6. The two ASCA sources AX J1812.2–1842 (left-hand panel) and AX J1811.2–1828 (right-hand panel) are clearly detected with *XMM–Newton* with a much improved source position. The larger circles are the ASCA error circles for these sources, whilst the smaller circles are the 15-arcsec source extraction regions we use here.

We extract the spectrum from the MOS1, MOS2 and pn detectors using a circular source region with radius 15 arcsec, and we extract the background from an annulus around the source with inner and outer radii of 1 and 2 arcmin, respectively. We bin the spectrum to 15 counts per bin, and fit the spectra simultaneously with an absorbed power-law model using XSPEC. The best-fitting model has parameters $N_{\text{H}} = 0.6^{+0.6}_{-0.3} \times 10^{22}$ cm $^{-2}$ and $\Gamma = 1.5^{+0.7}_{-0.3}$. The reduced χ^2 for this fit is 0.8. For comparison with the ASCA detection, we calculate the 2–10 keV flux, which we find to be $2.8^{+0.6}_{-0.4} \times 10^{-13}$ erg cm $^{-2}$ s $^{-1}$ and the unabsorbed 2–10 keV flux is $2.9 \pm 0.6 \times 10^{-13}$ erg cm $^{-2}$ s $^{-1}$. This flux is slightly lower than the flux detected by ASCA. From the spectral fits this source has a reddening smaller than the interstellar absorption in this direction, thus it is likely closer than a distance of 8 kpc. As the distance to this source is unknown it is hard to determine what type of source this is with just these spectral fits, and the power-law index does not rule out a HMXB, LMXB or cataclysmic variable. We conclude that we have detected AX J1811.2–1828 with *XMM–Newton* with a 2–10 keV flux which was slightly lower than that seen by ASCA.

We find several possible optical and IR counterparts to this X-ray source in the USNO-B1.0 and 2MASS catalogues. The USNO-B1.0 source named USNO-B1.00715–0555562 is 1.8 arcsec away and has a B2 magnitude of 19.8 and a R2 magnitude of 18.0. USNO-B1.00715–0555563 is 2.6 arcsec away, and has B2 and R2 magnitudes of 19.3 and 16.5, respectively. USNO-B1.00715–0555571 is 2.7 arcsec away and has B2 and R2 magnitudes of 19.11 and 15.35, respectively. The 2MASS source named 2MASS J18111351–1827565 is a distance of 3.1 arcsec away from the X-ray source and has *J*, *H* and *K* magnitudes of 12.3, 11.2 and 10.9, respectively. Optical/IR follow-up spectroscopy of these sources is needed to determine the spectral type, and hence the likely nature of this source.

5.4 AX J1812.1–1835

Sugizaki et al. (2001) identify AX J1812.1–1835 with the supernova remnant (SNR) G12.0–0.1 ($\alpha = 18^{\text{h}}11^{\text{m}}53^{\text{s}}.3$, $\delta = -18^{\circ}35'54''$; Clark, Caswell & Green 1975; Altenhoff et al. 1979; Kassim 1992; Dubner et al. 1993; Green et al. 1997; Green 2004). Their spectral fits to AX J1812.1–1835 lead to a flux for the source of $\sim 9 \times 10^{-13}$ erg cm $^{-2}$ s $^{-1}$ in the 0.7–10 keV band. No source is detected within the error circle of AX J1812.1–1835. To determine an upper

¹ Available from <http://xmm.vilspa.esa.es>

limit for this source, we extract the count rate from the faintest source detected at a similar off-axis angle. Using an extraction radius of 15 arcsec for a source located at $\alpha = 18^{\text{h}}12^{\text{m}}43^{\text{s}}.7$, $\delta = -18^{\circ}31'58''$ (J2000), we get 51 ± 10 background-subtracted counts in the 0.2–12 keV range for the pn detector. A source-free annulus between 40 and 100 arcsec was used to determine the background counts. The average effective exposure in the source extraction region is 2429 s, giving a net count rate of 0.021 ± 0.004 counts s^{-1} . The average interstellar absorption in the direction of this source is 2×10^{22} cm^{-2} . Using PIMMS and assuming this value, we get a 0.5–10 keV flux for the source in the range $(1.5\text{--}4.2) \times 10^{-13}$ $\text{erg cm}^{-2} \text{s}^{-1}$ depending on the power-law spectral index which we vary from 0 to 3. Thus, if AX J1812.1–1835 is a point source, and has not varied in flux since the ASCA observation it should be detected, which it is not.

There are several possible explanations for the lack of detection. First, if this source is not associated with the SNR, then the flux from this source has dropped by a factor of at least 2 between the ASCA and XMM-Newton observations, identifying the source as variable. However, it is likely that this source is associated with the SNR G12.0–0.1. Studying the VLA 1465 MHz radio observation of this SNR (see fig. 3 from Dubner et al. 1993), we find that although the bright thermal source is located several arcmin away from AX J1812.1–1835, the ASCA source does lie within the incomplete shell to the east. If AX J1812.1–1835 is associated with this SNR, then the emission will be diffuse, and therefore below our detection limit.

5.5 AX J1812.2–1842

AX J1812.2–1842 was first detected by the ASCA Galactic Plane Survey. No spectral model was fit to data from AX J1812.2–1842 by Sugizaki et al. (2001), though these authors quote a 2–10 keV count rate of 5.4 counts $\text{ks}^{-1} \text{GIS}^{-1}$. Depending on the spectral model, this converts to a 2–10 keV flux between 3×10^{-13} and 10×10^{-13} $\text{erg cm}^{-2} \text{s}^{-1}$.

Although not in the field-of-view of the XMM-Newton observation containing the transient, it is in the data with ObsID 0152835401. The source is detected with a position $\alpha = 18^{\text{h}}12^{\text{m}}10^{\text{s}}.5$, $\delta = -18^{\circ}42'10''$ (J2000) and a positional error of 2.3 arcsec (including both statistical and boresight errors). The improved source position can be seen in a combined MOS1, MOS2 and pn image (Fig. 6, left-hand panel).

We extract the spectrum from the MOS1, MOS2 and pn detectors using a circular source region with radius 15 arcsec, and we extract the background from an annulus around the source with inner and outer radii of 1 and 2 arcmin, respectively. We bin the spectrum to 15 counts per bin for the MOS1 and MOS2 data and 40 counts per bin for the pn data. The spectra are fit simultaneously with an absorbed power-law model using XSPEC. The best-fitting model has parameters $N_{\text{H}} = 0.5_{-0.5}^{+1.3} \times 10^{22}$ cm^{-2} and $\Gamma = 0.0_{-0.5}^{+0.7}$. The reduced χ^2 for this fit is 0.7. For comparison with the ASCA detection, we calculate the 2–10 keV flux, which we find to be $9.4_{-1.7}^{+2.0} \times 10^{-13}$ $\text{erg cm}^{-2} \text{s}^{-1}$ and the unabsorbed 2–10 keV flux is $9.6_{-1.2}^{+1.8} \times 10^{-13}$ $\text{erg cm}^{-2} \text{s}^{-1}$. This flux is consistent with the source flux detected by ASCA. The low N_{H} suggests that the source may be closer than 8 kpc, however, we note the large uncertainty in this parameter. The power-law spectral index suggests that this source could be a cataclysmic variable, though we cannot be conclusive.

A search for optical and IR counterparts finds no sources within the errors for the position. The nearest optical source, USNO-B1.00712–0536743, is a distance of 3.5 arcsec away, and has B2 and R2 magnitudes of 18.4 and 17.8, respectively. The nearest IR

source, 2MASS J18121061–1842133, is a distance of 4.3 arcsec, and has *J*, *H* and *K* magnitudes of 14.8, 12.3 and 10.9, respectively.

6 CONCLUSIONS

We have discovered an X-ray transient with position $\alpha = 18^{\text{h}}12^{\text{m}}27^{\text{s}}.8$, $\delta = -18^{\circ}12'34''$ (J2000) detected by XMM-Newton and the RXTE ASM. This source was not previously detected by either ROSAT or ASCA, though it may be associated with 1H1812–182 detected by HEAO 1 which has a large extended error box that is close to the transient source position. The X-ray spectrum of this transient is fitted equally well by an absorbed power law or multicolour disc blackbody model, where we find that the source is highly absorbed. We detect an unabsorbed 0.5–10 keV flux in the range $(2\text{--}5) \times 10^{-9}$ $\text{erg cm}^{-2} \text{s}^{-1}$, which at a distance of 8 kpc corresponds to a 0.5–10 keV luminosity of $(1\text{--}4) \times 10^{37}$ erg s^{-1} , quite typical of other Galactic X-ray binary transients. Using a periodogram analysis, we searched for pulsations to help with source identification, but none were found. The high column density to the source and lack of optical/IR counterpart in the USNO-B1.0 and 2MASS catalogues suggests that the source is not a nearby source with lower luminosity. A colour–colour diagram from RXTE ASM data of different accreting objects suggests that the transient is a high-mass X-ray binary, as is also suggested by the high absorption compared to the average interstellar value in the direction of the source. However, the power-law spectral index is far more typical of a low-mass X-ray binary. Thus, we are unable to conclusively identify the nature of the transient. Follow-up IR observations are required to help determine the nature of this source.

We also report on three sources first detected by the ASCA Galactic Plane Survey that lie within the field-of-view of this observation. We do not detect AX J1812.1–1835, though this is likely because it is associated with the SNR G12.0–0.1 and so the diffuse emission would not be bright enough to be detected with this observation. For AX J1811.2–1828, we detect a source towards the edge of the error circle with a 2–10 keV flux of $2.8_{-0.4}^{+0.6} \times 10^{-13}$ $\text{erg cm}^{-2} \text{s}^{-1}$, slightly lower than the ASCA observation of this source. The spectrum is fit with an absorbed power law with $N_{\text{H}} = 0.6_{-0.3}^{+0.6} \times 10^{22}$ cm^{-2} and $\Gamma = 1.5_{-0.3}^{+0.7}$. We also detect AX J1812.2–1842 with a 2–10 keV flux of $9.4_{-1.7}^{+2.0} \times 10^{-13}$ $\text{erg cm}^{-2} \text{s}^{-1}$, consistent with the ASCA observation of this source. The spectrum is fit with an absorbed power law with $N_{\text{H}} = 0.5_{-0.5}^{+1.3} \times 10^{22}$ cm^{-2} and $\Gamma = 0.0_{-0.5}^{+0.7}$. We cannot conclusively determine the nature of AX J1811.2–1828 or AX J1812.2–1842 with these observations and optical/IR follow-ups are needed. Many sources from the ASCA Galactic Plane Survey remain unidentified in nature, and require further X-ray observations with XMM-Newton and Chandra to provide subarcsec source positions allowing for optical/IR follow-up to aid with source identification as well as providing further X-ray temporal and spectral information.

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